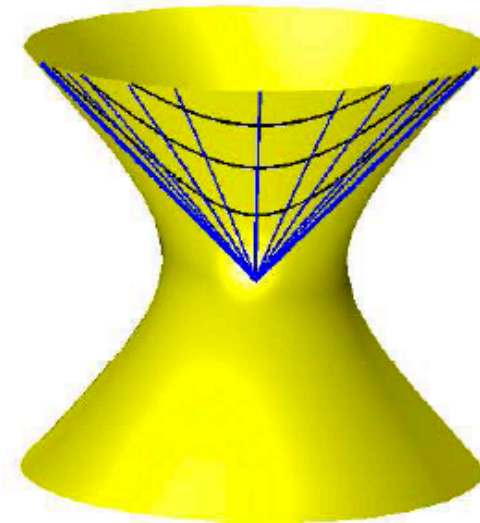
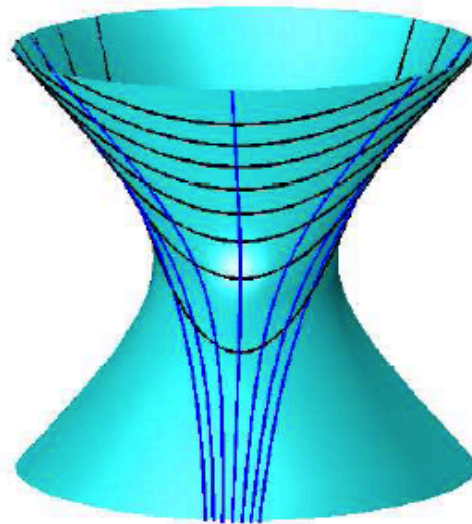
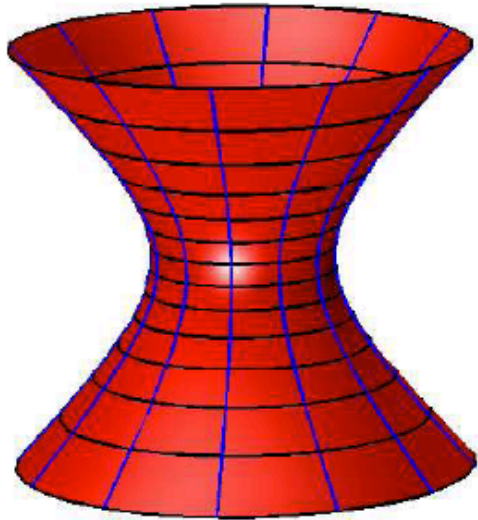


# Is our universe in a De Sitter phase?



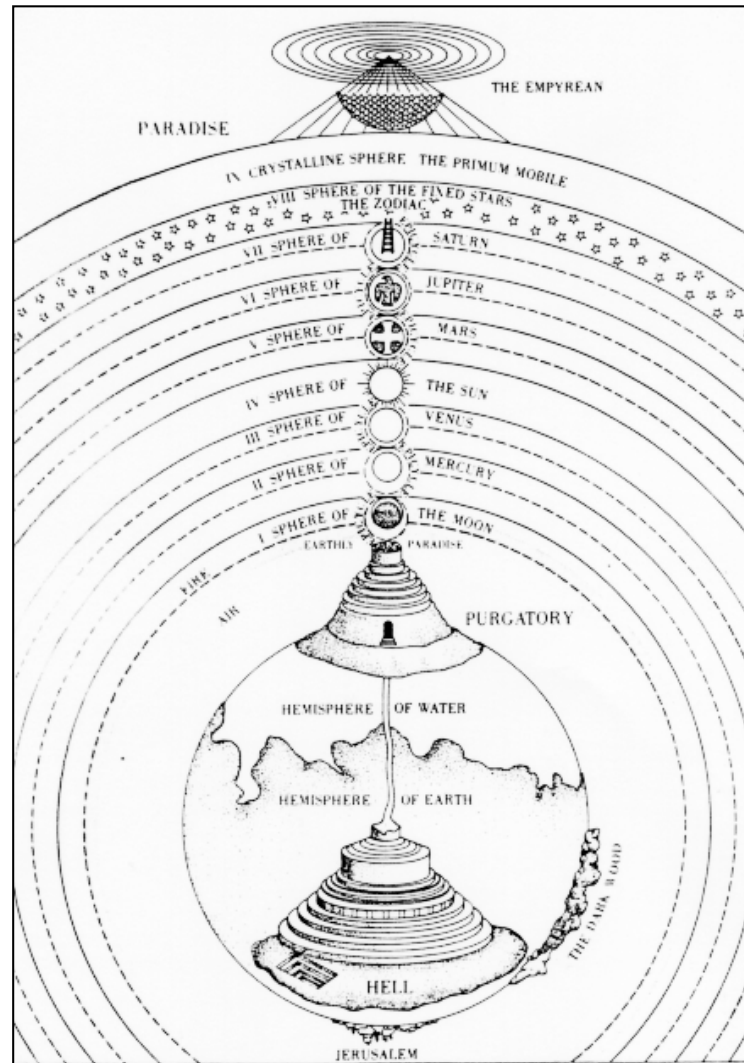
<http://www.bourbaphy.fr/moschella.pdf>

**Subir Sarkar**

*Rudolf Peierls Centre for Theoretical Physics, University of Oxford*

“The Physics of De Sitter Space-time”, Albert Einstein Institute, Hannover, 11-14 September 2012

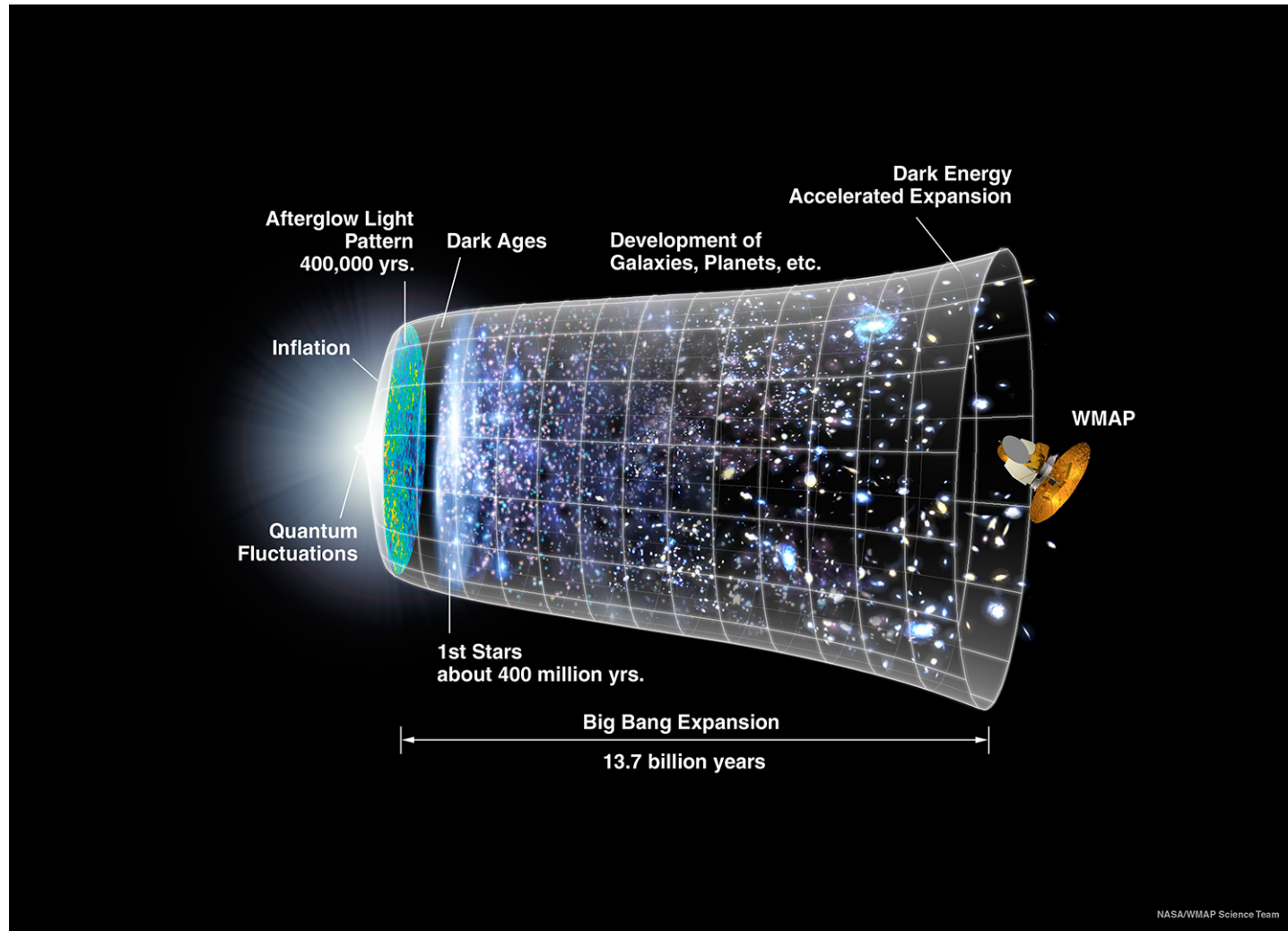
In the Aristotelean 'standard model' of cosmology (circa 350 BC) the universe was static and finite and centred on the Earth



*The Divine Comedy, Dante Alighieri (1321)*

This was a 'simple' model and fitted all the observational data ... but the underlying principle was unphysical

Today we have a new standard model of the universe ... dominated by dark energy and undergoing accelerated expansion



Courtesy: NASA/WMAP Science Team (2007)

It too is 'simple' and fits all the observational data ... but lacks an underlying physical basis

The Standard  $SU(3)_c \times SU(2)_L \times U(1)_Y$  Model provides an exact description of all microphysics (upto some cut-off  $M$ , when viewed as an effective field theory)

*Cosmological constant*

*Higgs mass divergence*

$$\mathcal{L}_{\text{eff}} = F^2 + \bar{\Psi} \not{D} \Psi + \bar{\Psi} \Psi \Phi + (D\Phi)^2 + \Phi^2 + \frac{\bar{\Psi} \Psi \Phi \Phi}{M} + \frac{\bar{\Psi} \Psi \bar{\Psi} \Psi}{M^2} + \dots$$

super-renormalisable  
renormalisable  
non-renormalisable

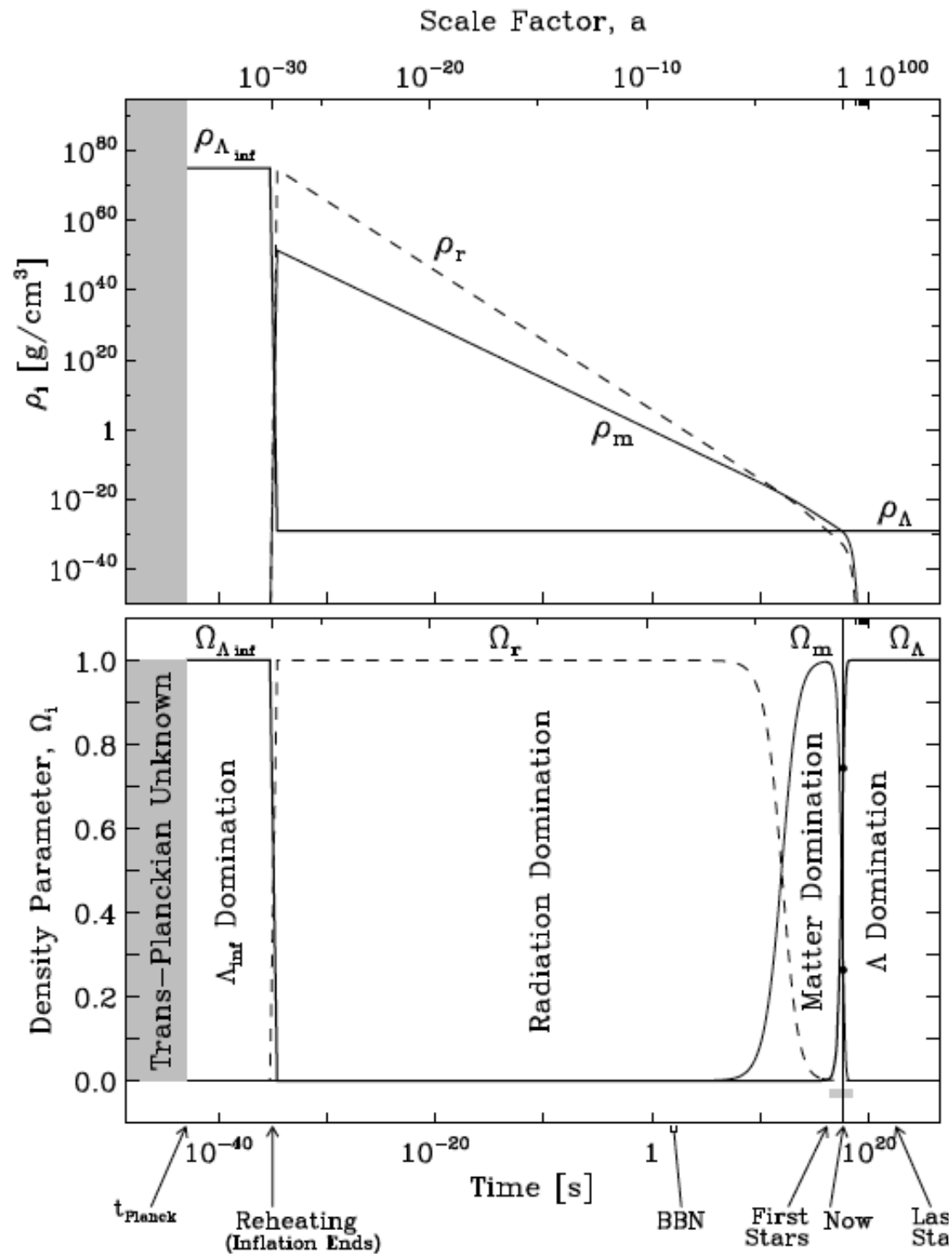
New physics beyond the SM (neutrino mass, nucleon decay, FCNC ...)  $\Rightarrow$  non-renormalisable operators suppressed by  $M^n$  ... which 'decouple' as  $M \rightarrow M_p$

But as  $M$  is raised, the effects of the super-renormalisable operators are exacerbated

There are possible solutions for the Higgs mass divergence, e.g. 'softly broken' supersymmetry at  $M \sim 1$  TeV (or perhaps the Higgs is composite e.g. as in technicolour)

But the 1<sup>st</sup> term couples to gravity so the natural expectation is  $\rho_\Lambda \sim (1 \text{ TeV})^4$  i.e. the universe should have been inflating since (or collapsed at)  $t \sim 10^{-12}$  s

Why did this not happen ... did  $\Lambda \rightarrow 0$  or does vacuum energy not couple to gravity?



Egan, arXiv:1005.0745

Either way this also raises the question of whether primordial inflation was really driven by the vacuum energy of a scalar field (fine-tuned at its potential minimum)

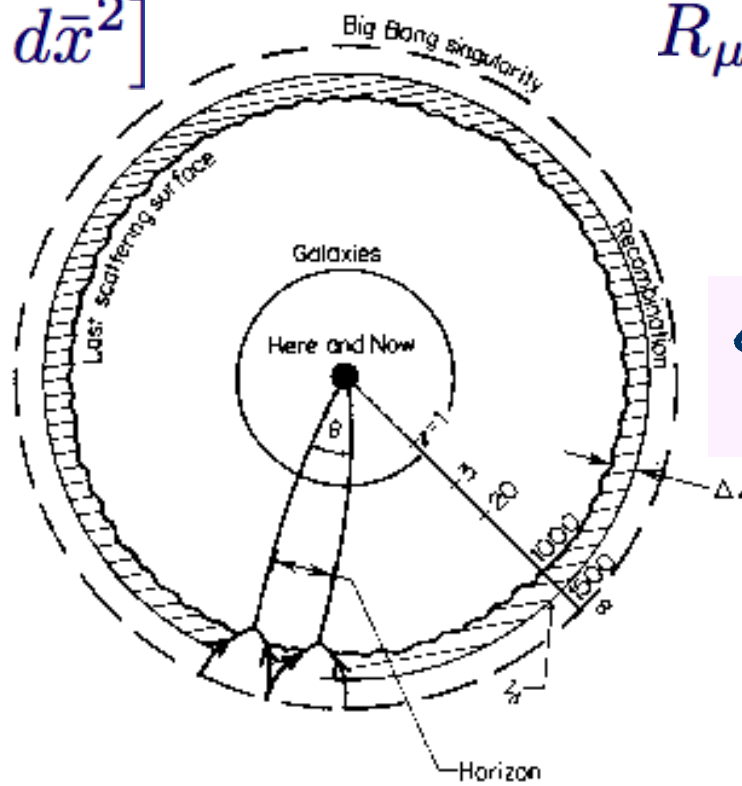
We are fascinated by De Sitter space-time ... but have no real understanding of the physical mechanism which can generate it ... and get us out of it!

