

FIVE-BRANE CONFIGURATIONS,  
CFT's and the  
STRONG-COUPING PROBLEM

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1. General context

- string theory  $\Rightarrow$  gauge interactions & gravity
  - ↳ tradition  $\rightarrow$  string phenomenology  $\rightarrow$  SM and beyond
  - $\rightarrow$  string gravity  $\rightarrow$  search for backgrounds
  - string cosmology  $\leftarrow$  time-dependent  $\rightarrow$
- quantum effects, singularities, bh  $\leftarrow$  short-distance behaviour  $\leftarrow$
- Branes  $\rightarrow$  new landscape of interests

Fundamental objects : F1 NSS DP

- conceptual achievement → understand the string spectrum beyond perturbation theory
  - unify the string vacua (dualities), introduce M-theory  $\ni M2 \& M5$
- merging of phenomenology and gravity motivations and methods
  - branes act like "impurities"
  - alter the string spectrum
  - modify the background
- Use exact CFT approaches  
(orbifold, fermionic, ...)
- $\rightarrow \text{SM}$
- find old'  $\sigma$ -models and interpret them geometrically  $\rightarrow \text{GRAVITY}$
- Find brane configurations, their spectrum, their geometry and CFT interpretation
- concrete and celebrated example: the Randall-Sundrum scenario
- other drawback: DECOUPLING LIMITS  $\neq$  low-energy
  - string spectrum is highly constrained (GSO, modular invariance, ...) and contains gravity and gauge interactions in an intricate way
  - much like orbifold fixed points, branes carry part of the spectrum i.e. one among many sectors, all necessary (see exact CFT)
  - however . . .

... there are limits where

- the spectrum is dominated by excitations leaving **on** the brane, **not** described by QFT
  - gravitational sector decouples
  - holography is at work
- ... occurring in regions of space-time where  $g_{\text{string}}$  diverges!

- motivation: illustrate these issues in the celebrated example of **5-branes** and show how to overcome the strong-coupling problem within exact CFT framework.

## 2. The 5-brane solutions in type II [the original CHS<sup>91</sup>]

- $S_{\text{II 10-dim}}^{\text{Low-energy}} = \frac{1}{2k_{10}^2} \int d^{10}x \sqrt{-g^{(10)}} \left( R^{(10)} - \frac{1}{2} (\partial \phi_j)^2 - \frac{1}{12} e^{-\gamma \phi_j} H^2 \right)$

→  $H = dB$        $\gamma = 1$  NS-NS ;  $\gamma = -1$  R-R (IIB)

→ string coupling       $g_s = \exp \phi_j$

- $\beta = 0$  [ $O(\lambda')$ ] → five-brane type of solutions:

$$\frac{ds^2}{d'} = h(z)^{-1/4} \underbrace{(-dt^2 + d\vec{z}^2)}_{\sim 5+1} + h(r)^{1/4} \underbrace{(dr^2 + r^2 d\Omega_3^2)}_{\sim 4}$$

wald volume<sub>eff</sub>  
6+1 - Poincaré

transverse  $\perp$   
 $SO(4)$

$$\rightarrow H = -r^3 h' \times \text{volume form of } S^3$$

$$\square_{(4)} h = 0$$

$$\rightarrow \phi_f = \frac{X}{2} \log h$$

$$Q_5 \text{ 5-branes at } r=0 \quad \leftarrow \quad h = h_0 + Q_5/r^2$$

$\frac{1}{2}$  \* if  $Q_5 = 0$  flat space exact CFT (torus)

generic  $h_0, Q_5$  \* if  $h_0 = 0$   $\begin{cases} 6+1 \text{ flat space} \oplus S^3 & \text{rad} = \sqrt{Q_5} \\ \text{linear dilaton in one direction} \end{cases}$

( $\gamma=1$ ) EXACT CFT

$$N=4 \text{ superconformal} \rightarrow SU(2)_{Q_5-2} \text{ WZW} \times U(1) \times \text{Flat}_{5+1}$$

