

de Sitter Holography with Higher Spin Gravity

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The Physics of de Sitter Spacetime
Albert Einstein Institut, Hannover
September 13, 2012

Motivation

de Sitter has many puzzles, classical and quantum:

- definition of observables
- perturbative and non-perturbative stability
- entropy of the de Sitter horizon
- etc.

The Goal

In this talk I will describe a consistent theory of gravity in de Sitter space, and ‘solve’ it.

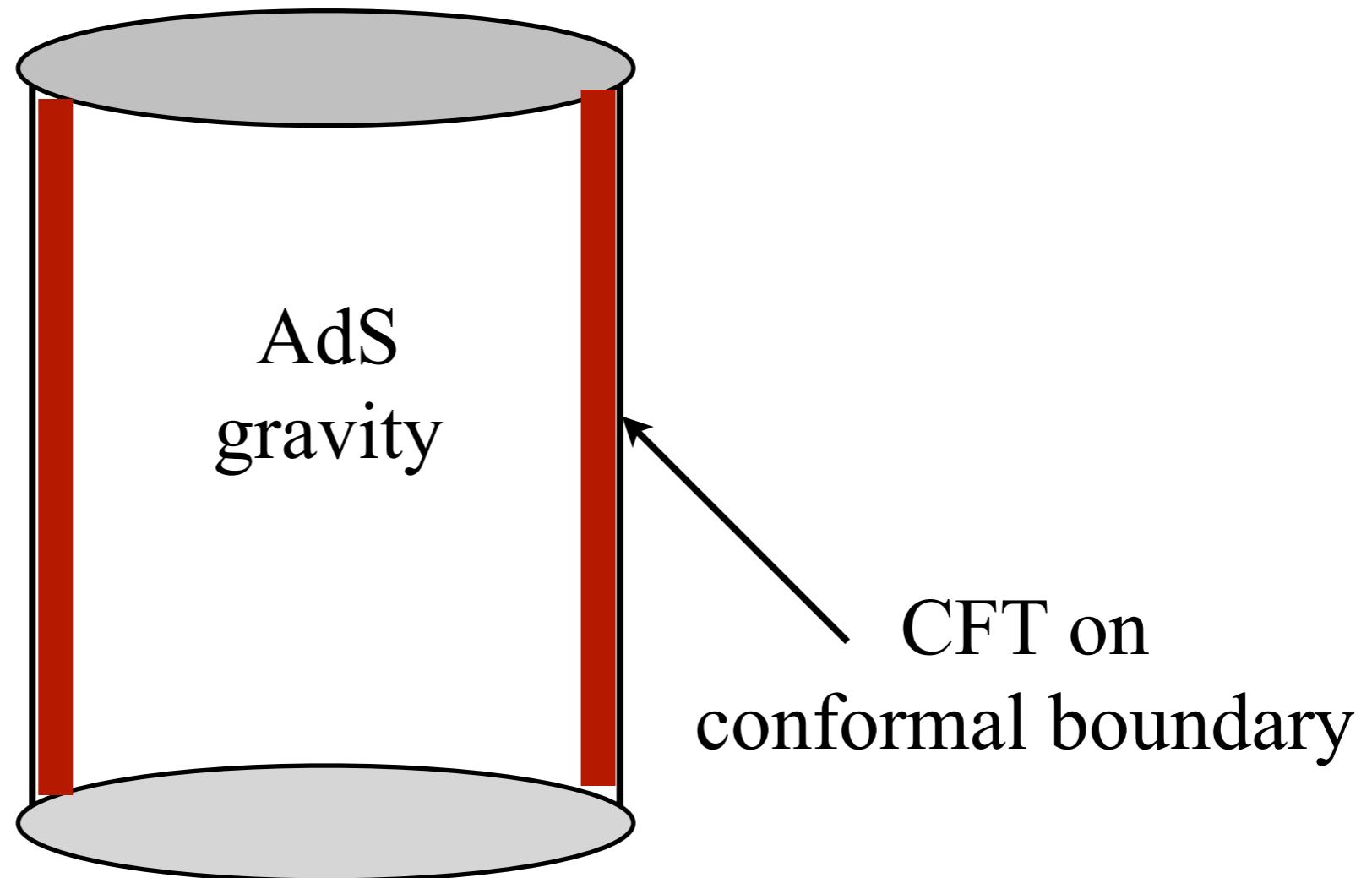
- Non-perturbative quantum gravity
- UV-complete
- Interacting but exactly solvable

however,

- It is a toy model
- In addition to gravity, it has an infinite number of massless particles of spin > 2
- (But is smaller than a string theory)
- meaning of “solve” is limited