The standard cosmological model is based on several key assumptions:  
*maximally symmetric space-time + general relativity + ideal fluids*

$$ds^2 = a^2(\eta) [d\eta^2 - d\bar{x}^2]$$

$$a^2(\eta)d\eta^2 \equiv dt^2$$

Space-time metric  
 Robertson-Walker



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

Geometrodynamics  
 Einstein

$$\Rightarrow H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G_N \rho_m}{3} - \frac{k}{a^2} + \frac{\Lambda}{3}$$

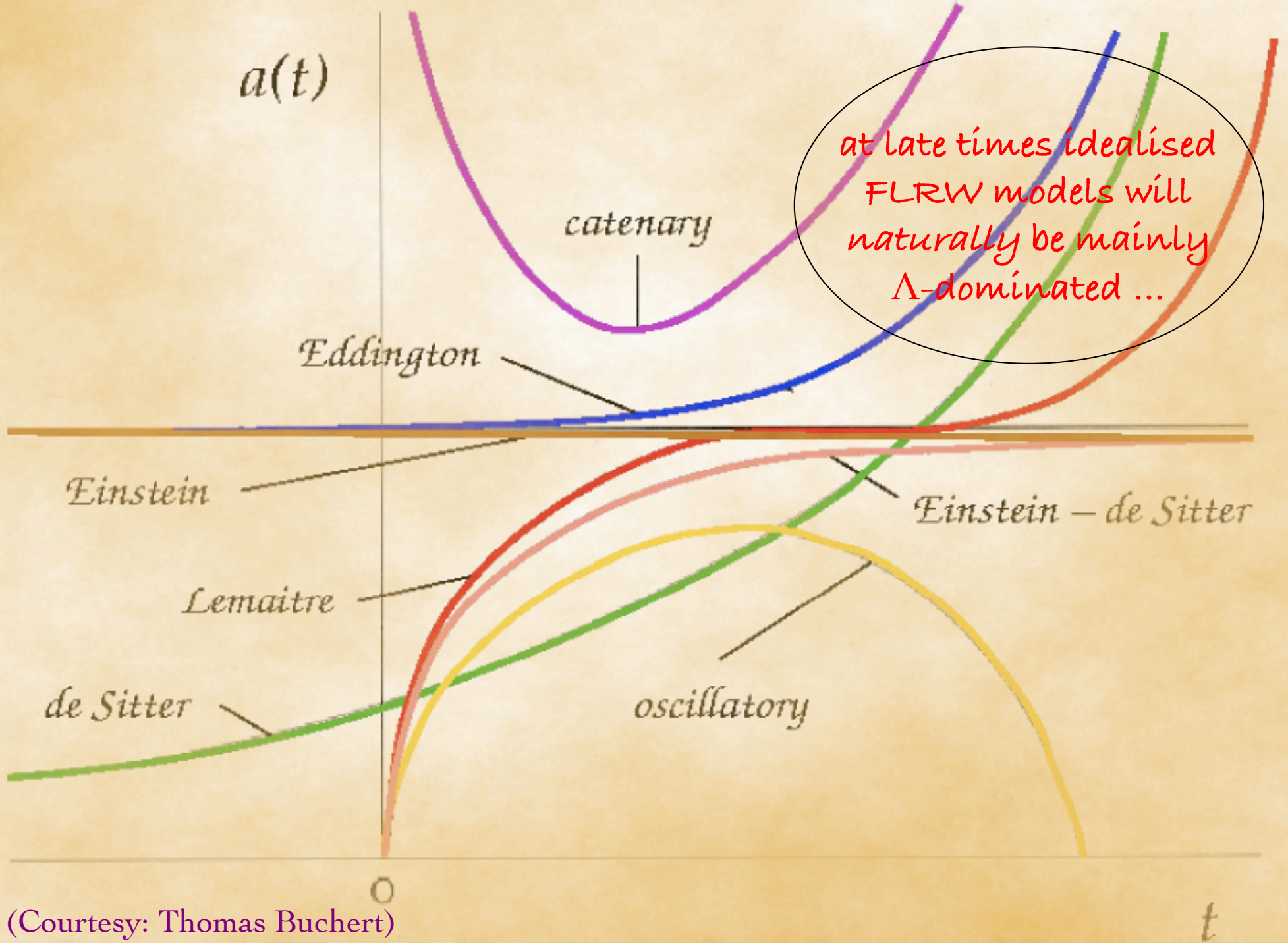
$$\equiv H_0^2 [\Omega_m(1+z)^3 + \Omega_k(1+z)^2 + \Omega_\Lambda]$$

$$\Omega_m \equiv \rho_{m0} / \frac{3H_0^2}{8\pi G_N}$$

$$\Omega_k \equiv -k/a_0^2 H_0^2$$

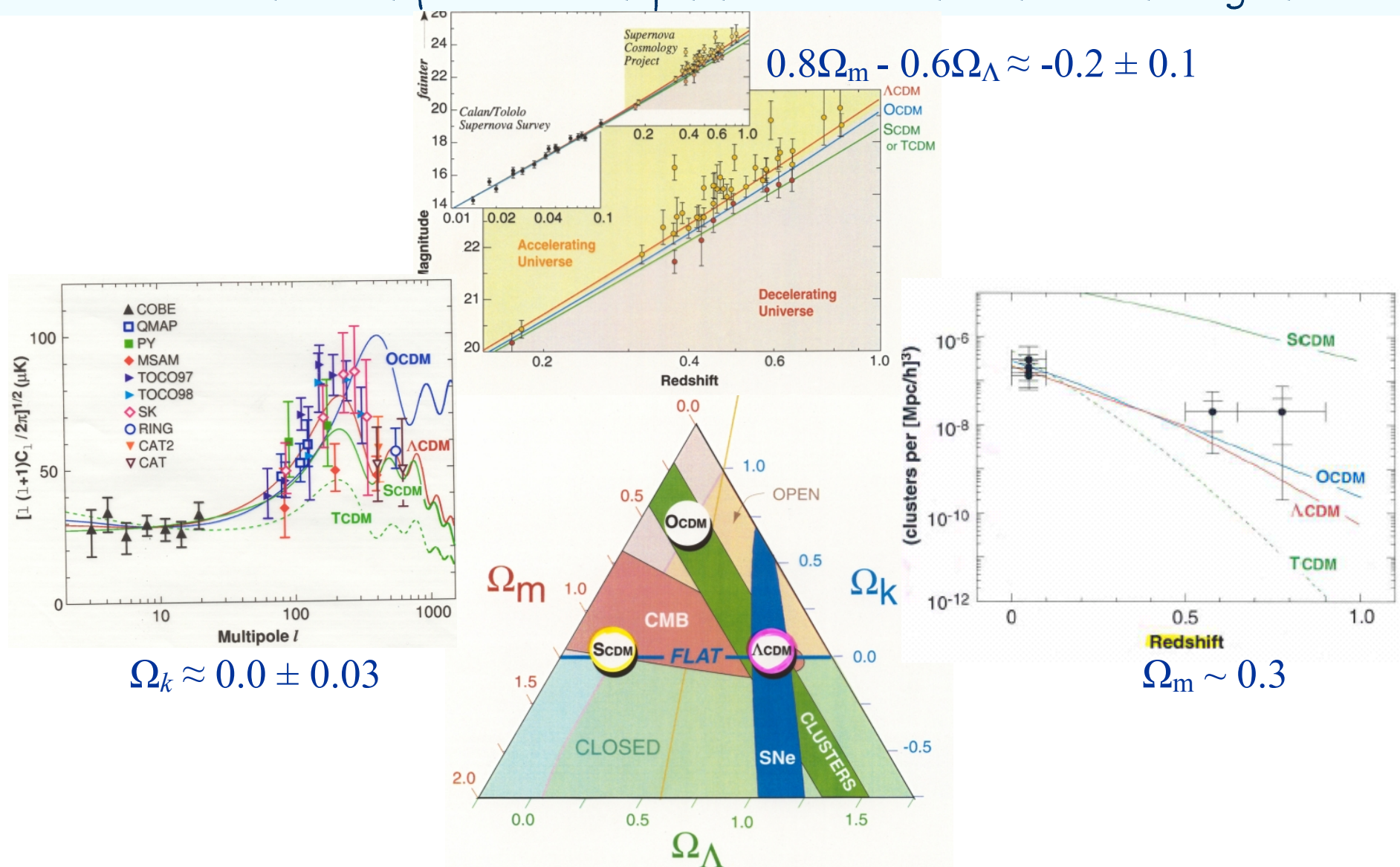
$$\Omega_\Lambda \equiv \Lambda/3H_0^2$$

... so has apparent late-time acceleration built into it!



(Courtesy: Thomas Buchert)

It is almost inevitable for data interpreted in this idealised model to give  $\Lambda \sim H_0^2$

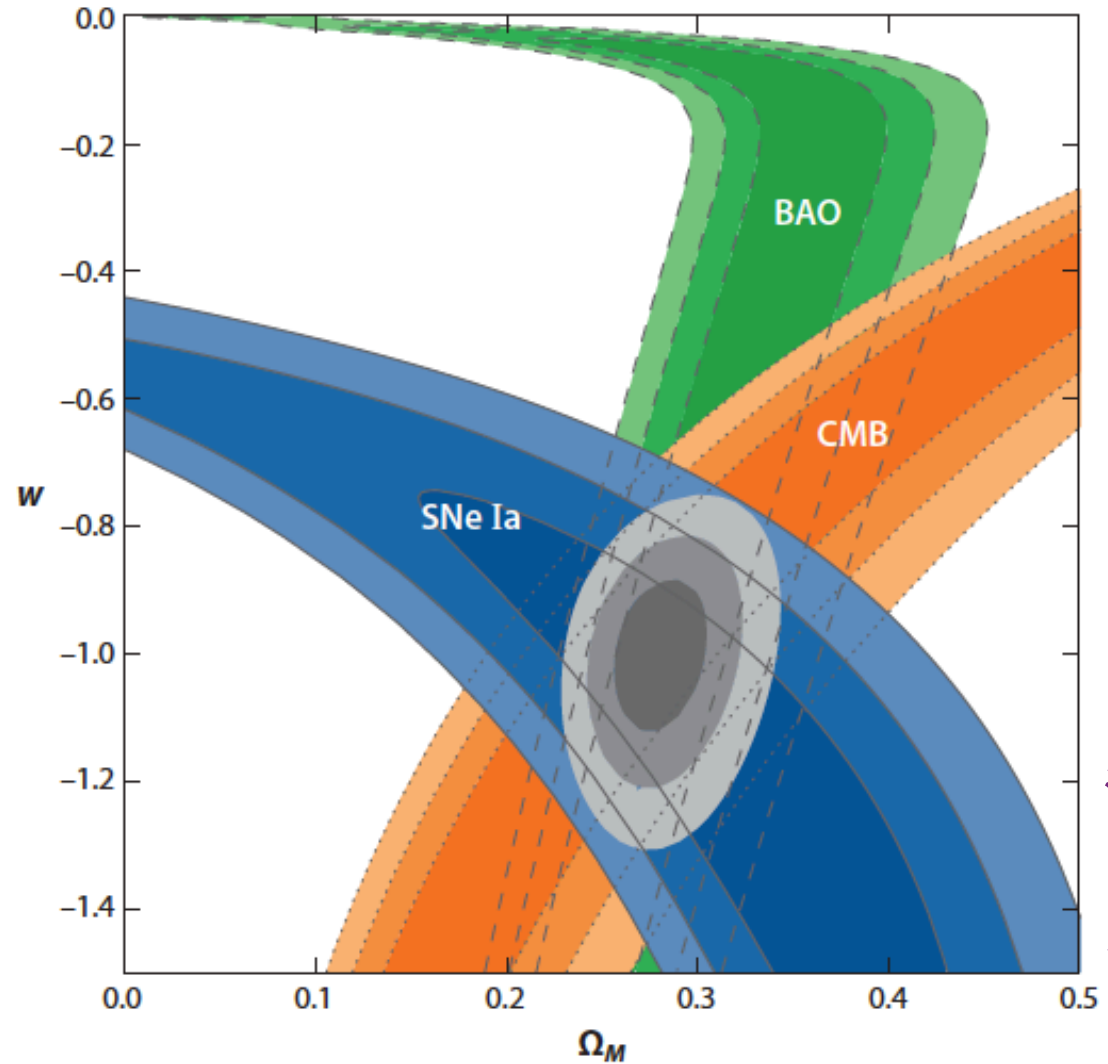
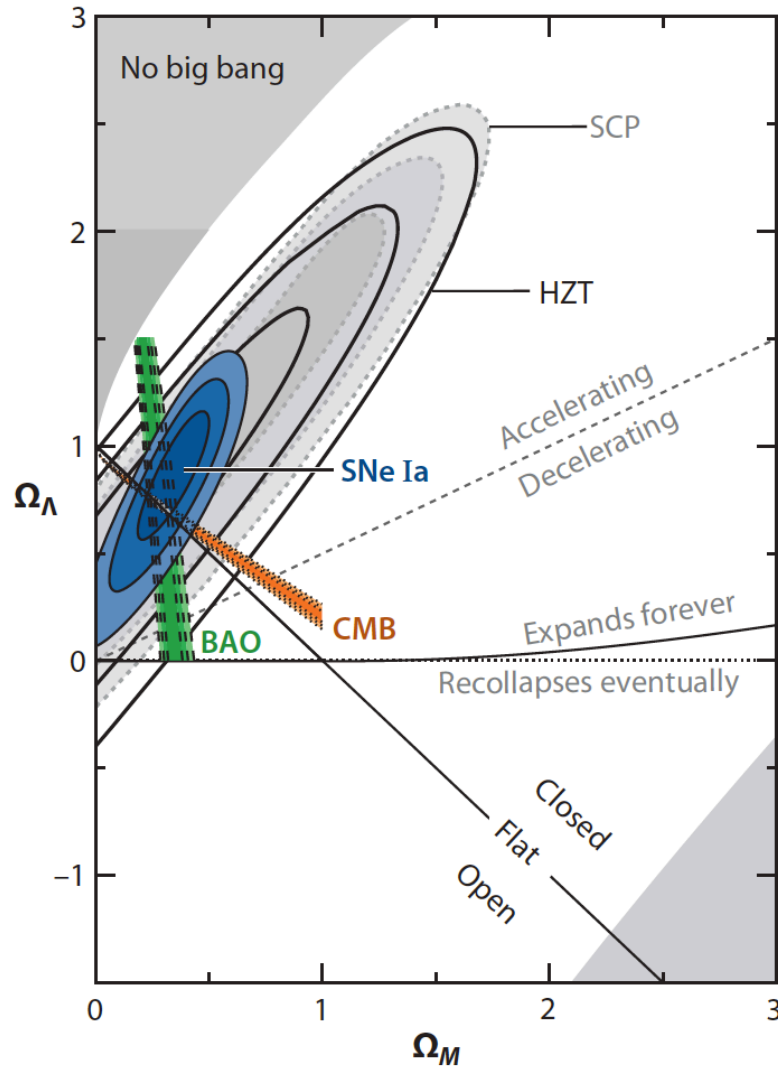


Bahcall *et al*, Science 284:1481, 1999

... Is it then so surprising that we infer  $\Omega_\Lambda (\equiv \Lambda/3H_0^2)$  to be of  $O(1)$  from the cosmic sum rule, given the uncertainties in measuring  $\Omega_m$  and  $\Omega_k$ , and the possibility of other components ( $\Omega_x$ ) which are not accounted for?



The 'concordance' of evidence for a Cosmological Constant today with  $\Lambda = 2H_0^2 \Rightarrow \Omega_\Lambda \approx 0.7$  is based entirely on geometrical measures (luminosity distance, angular diameter distance) which are interpreted assuming exact homogeneity



Goobar, Leibundgut, ARNPS 61:251,2011

The data have been interpreted more generally as implying 'dark energy' with negative pressure ( $w = p/\rho \approx -1$ ) but there is no direct evidence yet (late ISW effect) for this property

Quantities averaged over a domain  $\mathcal{D}$  obey modified Friedmann equations  
Buchert 1999:

$$3 \frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = -4\pi G \langle \rho \rangle_{\mathcal{D}} + Q_{\mathcal{D}} ,$$
$$3 \left( \frac{\dot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} \right)^2 = 8\pi G \langle \rho \rangle_{\mathcal{D}} - \frac{1}{2} \langle {}^{(3)}R \rangle_{\mathcal{D}} - \frac{1}{2} Q_{\mathcal{D}} ,$$

where  $Q_{\mathcal{D}}$  is the backreaction term,

$$Q_{\mathcal{D}} = \frac{2}{3} (\langle \theta^2 \rangle_{\mathcal{D}} - \langle \theta \rangle_{\mathcal{D}}^2) - \langle \sigma^{\mu\nu} \sigma_{\mu\nu} \rangle_{\mathcal{D}} .$$

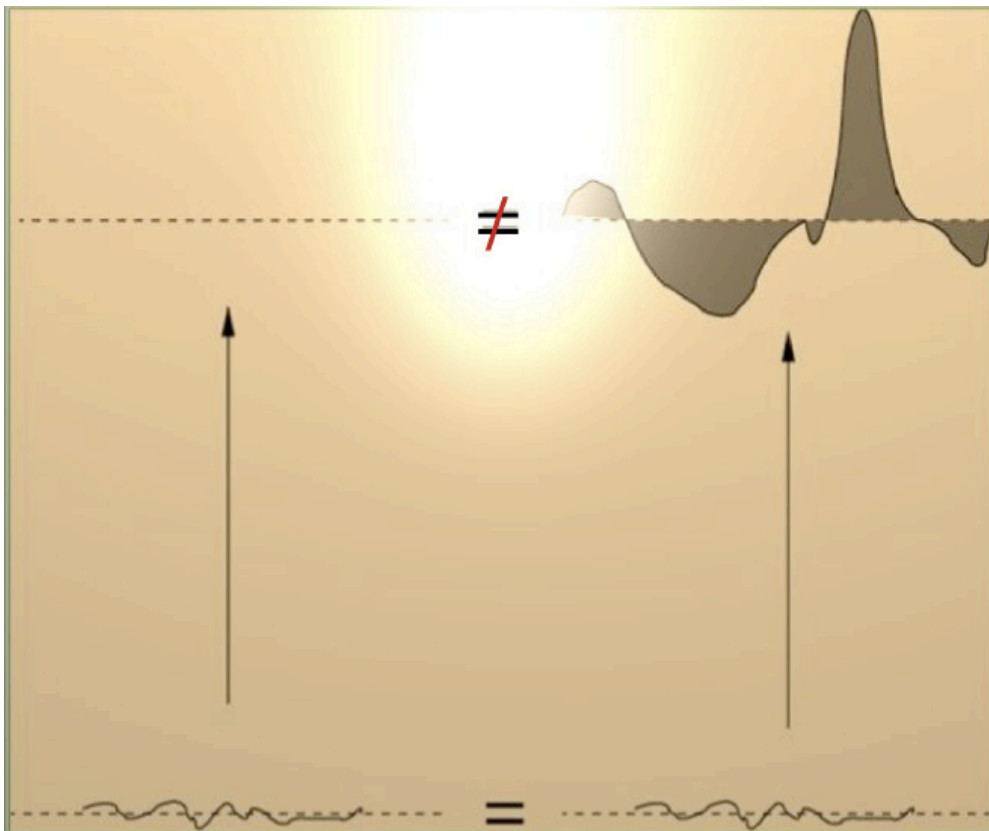
Variance of the expansion rate.