[Antoniadis, Ferrara,  
Kounnas '94;  
Kiritsis, Kounnas '95]

spectrum known (combination of characters  
of  $SU(2)$  and Liouville)  $\rightarrow$  discrete  $\rightarrow$  short  $N=4$   
 $\rightarrow$  continuous  $\rightarrow$  long(massive)

“brane”

“bulk”

$\rightarrow$  at large  $k$

survive in the decoupling limit: LST

- The string coupling:



NS  $\rightarrow$  strong coupling at  $r \sim 0$

D  $\rightarrow$  If  $h_0 = 0$ , strong coupling at  $r \rightarrow \infty$

- This behaviour is not pathological: consistent with LST

→ type IIB  $\nu \rightarrow 0 \equiv$  IR in the LST of NS5  
III

SYM  $(1,1)_{6D}$  ~ free

dual by holography

D5-brane background string theory in the above background  
at  $\underline{g_s \rightarrow 0}$        $\text{dual}$       at  $\underline{g_s \rightarrow \infty}$   
 ↳ no exact CFT techniques

→ type IIA LST in IR  $\equiv$  superconformal  $(2,0)_{6D}$  fixed-point  
dual by holography

M5-spread over  $0 \rightarrow 0 \cdots 0$

$\text{AdS}_7 \times S^4$

### 3. Scanning the web of NS5 (or D5)

idea: distribute the branes in the transverse space

better understanding of the strong-coupling ↳ new backgrounds  
new decoupling limits  
new web of CFT's

well known example: S-branes on a circle

[Sfetsos '98,

Gireon et al '99  
Bakano et al '00]

- in the LST higgs  $\rightarrow \chi_N$

- simpler setup: continuous distribution in the NH limit (decoupling)



(6)

- geometry

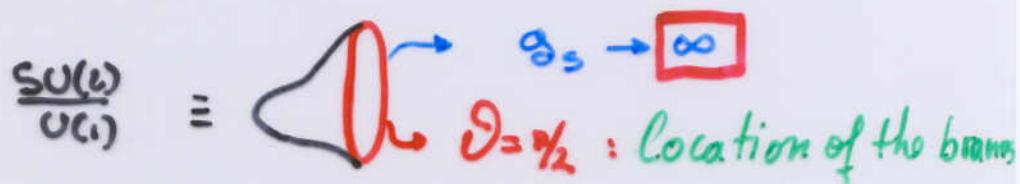
$$ds^2 = Q_5 \left\{ d\rho^2 + d\vartheta^2 + \frac{\tan^2 \vartheta d\psi^2 + \tanh^2 \rho d\zeta^2}{1 + \tanh^2 \rho \tan^2 \vartheta} \right\}; H = dB$$

$$g_s^{-2} = e^{-2\phi} = e^{-2\phi_0} \left( \cosh^2 \rho \cos^2 \vartheta + \sinh^2 \rho \sin^2 \vartheta \right) \text{ for NS5}$$

- rich content in terms of exact CFT's (NS 5's)

→  $\rho = \text{cons.}$  squashed sphere  $SU(2)_{Q_5^2}$ -deformed  $SU(2)$   
 ↳ dynamical deformation

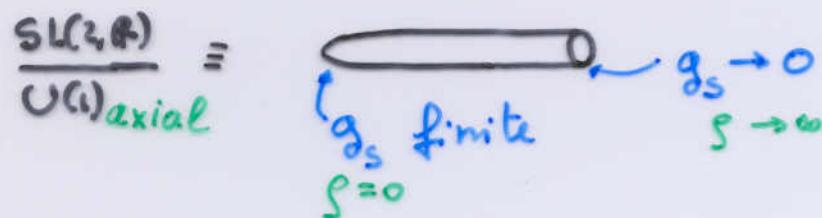
↳  $\rho \rightarrow 0$  Bell  $\times U(1)$



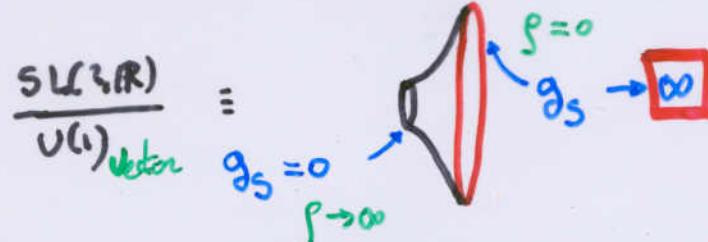
↳  $\rho \rightarrow \infty$   $SU(2)_{Q_5-2}$  undeformed