Average shear.

If  $Q_{\mathcal{D}} > 4\pi G \langle \rho \rangle_{\mathcal{D}}$  then  $a_{\mathcal{D}}$  accelerates.

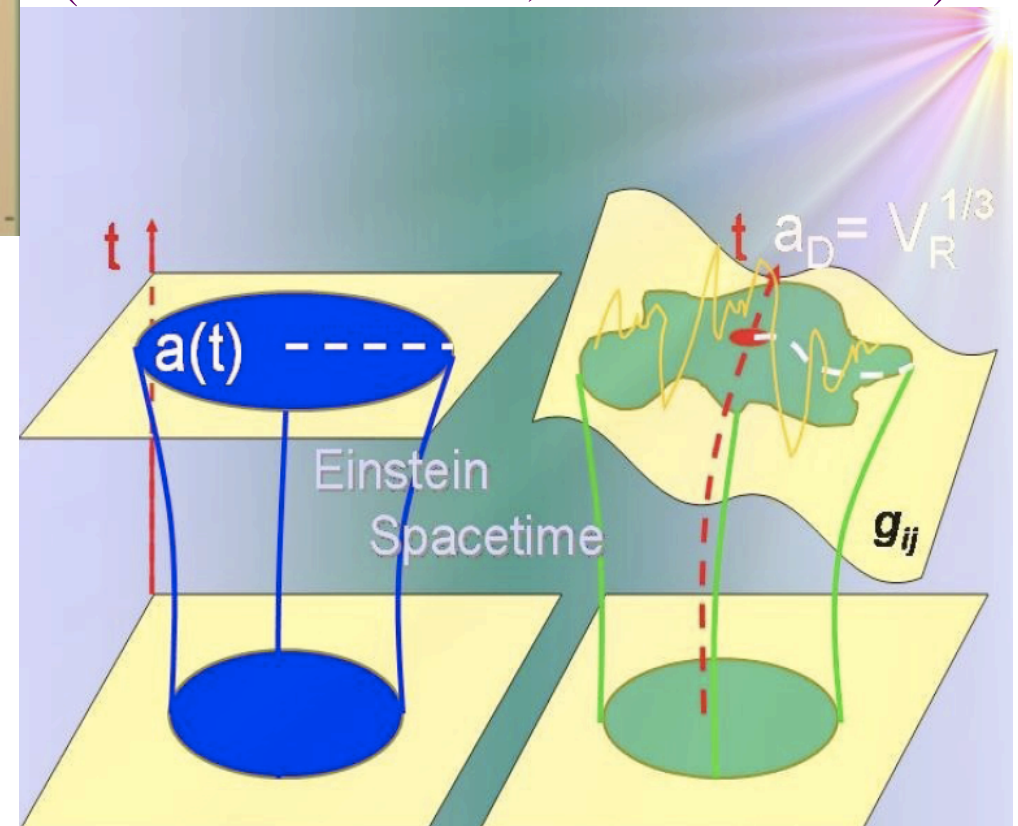
Can mimic a cosmological constant if  $Q_{\mathcal{D}} = -\frac{1}{3} \langle {}^{(3)}R \rangle_{\mathcal{D}} = \Lambda_{\text{eff}}$ .

Whether the backreaction can be sufficiently large is an open question



'Back reaction' is hard to compute because spatial averaging and time evolution (along our past light cone) do not commute (Ellis 1982) ... hence the ongoing controversy (Buchert & Rasanen, arXiv:1112.5335)

Due to structure formation, the homogeneous solution of Einstein's eqs. is distorted - its average must be taken over the actual geometry ... the result is different from the standard FLRW model



Courtesy: Thomas Buchert

Interpreting  $\Lambda$  as vacuum energy raises the coincidence problem:

*why is  $\Omega_\Lambda \approx \Omega_m$  today?*

Option 1: invent an ultralight scalar field ('quintessence') with  $V(\phi)^{1/4} \sim 10^{-12}$  GeV but  $\sqrt{d^2V/d\phi^2} \sim H_0 \sim 10^{-42}$  GeV, which displays 'tracking' behaviour

*... but this is just as much fine-tuning as a bare cosmological constant*

Option 2: Modify gravity on the scale of the present Hubble radius so as to mimic vacuum energy ('DGP brane-world') taking care to avoid instabilities

*... this scale is unnatural in a fundamental theory and is just put in by hand*

Option  $n \gg 1$ : chameleon/ $f(R)$  models, symmetron fields, massive gravity ...

*All 'explanations' for cosmic accn. insert the scale  $H_0 \sim 10^{-42}$  GeV by hand!*

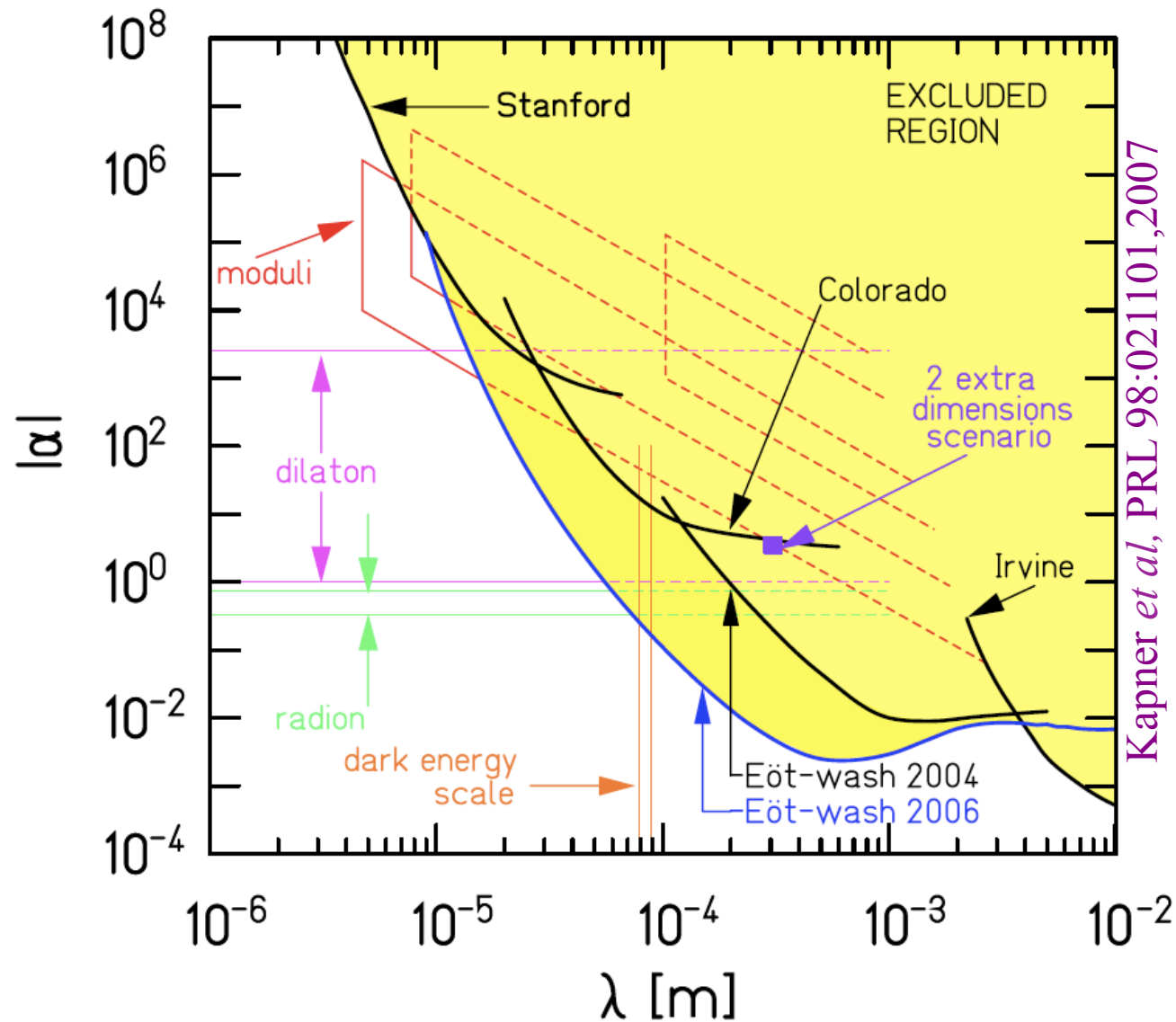
The only natural option is if  $\Lambda \sim H^2$  always (as attempted in e.g. causal set models), but this is just a renormalisation of  $G_N$  (recall:  $H^2 = 8\pi G_N/3 + \Lambda/3$ ) (ruled out by Big Bang nucleosynthesis which requires  $G_N$  to be within few% of its laboratory value ... and in any case this will not yield accelerated expansion)

*There cannot be a natural explanation for the coincidence problem*

Do we infer  $\Lambda \sim H_0^2$  simply because  $H_0$  is the only scale in the FLRW model and enters in every observation (through the distance scale)?

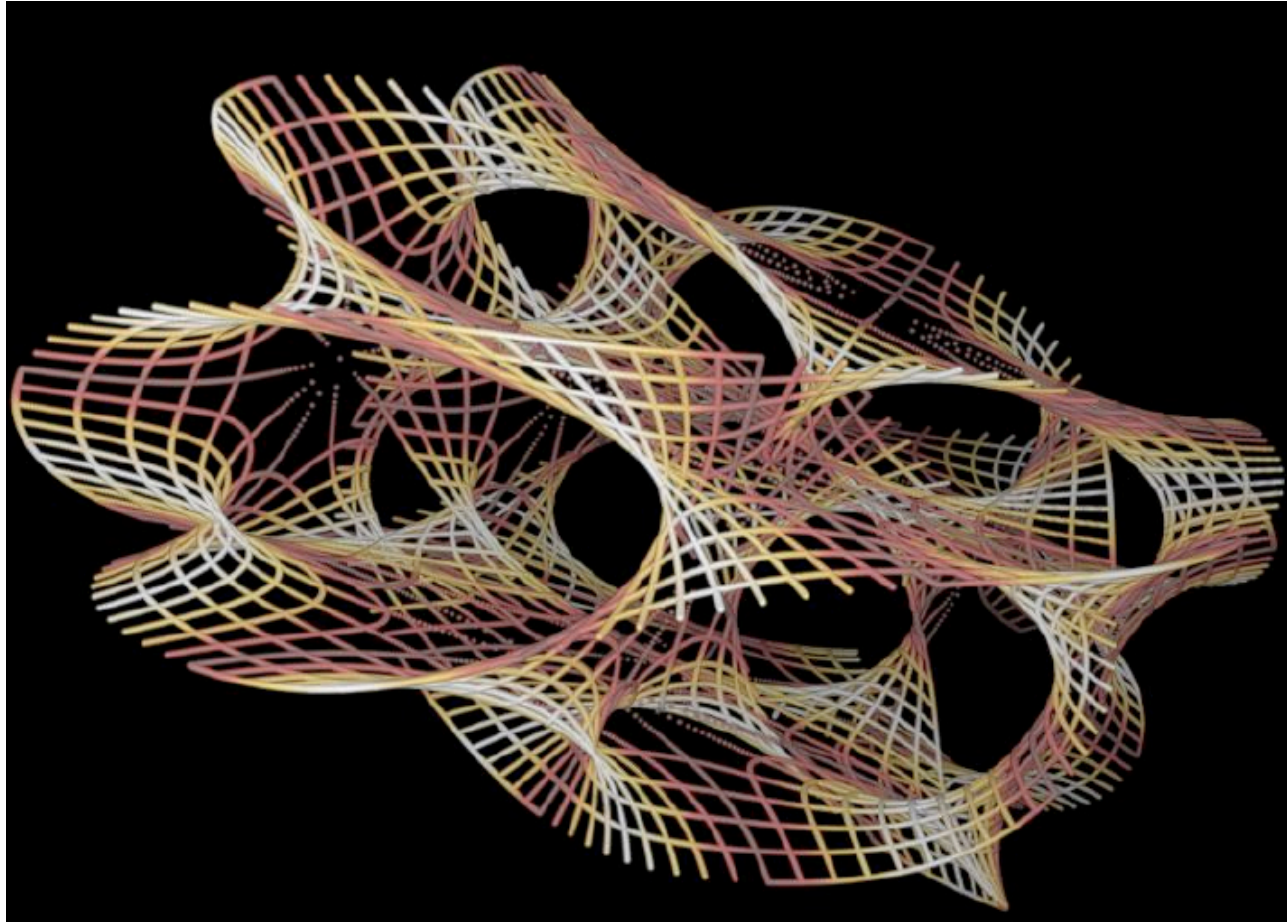
There is no evidence for any change in the inverse-square law of gravitation at the 'dark energy' scale:  $\Omega_\Lambda^{-1/4} \sim (H_0 M_P)^{-1/2} \sim 0.1 \text{ mm}$

$$V(r) = -G \frac{m_1 m_2}{r} [1 + \alpha \exp(-r/\lambda)]$$



Kapner et al, PRL 98:021101, 2007

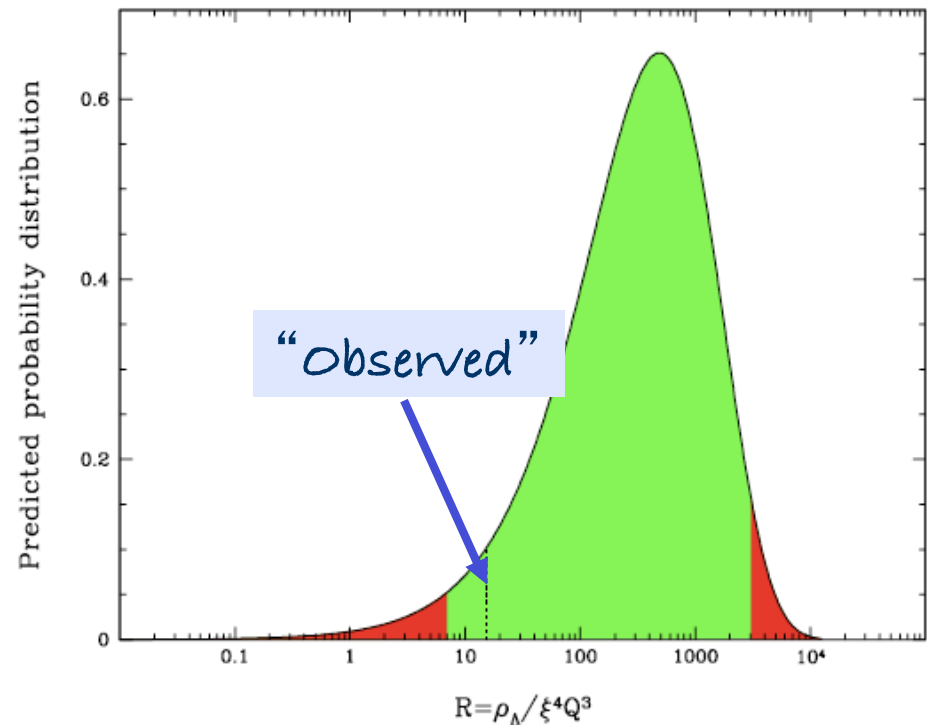
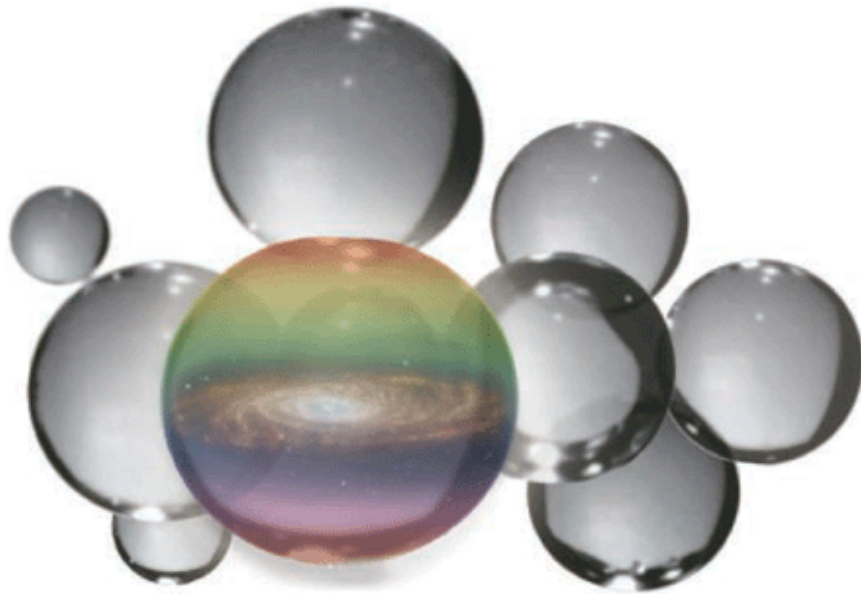
In string/M-theory, the sizes and shapes of the extra dimensions ('moduli') must be stabilised ... e.g. by turning on background 'fluxes'



Given the variety of flux choices and the number of local minima in the flux potential, the total number of vacua is very large - perhaps  $10^{500}$

The existence of the huge landscape of possible vacua in string theory (with moduli stabilised through background fluxes) has remotivated attempts at an 'anthropic' explanation for  $\Omega_\Lambda \sim \Omega_m$

Perhaps it is just "observer bias" - galaxies would not have formed if  $\Lambda$  had been higher (Weinberg 1989, Efstathiou 1995, Martel, Shapiro, Weinberg 1998 ...)



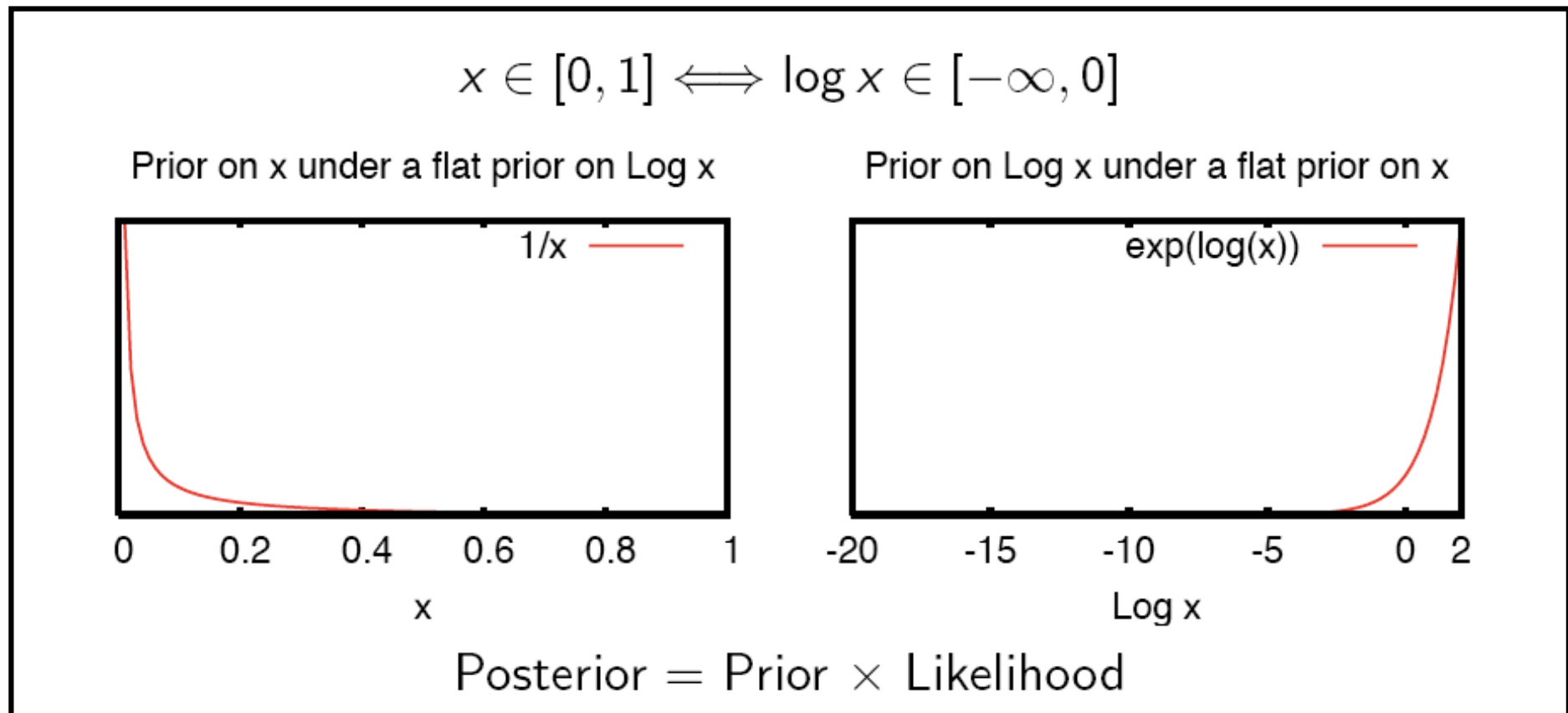
Tegmark, Aguirre, Rees, Wilczek,  
Phys.Rev. D73:023505, 2006

But the 'anthropic prediction' of  $\Lambda$  from considerations of galaxy formation is significantly higher than the observationally inferred value ... not surprising since galaxy formation occurred at redshift  $z \sim 3-5$  when the matter density was  $(1+z)^3$  times higher

Moreover this assumes the prior to be flat in the range  $0 \rightarrow 10^{-120} M_p^4$

Since we have no physical understanding of  $\Lambda$ , this may not be reasonable

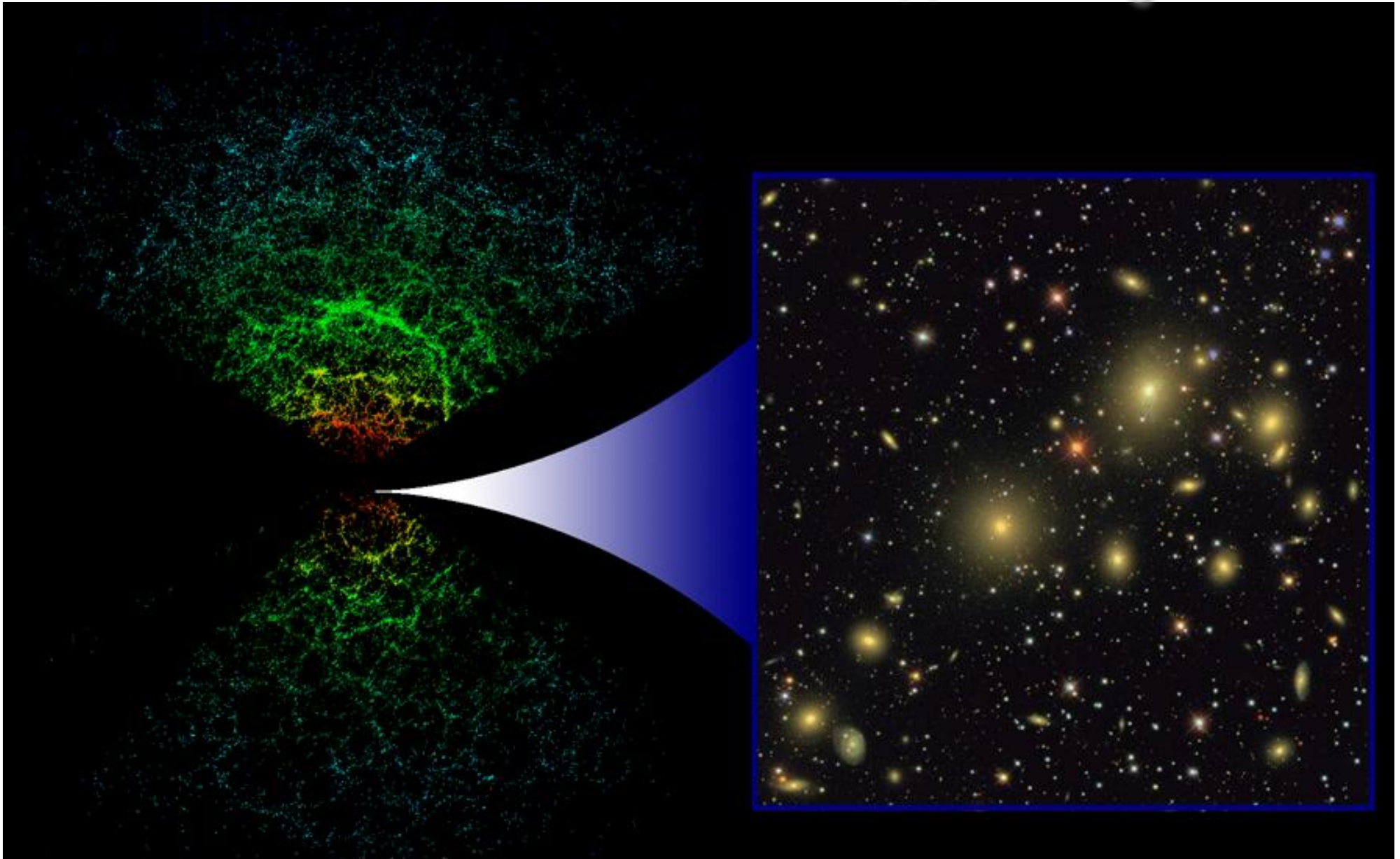
if the relevant physical variable is e.g.  $\log \Omega_\Lambda$ , then  $\Omega_\Lambda = 0$  would be favoured!



So it is far from clear that  $\Lambda \sim H_0^2$  has an anthropic explanation



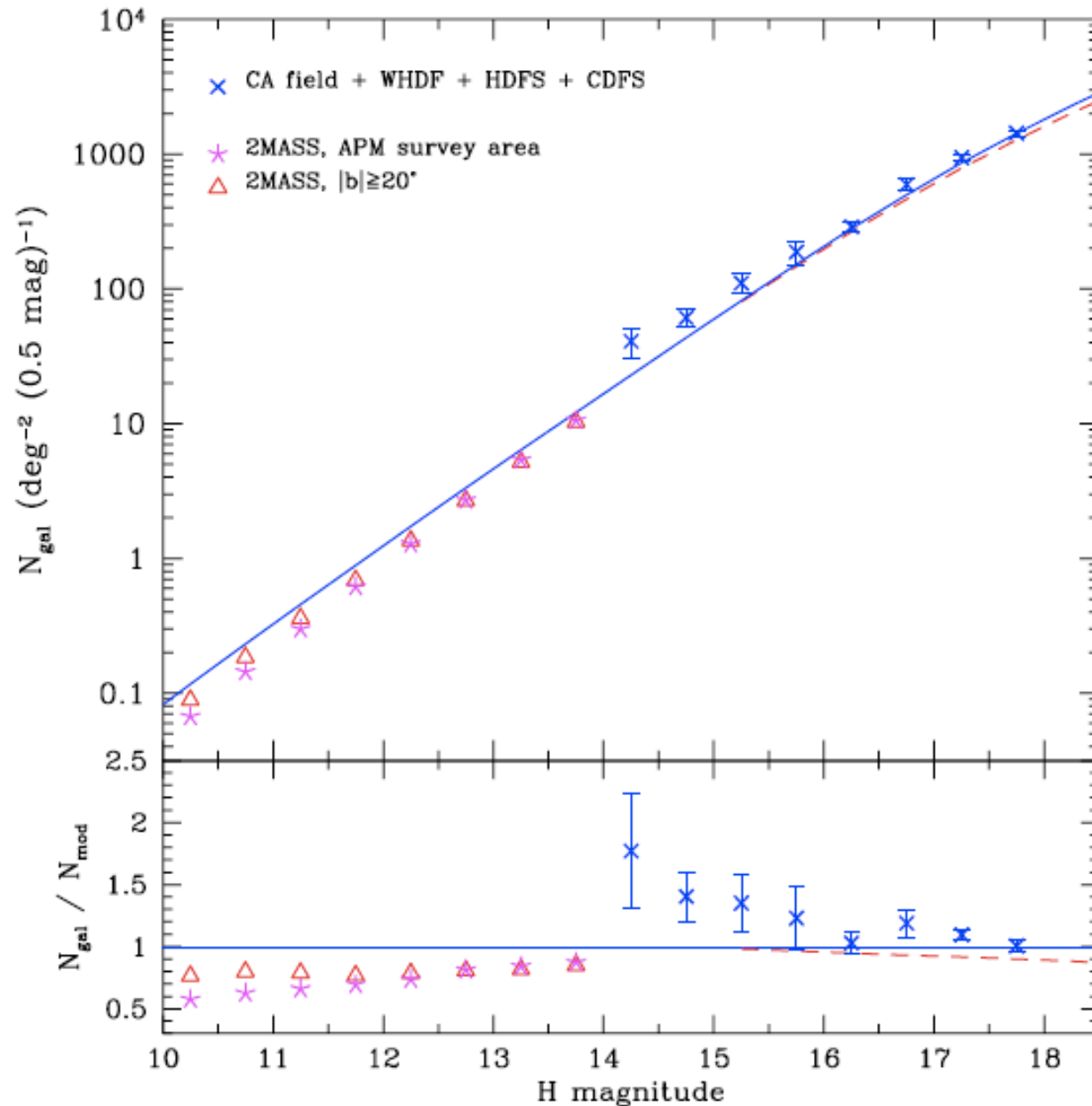
Galaxies are seen to trace out a cosmic 'web' of filamentary structure



Averaged on large scales the universe may be homogeneous but how would it bias cosmological inferences if we are located in a void?

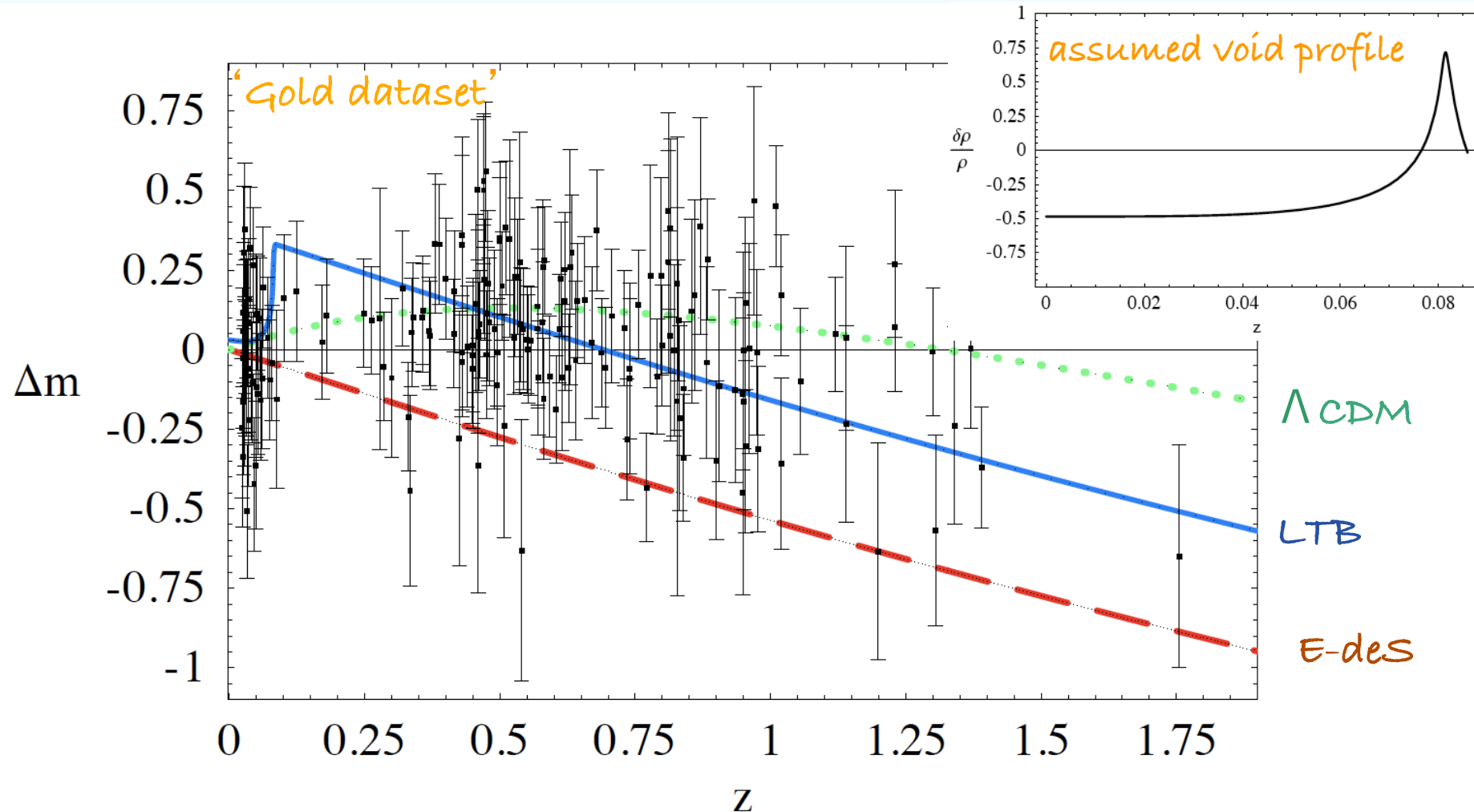
# New H-band Galaxy Number Counts

Are we located in an underdense region in the galaxy distribution?