→  $\vartheta = \text{cons.}$

↳  $\vartheta = 0$  cigar  $\times U(1)$



↳  $\vartheta = \pi/2$  trumpet  $\times U(1)$



- strong coupling avoided by chain of  $T$ -dualities +  $\mathbb{Z}_{Q_5}$

$$\hookrightarrow \mathbb{Z}_{Q_5}^{\mathbb{C}} \longrightarrow T^{\mathbb{C}} \longrightarrow \mathbb{Z}_{Q_5}^4 \longrightarrow T^4 \downarrow$$

$N=4$  superconformal world-sheet

$$\left\{ \begin{array}{c} \text{cigar} \times \text{bell} \\ \frac{SL(2, \mathbb{R})_{Q_5+2}}{U(1)_{\text{axial}}} \times \frac{SU(2)_{Q_5-2}}{U(1)} \end{array} \right\} \text{no strong coupling}$$

- which is also  $T$ -dual to the original  $SU(2)_k \times U(1)_Q$   
→ needs further investigation

#### 4. Adding F1 (or D1) & null deformation

(see talk by Dan Freed)

- NS5/F1 (IIA, B) or D5/D1 (IIB)

$N=2$  backgrounds ( $N=4$  in the NHG)

$$AdS_3 \times S^3 \times T^4 \xleftarrow{\text{decoupling limit}}$$

$$\begin{matrix} M^2 \rightarrow 0 & V \sim d'^{1/2} v & U = r/d' \text{ fixed} \\ \uparrow & \nearrow & \downarrow \\ \end{matrix}$$

$$d' \rightarrow 0$$

$$U = r/d' \text{ fixed}$$

$$v \text{ fixed}$$

marginal  
null deformation

$$(J_1 + J_3)(\bar{J}_1 + \bar{J}_3)$$

another decoupling limit

$$d' \rightarrow 0$$

$$U = r/d' \text{ fixed}$$

dilution  $g_s, d'$  fixed

of D1's  $v, d'^{1/2}$  fixed

$$M^2 = \frac{g_s N_2}{d' U}$$

$$ds^2 = Q_5 \left( \frac{du^2}{U^2} + \frac{dx^2 - dt^2}{U^2 + \eta_M^2} + d\Omega_3^2 \right) + d(T^4)^2$$

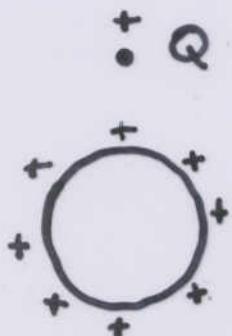
$$g_s^2 = e^{2\phi} = \frac{1}{g_{10}^2} \frac{u^2}{u^2 + 1/M^2} \quad (\text{NS5/F1})$$

- at  $M^2 \rightarrow 0$  (F1 very diluted):  $SU(2)_k \times U(1)_Q$   
 $g_s^2$  diverges at  $u \rightarrow \infty$  (horizon)
- at finite  $M^2$ : no divergence

[Israël, Kounnas, Petropoulos]

## 5. Spreading the NS5 or D5 over an $S^3$

- the guide: electrostatic analogue



$$V = Q/r \quad \text{divergent}$$

$$V = Q/r \quad \text{outside} \quad \text{finite}$$

$$V = Q/r_0 \quad \text{inside}$$

- translated in our brane system: distribute the  $Q_s$  branes over an  $S^3$ , at radius  $r=1$  instead [Keritsis, Kounnas, Petropoulos, Rizos '02] of having them all at  $r=0$