Frith, Metcalfe & Shanks (2006)

If so, the SN Ia Hubble diagram can be explained without invoking acceleration, since distant supernovae would be in a slower Hubble flow than the nearby ones within the local void (inhomogeneous Lemaitre-Tolman-Bondi model)



Fits the SN data with  $h_{\text{out}} \sim 0.45$ ,  $0.51 < h_{\text{in}} < 0.59$ , void radius  $\sim 150\text{-}250 \text{ Mpc } h_{\text{in}}^{-1}$

However subsequent SN data has filled in the gap at  $z \sim 0.1\text{-}0.4$  and may have ruled out this model ... so now one needs to consider a larger void of  $\sim 4\text{ Gpc}$  size

Toy model that fits the SN Ia Hubble diagram without cosmic acceleration/dark energy

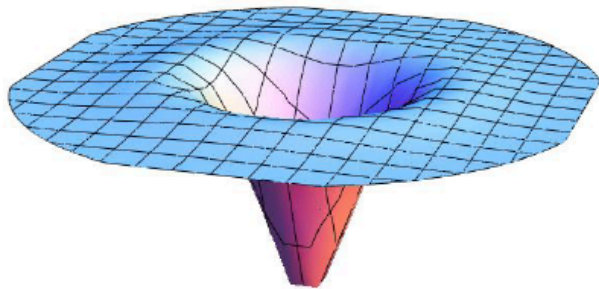
LTB metric:  $ds^2 = -c^2 dt^2 + \frac{A'^2(r,t)}{1+K(r)} dr^2 + A^2(r,t) d\Omega^2$

Two Hubble rates:  $H_T(r,t) \equiv \frac{\dot{A}}{A}$  and  $H_L(r,t) \equiv \frac{\dot{A}'}{A'}$

Obtain modified version of Friedmann equation -

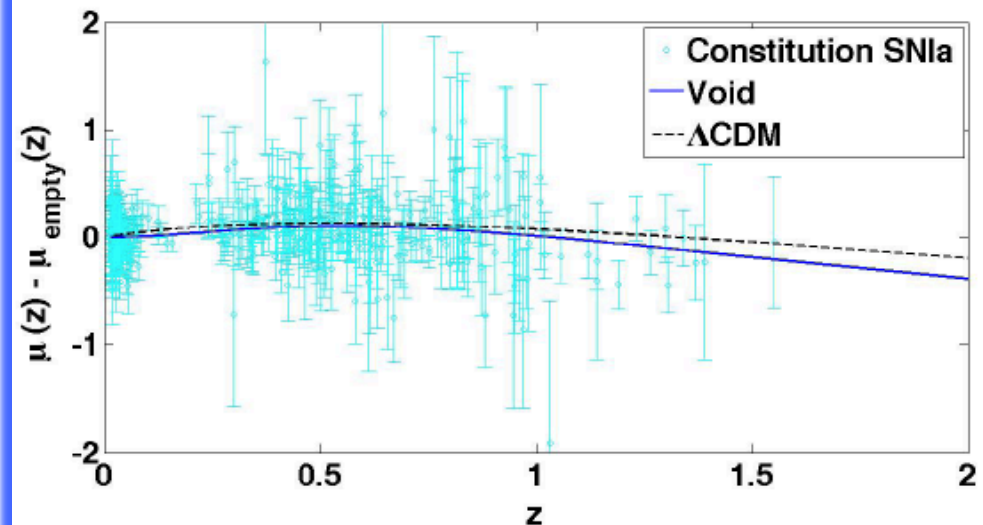
$$H_T^2(r,t) = H_0^2(r) \left[ \Omega_m(r) \left( \frac{A_0(r)}{A(r,t)} \right)^3 + \Omega_K(r) \left( \frac{A_0(r)}{A(r,t)} \right)^2 \right]$$

Choose a density profile

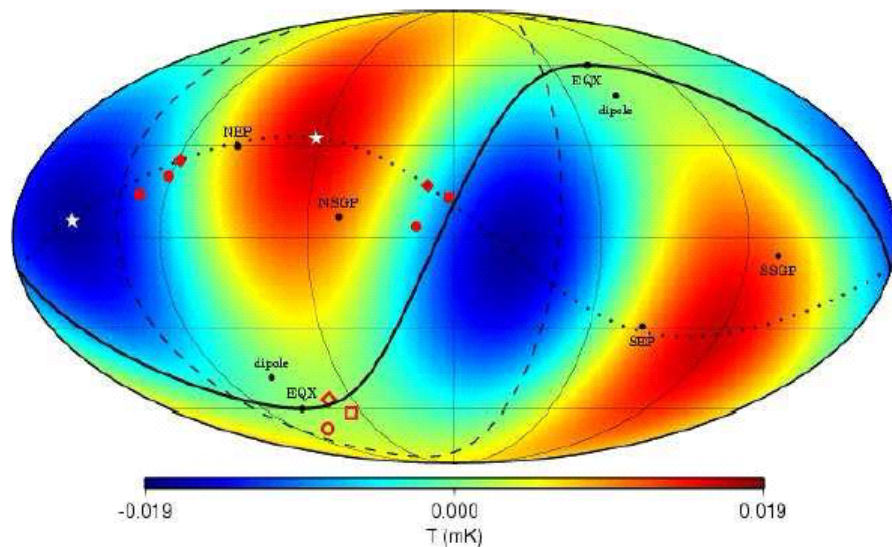


- Solve modified Friedmann equation numerically for  $A(r,t)$
- Luminosity distance  
 $D_L(z) = (1+z)^2 A(r,t)$

Can get a very good fit to supernovae data:



*e.g. Nadathur & Sarkar, Phys. Rev. D83:063506,2011*

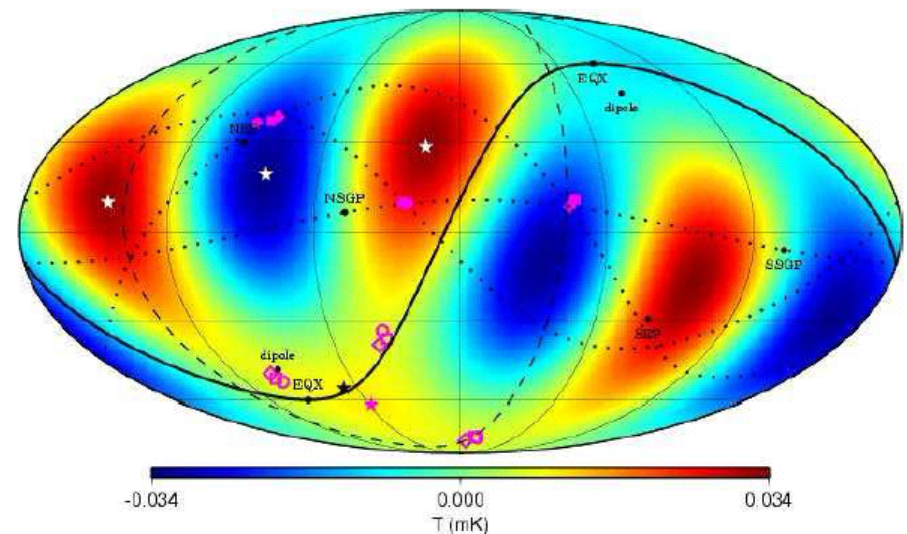


The local void need not be exactly spherical ... nor would we expect to be exactly at its centre

So might expect (low  $l$ ) CMB anisotropies to be generated by the 'Rees-Sciama effect' (must be within  $\sim$  few % of the centre so as to not generate excessive dipole)

The CMB quadrupole and octupole are indeed very well-aligned!

This requires us to be located at the boundary between two voids (to yield the observed planar - rather than linear - alignment)  
Inoue & Silk, APJ 648:23, 2006

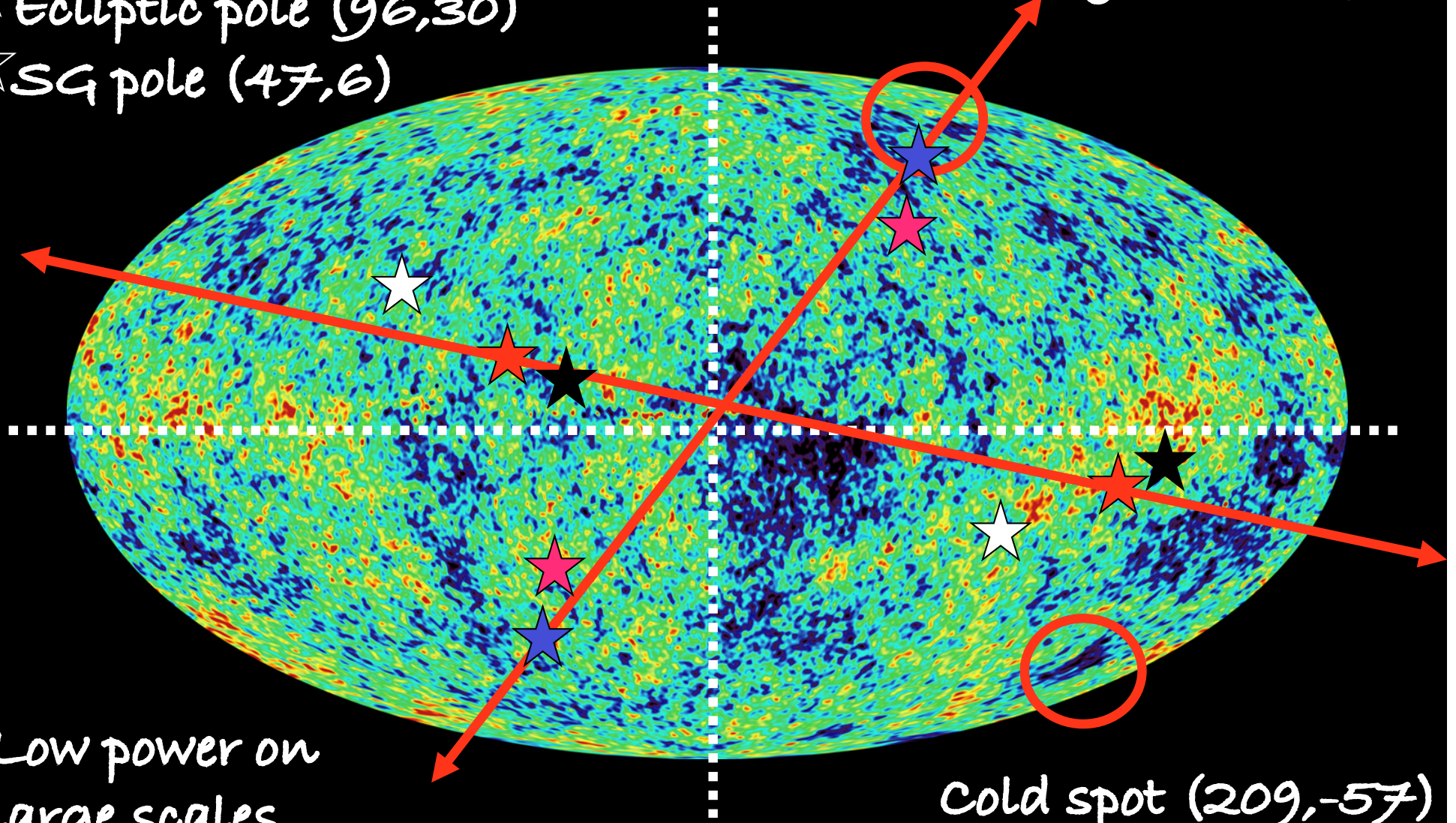


Can such a void be responsible for the CMB 'cold spot'?

★ Axis of Evil  $\sim (260, 60)$   
★ Dipole  $(264, 48)$

★ Max asym axis  $(57, 10)$   
★ Ecliptic pole  $(96, 30)$   
★ SG pole  $(47, 6)$

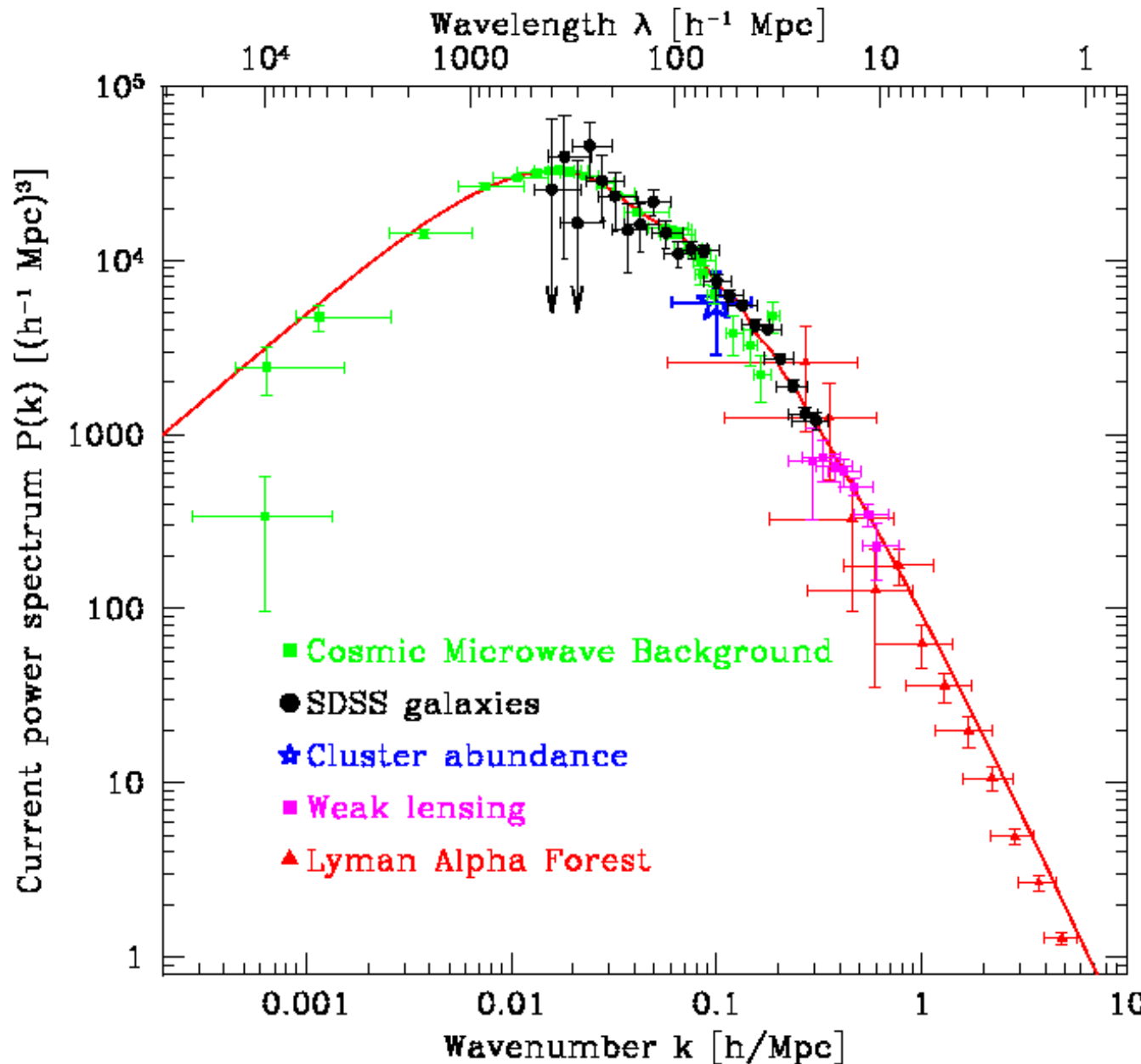
Virgo  $\sim (260, 70)$



Low power on large scales

Cold spot  $(209, -57)$

Observations of large-scale structure are consistent with the  $\Lambda$ CDM model if the primordial fluctuations are *adiabatic* and *~scale-invariant* (as “expected in the simplest models of inflation”)



Tegmark (2004)

The formation of large-scale structure is akin to a scattering experiment

**The Beam:** inflationary density perturbations

No 'standard model' - usually assumed to be *adiabatic* and *~scale-invariant*

**The Target:** dark matter (+ baryonic matter)

Identity unknown - usually taken to be *cold* (sub-dominant 'hot' component?)

**The Detector:** the universe

Modelled by a 'simple' *FRW cosmology* with parameters  $h, \Omega_{\text{CDM}}, \Omega_{\text{b}}, \Omega_{\Lambda}, \Omega_k \dots$

**The Signal:** CMB anisotropy, galaxy clustering ...

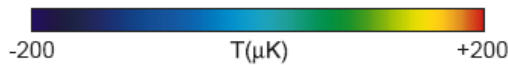
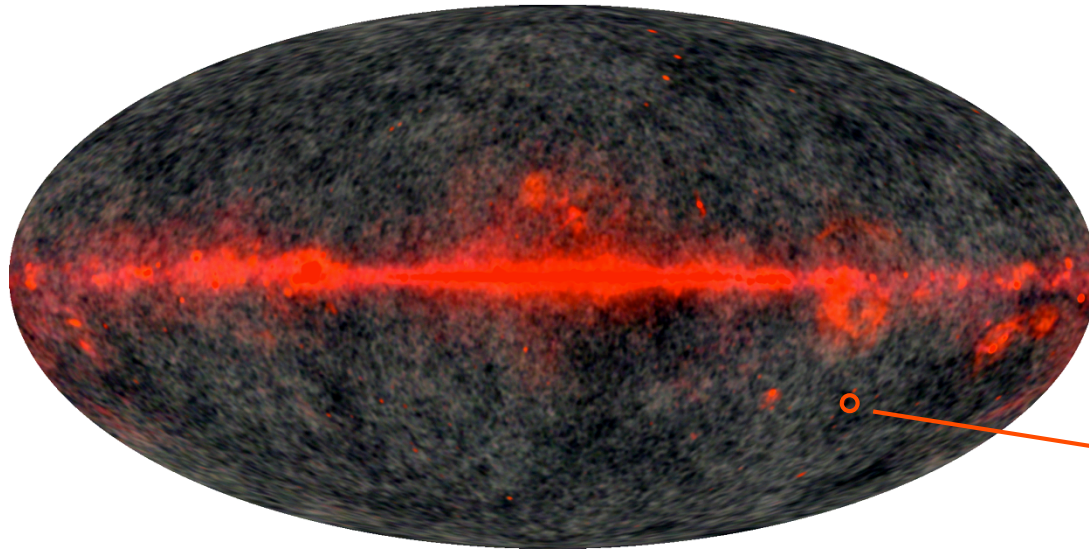
measured over scales from  $\sim 1 - 10000$  Mpc ( $\Rightarrow \sim 8$  e-folds of inflation)

We cannot simultaneously determine the properties of both the beam and the target with an unknown detector

... hence need to adopt suitable 'priors' on  $h, \Omega_{\text{CDM}}$ , etc in order to break inevitable parameter degeneracies



# 'Internal Linear Combination' map

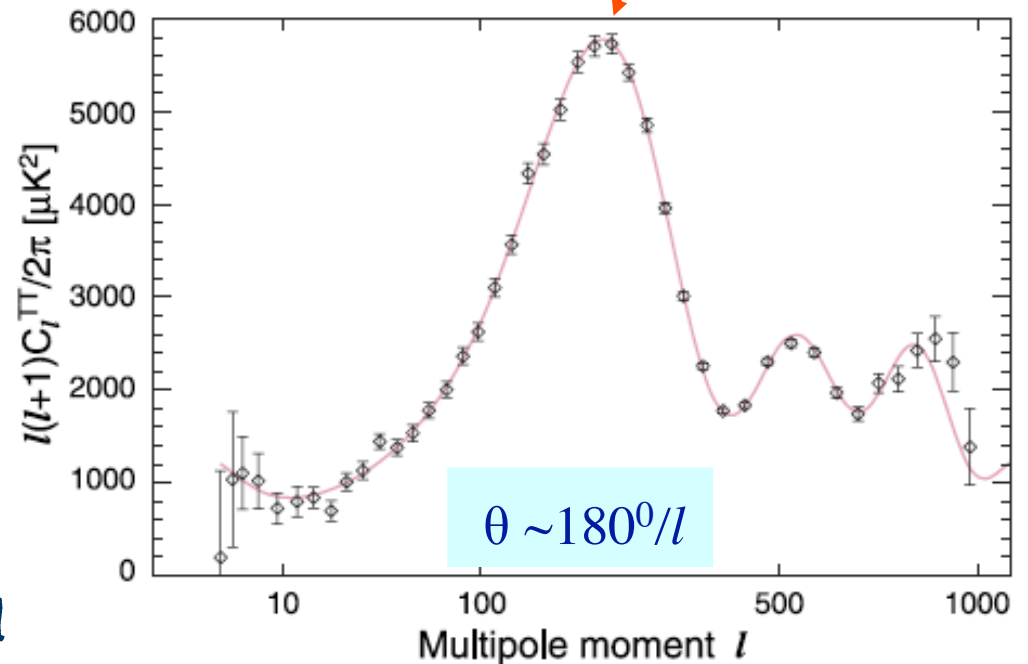


Coherent oscillations  
in photon-baryon  
plasma, excited by  
primordial density  
perturbations on  
super-horizon scales ...  
(Hubble radius at  $t_{rec}$ )

$$\Delta T(\mathbf{n}) = \sum a_{lm} Y_{lm}(\mathbf{n})$$

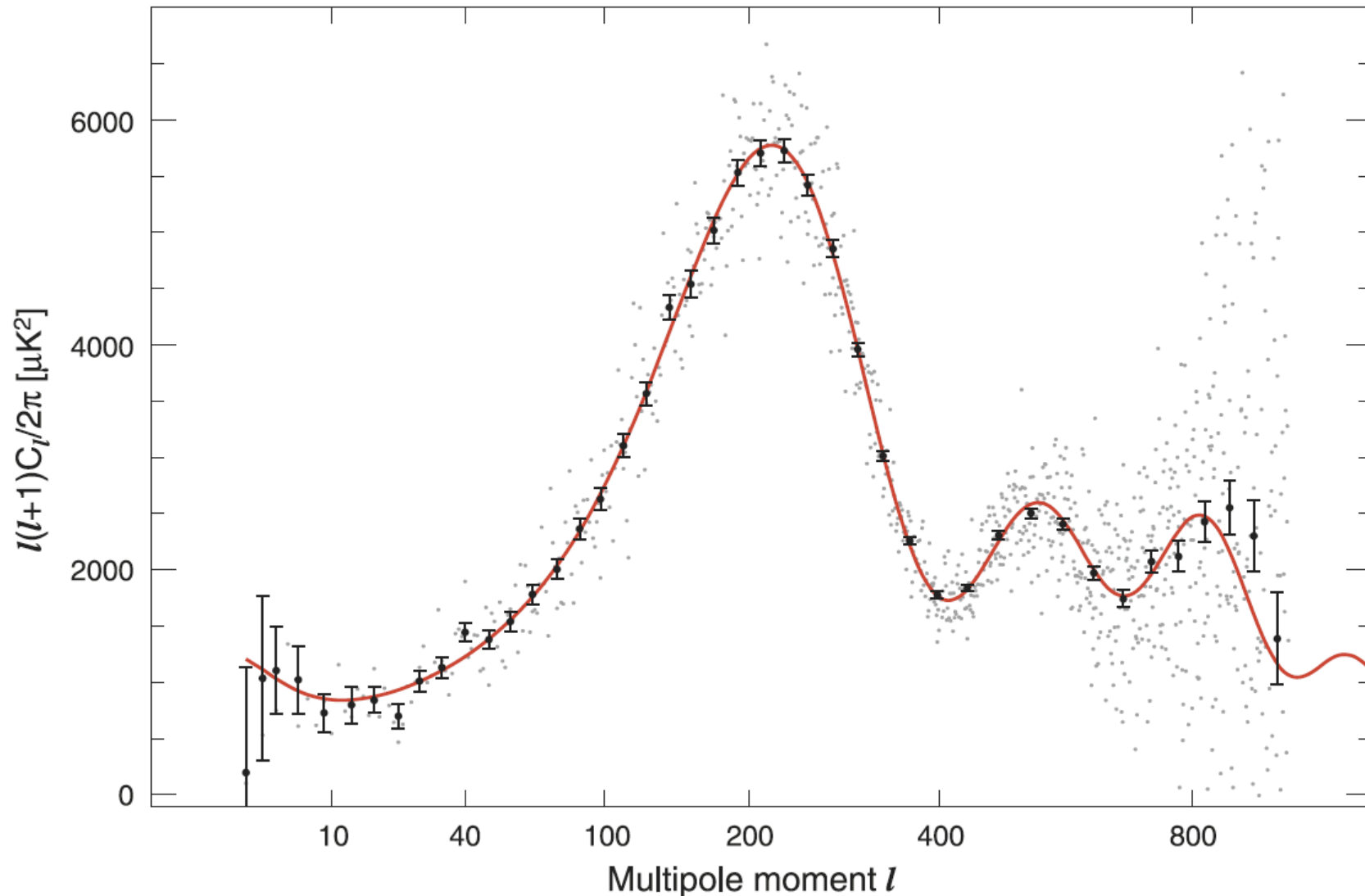
$$C_l \equiv \frac{1}{2l+1} \sum |a_{lm}|^2$$

$C_l$ 's mildly correlated since  
(due to Galactic foreground)  
only  $\sim 85\%$  of sky can be used



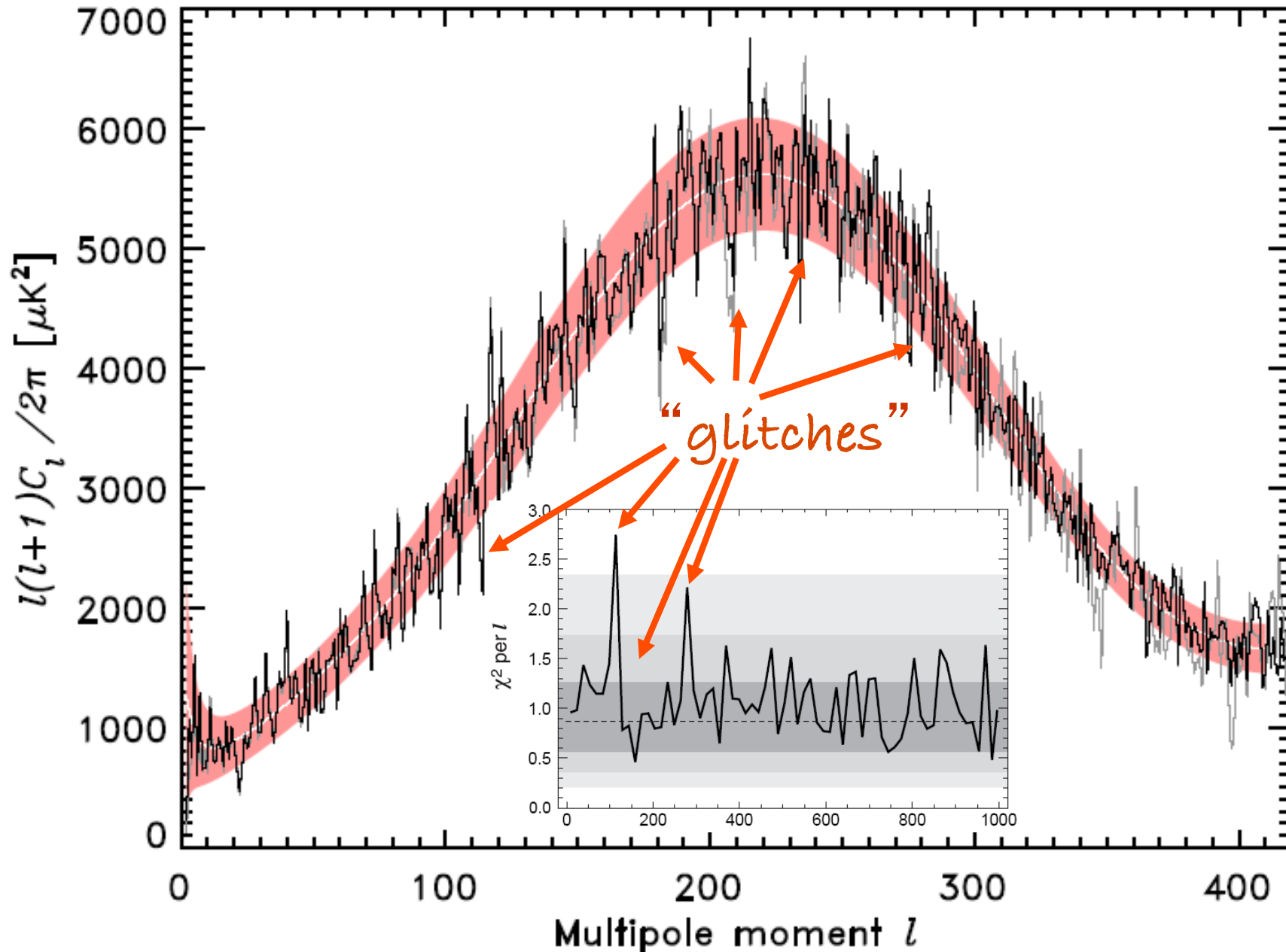
The 'power-law  $\Lambda$ CDM model' is believed to be confirmed by WMAP

Best-fit:  $\Omega_m h^2 = 0.11 \pm 0.01$ ,  $\Omega_b h^2 = 0.023 \pm 0.001$ ,  $h = 0.72 \pm 0.03$ ,  $n = 0.96 \pm 0.02$



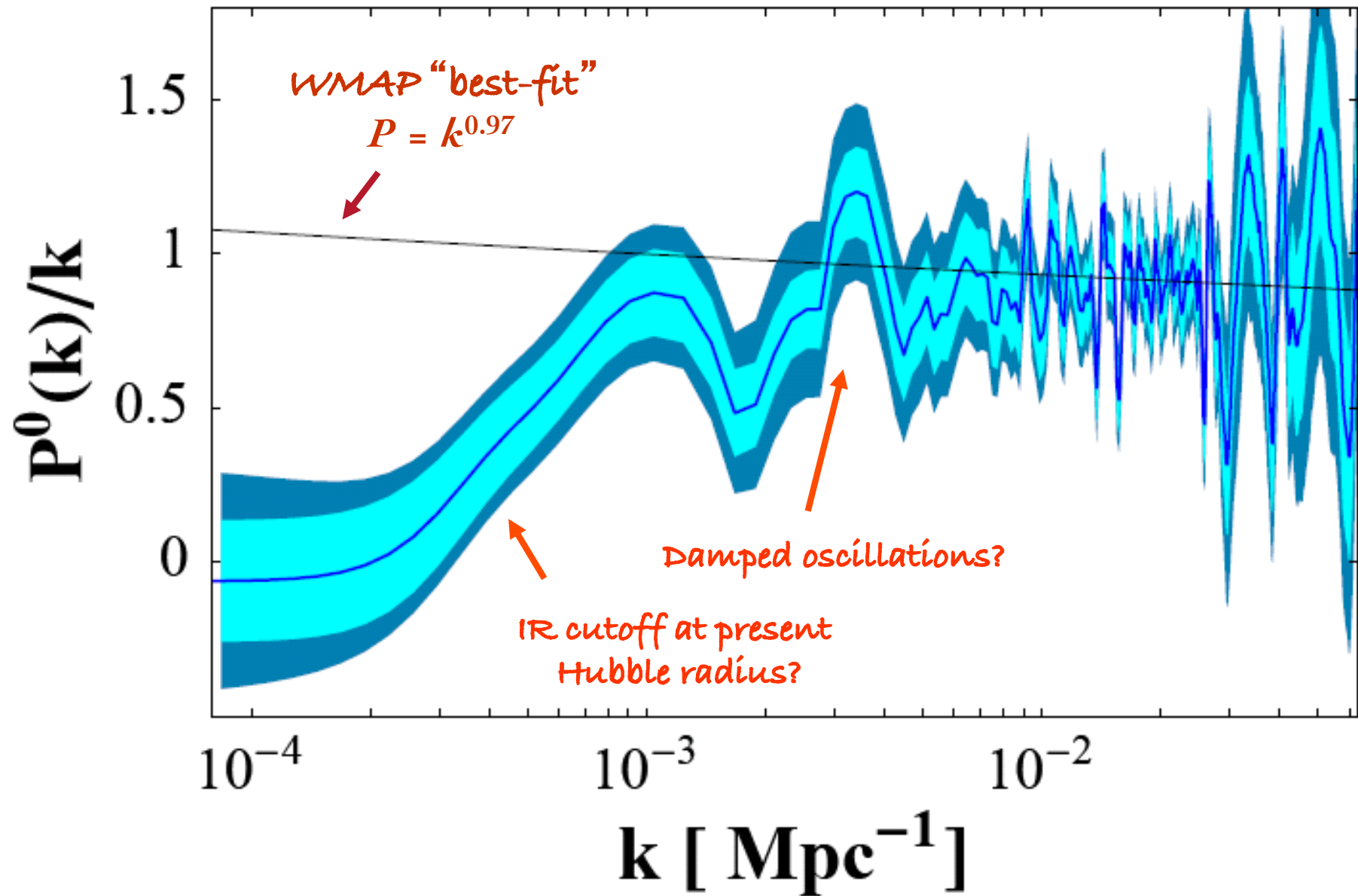
But the  $\chi^2/\text{dof} = 1049/982 \Rightarrow$  probability of  $\sim 7\%$  that this model describes the data

The excess  $\chi^2$  comes mostly from the outliers in the TT spectrum



Is the primordial density perturbation really scale-free?