- Eq.  $\partial_z^2 h + 2\partial_z h = 0 \quad (r = \exp z)$

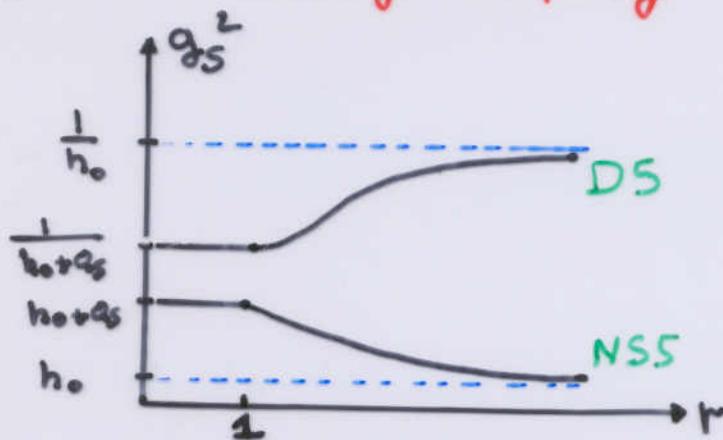
becomes  $\partial_z^2 h + 2\partial_z h = -2 \frac{Q_s}{\alpha'} S(z)$

from  $S = S_{\text{II load}}^{\text{low-energy}} + S_{\text{brane}}^{\text{dist}}$   $\hookrightarrow$  source at finite  $z$

$$S_{\text{sb}}^{\text{dist}} = -\frac{Q_s T_s}{2\pi r} \int d^4x \sin^2 \sin \varphi S(z) \sqrt{-g(z)} e^{-\frac{\phi}{2}} \quad \begin{matrix} \text{NS} \\ \gamma = \pm 1 \\ D \end{matrix}$$

9

- Solution  $h = h_0 + Q_5 \exp(-z - |z|)$
  - $z < 0$  ( $0 < r < 1$ )  $h = h_0 + Q_5$  flat
  - $z > 0$  ( $r > 1$ ) previous solution
  - antisym. tens.  $H = 2Q_5 \Theta(z) \times \text{volume form of } S^3$
  - dilaton → string coupling



- for  $h_0 = 0$   $\rightarrow$  the above solution "interpolates" between two exact CFT's (NS5) in two regions of the target space

*flat*

$S^3$  + linear dilaton

inside

outside

$\rightarrow$  transverse  
radial direction

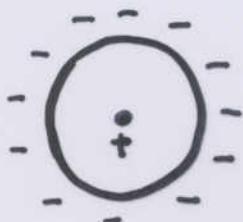
location of the thin shell of bones

- it solves the strong-coupling problem for NS branes; not for D.

- this pattern is quite generic:

## 6. Extensions and generalizations

- another electrostatic analogue: the screening



$$V = Q/r \quad \text{inside}$$

$$V = Q/r_0 \quad \text{outside}$$

- This is achieved in our framework by using orientifold planes ( $\gamma = -1$ ) or orbifold planes ( $\gamma = 1$ )

D5

NS5

Both are negative-tension objects allowing for new type II configurations

- we can obtain flat space outside the  $S^3$ 
  - ↳ solve the strong-coupling problem for D5
- we can have various exotic configuration of brane-shells with alternating flat and curved domains or D5 & NS5 ( $\rightarrow$  dyons)

- New kind of "CFT":
  - Target space made of several patches
  - Each patch is a piece of the target space of a 2-d Conformal  $\sigma$ -model
  - needs further investigation
  - low-energy spectrum by studying the fluctuations in the off. field  $H$ 
    - ↳ gravitational K-K sector
    - ↳ non-localized states with a mass gap
    - ↳ monopole states and moduli

## 7. Summary

standard 5-brane configuration



decoupling limit & LST      target space of  $\sigma$ -model (NS: exact CFT)

"strong-coupling region"

- \* embed the original configuration in a "web of T-dualities"
- \* regularize by adding F- or D-strings:  
marginal deformations of  $AdS_3 \times S^3 \times T^4$
- \* distribute the branes over  $S^3$
  
- further investigation is needed in all cases

- \* incorporate the full geometry including the asymptotic region
- \* control of T-dualities in curved space
- \* recast the decoupling limit for spread sources
- \* analyse the "patchwork CFT's": spectra, ...