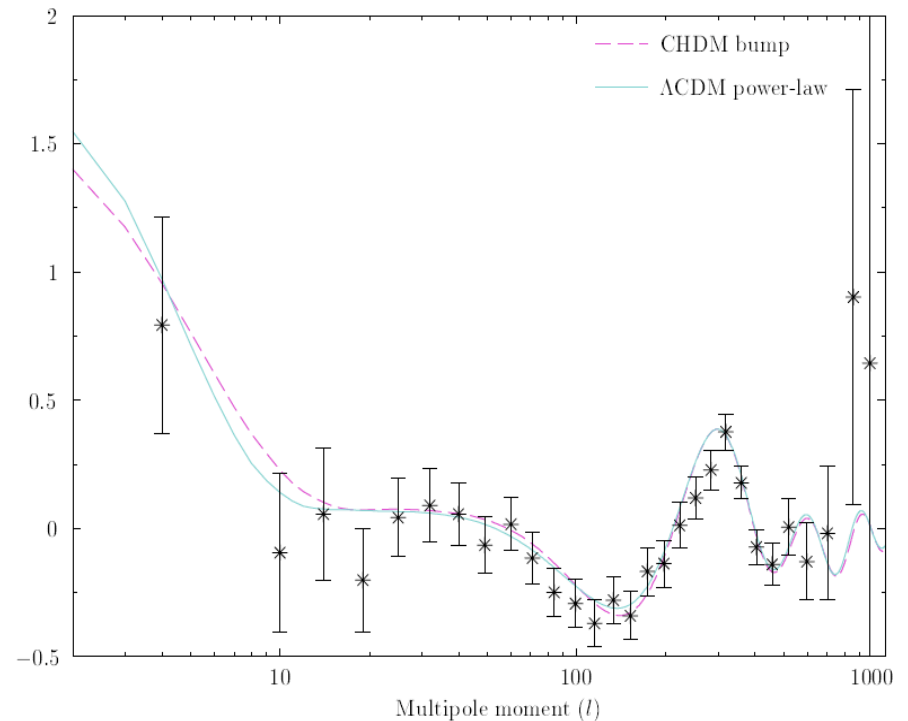
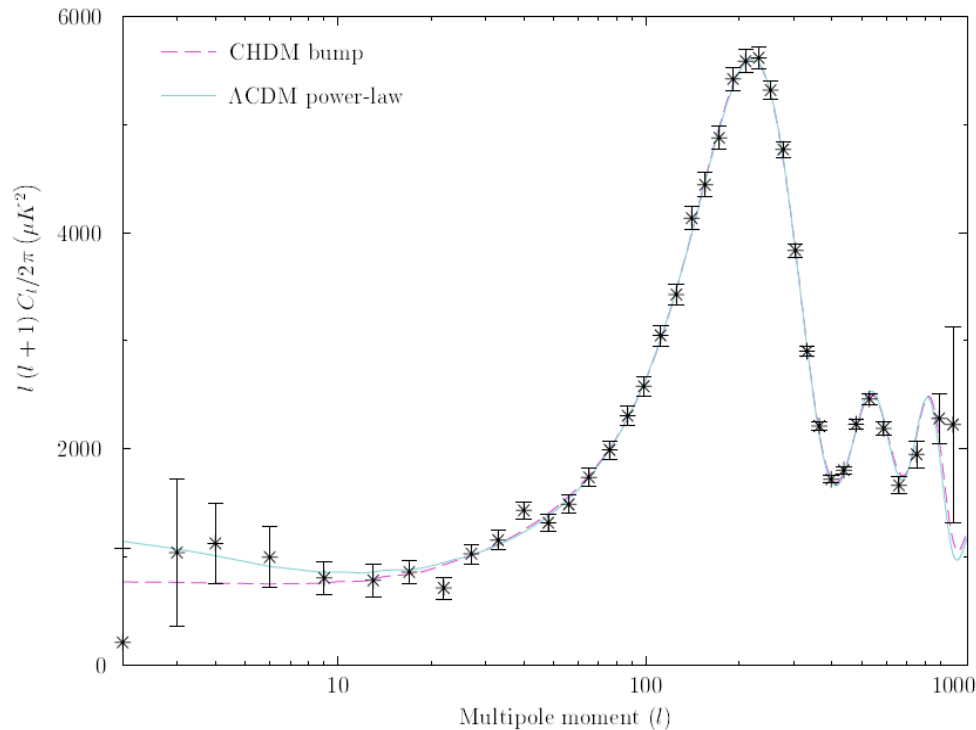
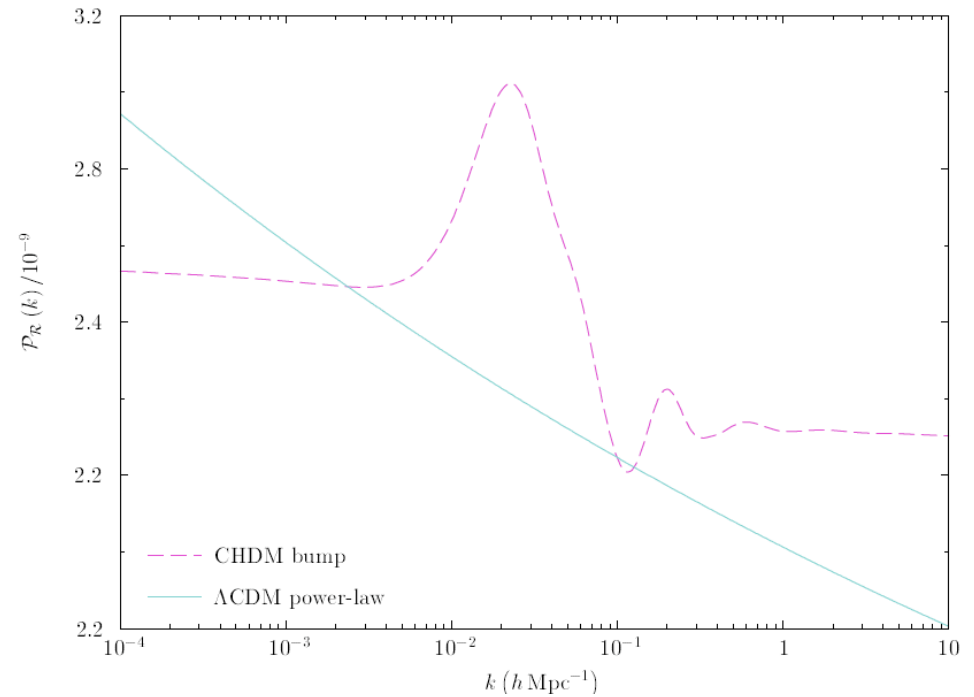
Many attempts made to reconstruct the primordial spectrum (assuming  $\Lambda$ CDM)  $\rightarrow$  indications for departures from a power-law spectrum



Tochhini-Valentini, Hoffman & Silk (2005)

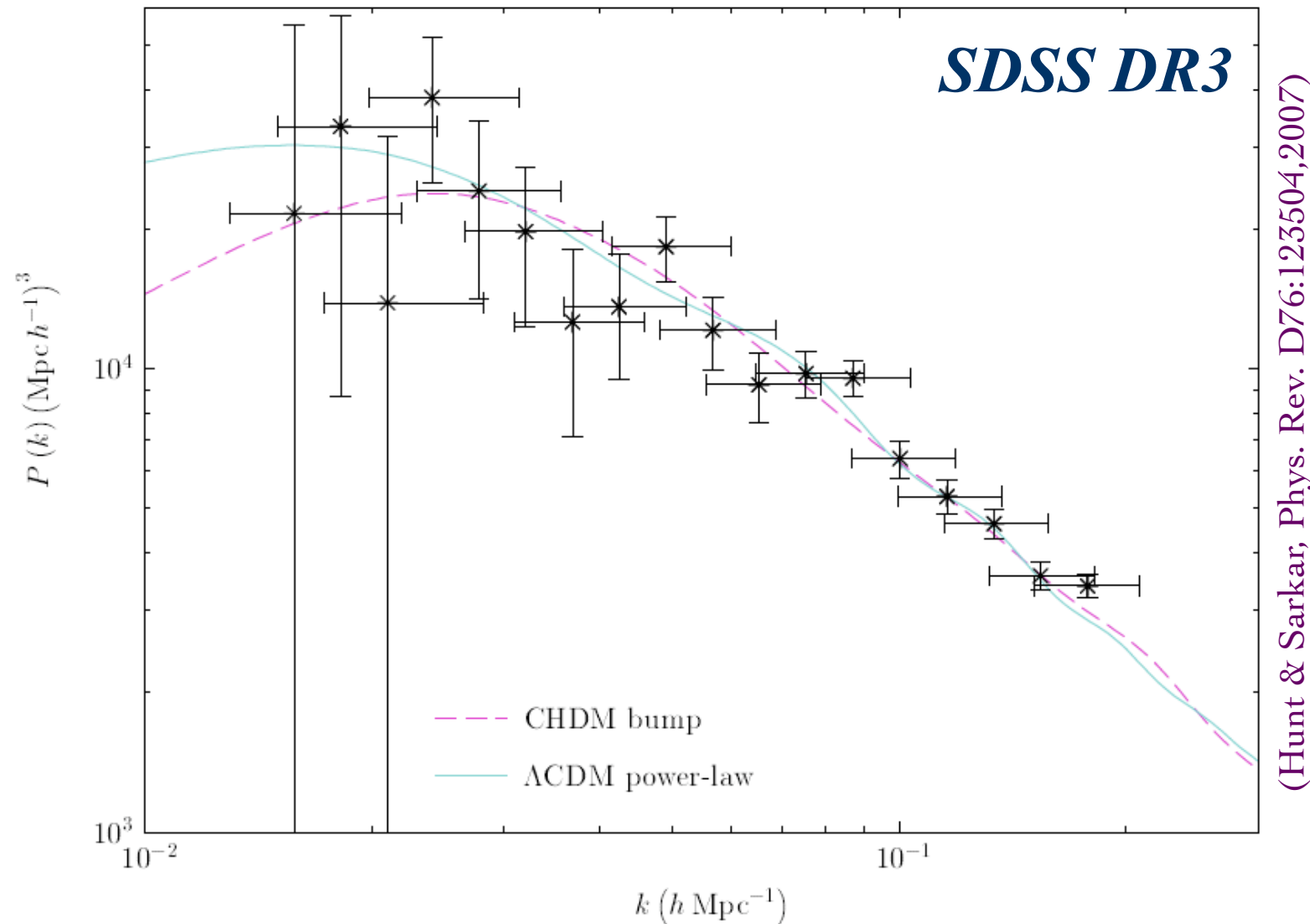
The primordial perturbation spectrum need not be scale-free as is commonly assumed

if there is a 'bump' in the spectrum, the WMAP data can be fitted with no dark energy ( $\Omega_m = 1, \Omega_\Lambda = 0$ ) if  $h \sim 0.44$



The small-scale power would be excessive unless damped by free-streaming ...

adding  $\nu$  of mass 0.5 eV ( $\Rightarrow \Omega_\nu \sim 0.1$ ) gives good match to large-scale structure



(*Hunt & Sarkar, Phys. Rev. D76:123504,2007*)

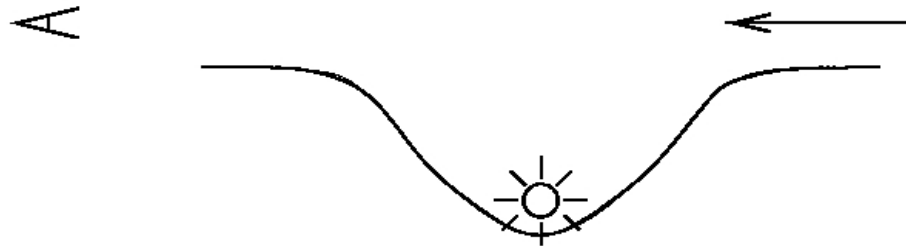
Fit gives  $\Omega_b h^2 \approx 0.018 \rightarrow$  BBN  $\checkmark \Rightarrow$  baryon fraction in clusters  $\sim 10\% \checkmark$   
... also excellent match to gravitational lensing signal in CMB

# Is there direct dynamical evidence for $\Lambda$ ?

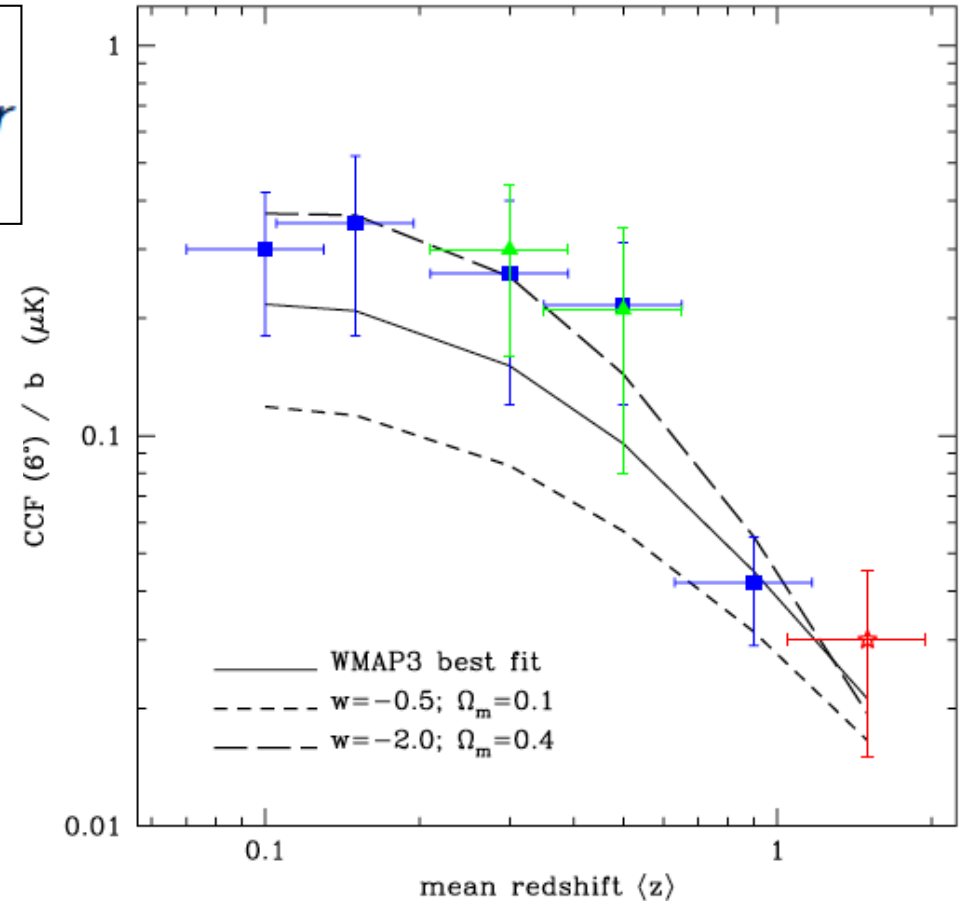
(‘late integrated Sachs-Wolfe effect’)

$$\frac{\Delta T(\hat{n})}{T_0} = \frac{2}{c^3} \int_0^{r_L} \dot{\Phi}(r, z, \hat{n}) a \, dr$$

gravitational potential traced by galaxy counts



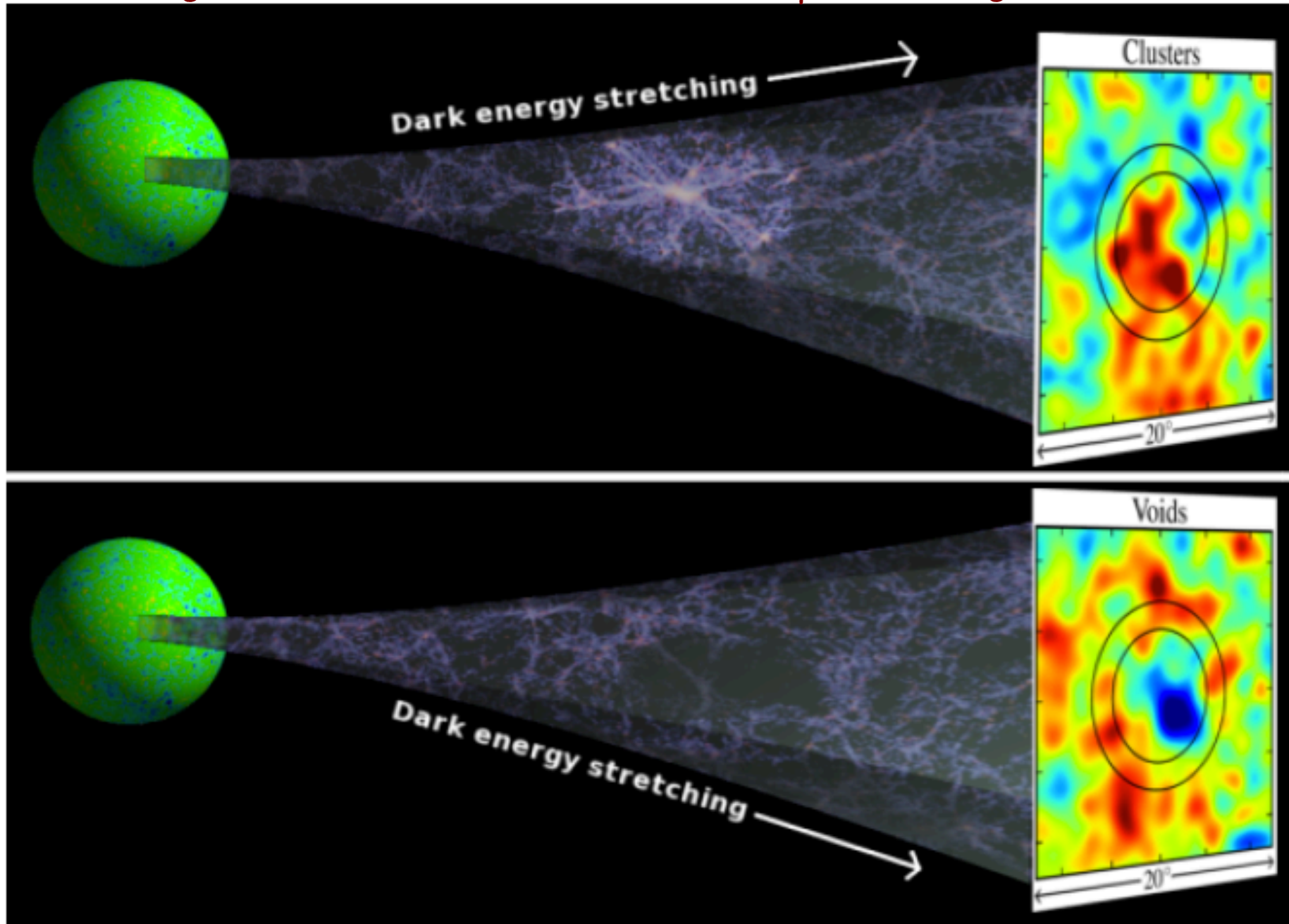
potential depth changes as CMB photons pass through



Gianantonio *et al* (2007)

Present detections are of low significance (2-3  $\sigma$ ) ... moreover the observed amplitude/ $z$ -dependence is higher/steeper than expected for  $\Lambda$

So it was big news when a  $>4\sigma$  detection was reported using SDSS DR6 LRGs



$$\Delta T = -11.3 \pm 3.1 \mu\text{K} \text{ for voids, } \Delta T = 7.9 \pm 3.0 \mu\text{K} \text{ for clusters}$$

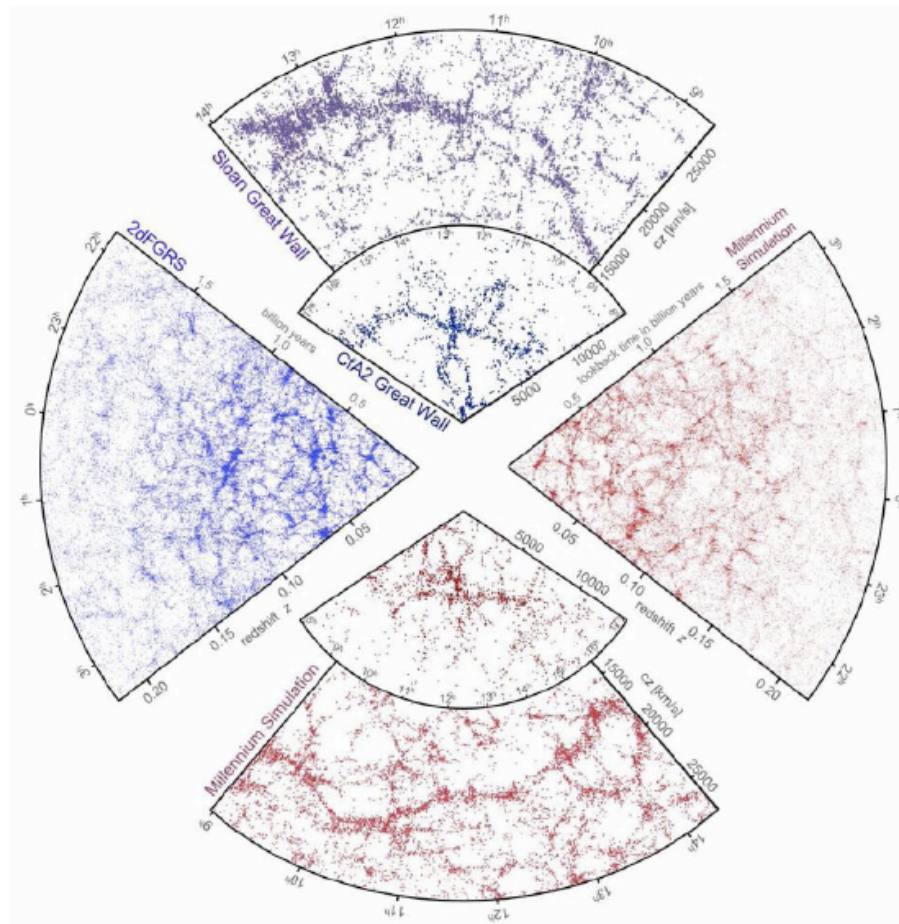


However expected signal (using compensated top-hat profile as expected for asymptotic evolution of large void) is only  $\Delta T \sim 0.1 \mu\text{K}$  ... to yield the observed signal the voids would have to be essentially empty - very unlikely in Gaussian density field! (Hunt & Sarkar, MNRAS, 401:547,2010)

More sophisticated treatment (using BBKS formalism for linear perturbations in Gaussian field) increases expectation somewhat but there is still  $>3\sigma$  discrepancy with observations (Nadathur et al, JCAP06:042,2012)

Although simulations of structure formation in  $\Lambda\text{CDM}$  are generally in good agreement with observations, the abundance of rare structures (rich clusters, giant voids) does seem to be a problem ... peculiar velocities too are higher than expected.

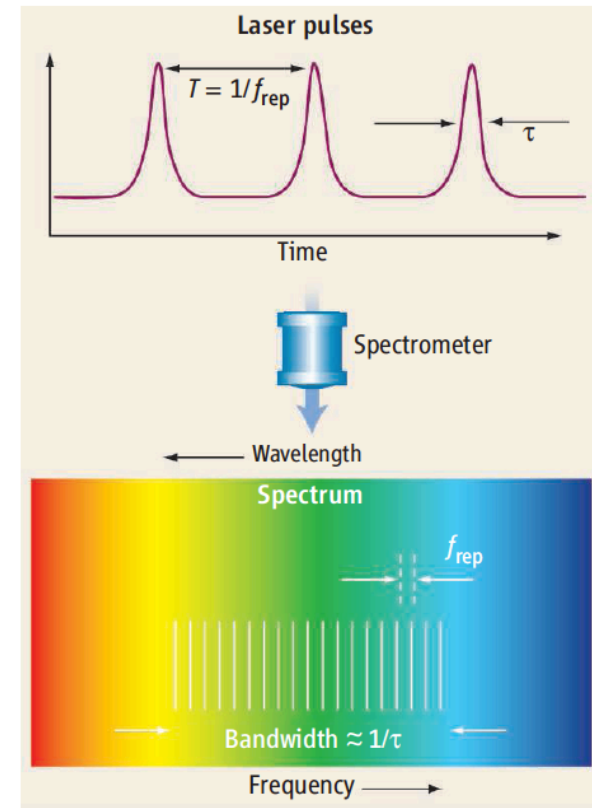
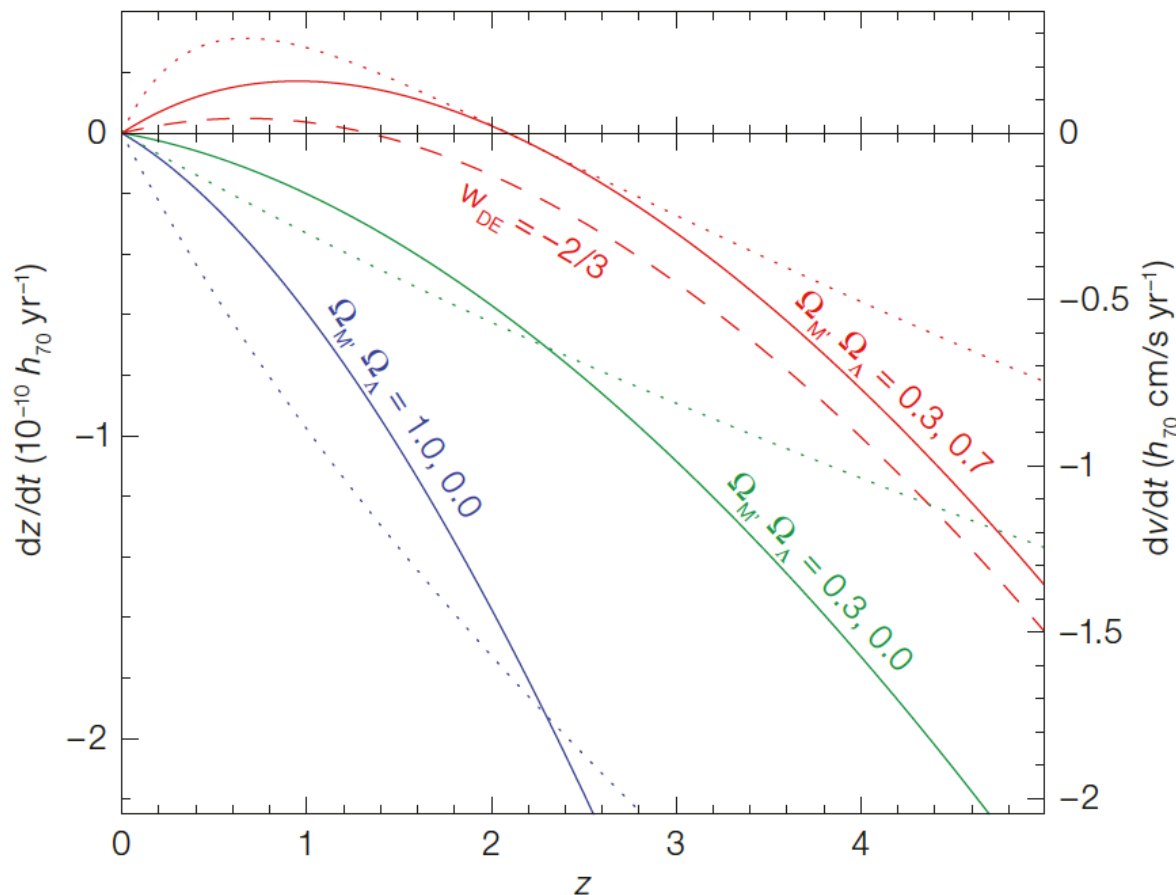
Forthcoming observations by *Euclid* of e.g. variations in the growth rate of structure with redshift will be crucial.



(Springel, Frenk, White 2007)

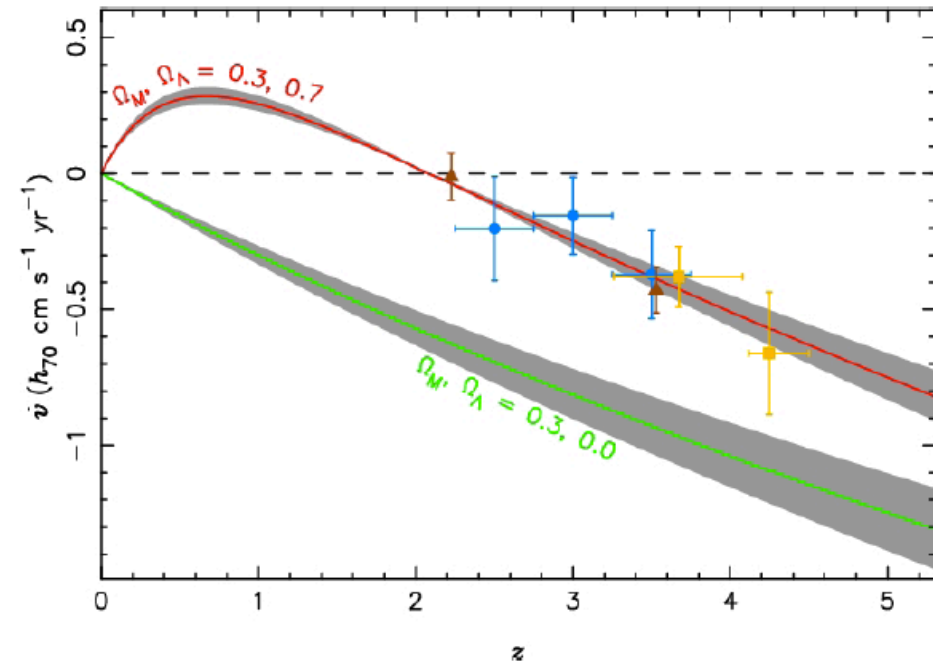
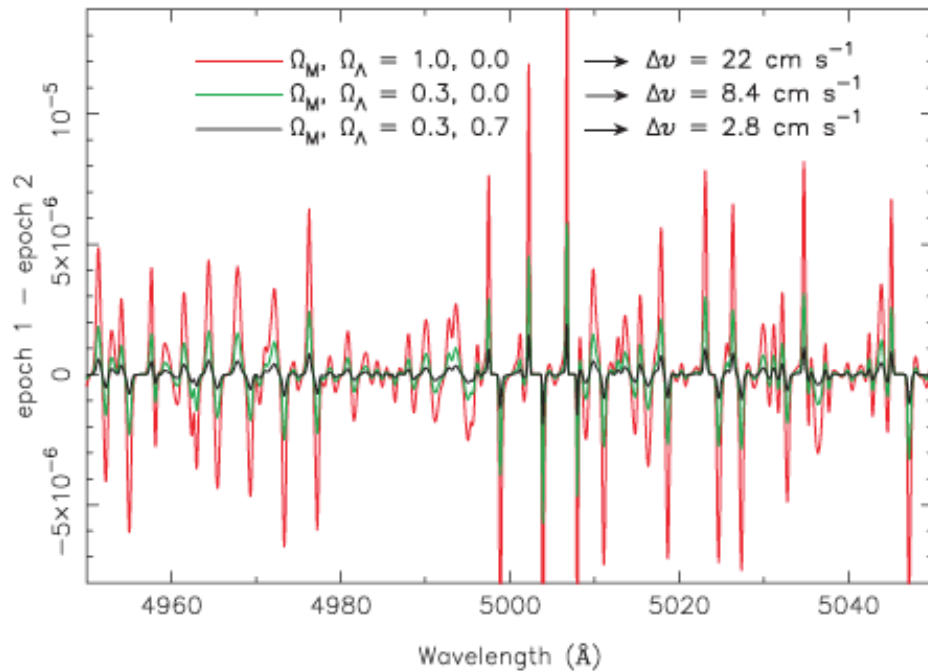
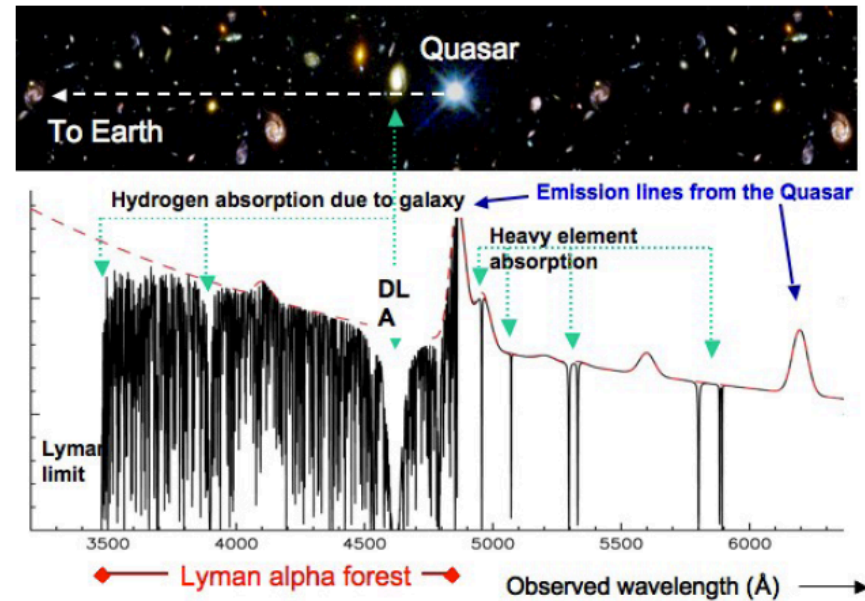
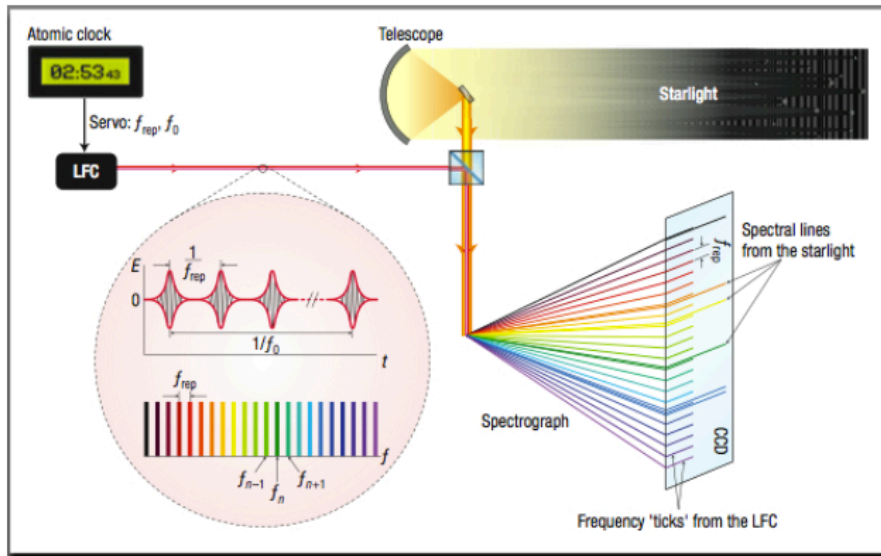
The evolution of the expansion rate can in principle be measured directly by monitoring the 'redshift drift' over time (Sandage, 1962):  $\dot{z} = (1+z)H_0 - H(z)$  ... however "With present optical techniques there is apparently no hope of detecting such small changes in redshifts for time intervals smaller than  $10^7$  years"

The development of the 'Laser Frequency Comb' for very accurate wavelength measurements has opened up the possibility of doing this in just  $\sim 10$ - $20$  yrs!



Loeb (1998) suggested that the 'Lyman-alpha forest' of absorption lines in the spectra of distant quasars provide the best target (peculiar motions are small)

Whether the expansion rate is indeed accelerating will be tested directly with the CODEX spectrograph on the European-Extremely Large Telescope



Liske et al, MNRAS 386:1192,2008

## Conclusions

There has been a renaissance in cosmology but modern data is still interpreted in terms of an idealised model whose basic assumptions have not been rigorously tested

The standard FRW model naturally admits  $\Lambda \sim H_0^2 \dots$   
and this is being interpreted as dark energy:  $\Omega_\Lambda \sim H_0^2 M_P^2$

Realistic models of our inhomogeneous universe may account for the SNIa Hubble diagram without acceleration

The CMB and LSS data can be equally well fitted if the primordial perturbations are not scale-free and  $m_\nu \sim 0.5 \text{ eV}$

Dark energy may just be an artifact of an oversimplified cosmological model