Introduction

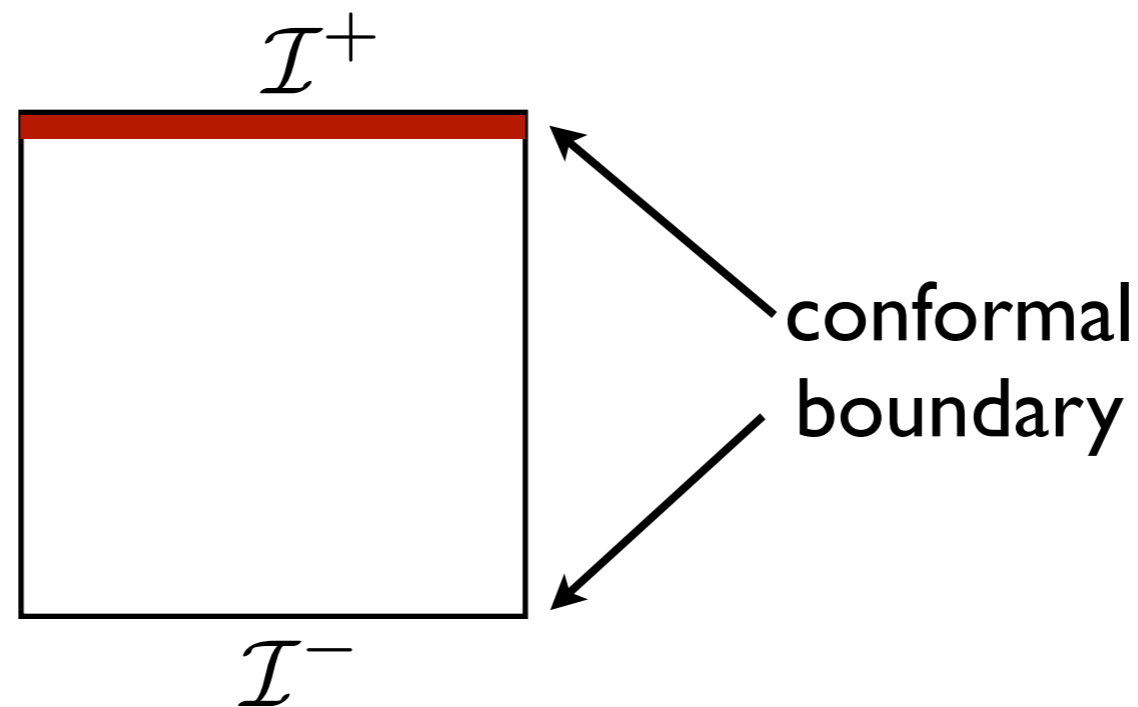
The AdS/CFT correspondence provides a non-perturbative definition of quantum gravity with negative cosmological constant.



But we live in (asymptotically) de Sitter space, in the past and future.

Introduction

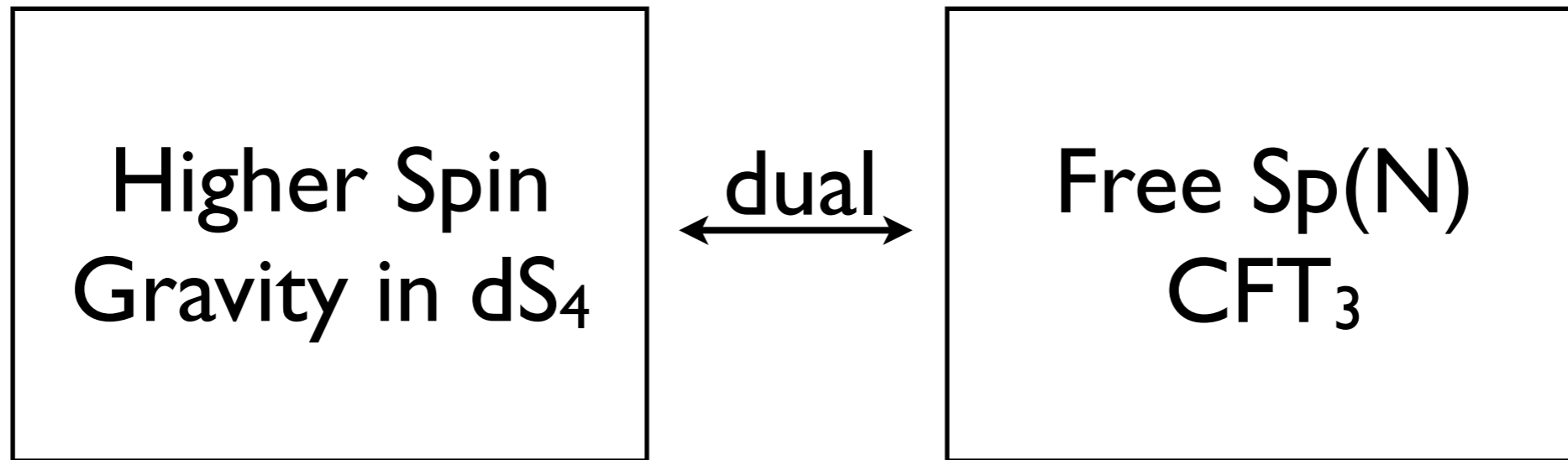
The dS/CFT correspondence, if it exists, is very different from AdS/CFT. Time is emergent:



A general dS/CFT dictionary has been proposed, but an explicit example was lacking.

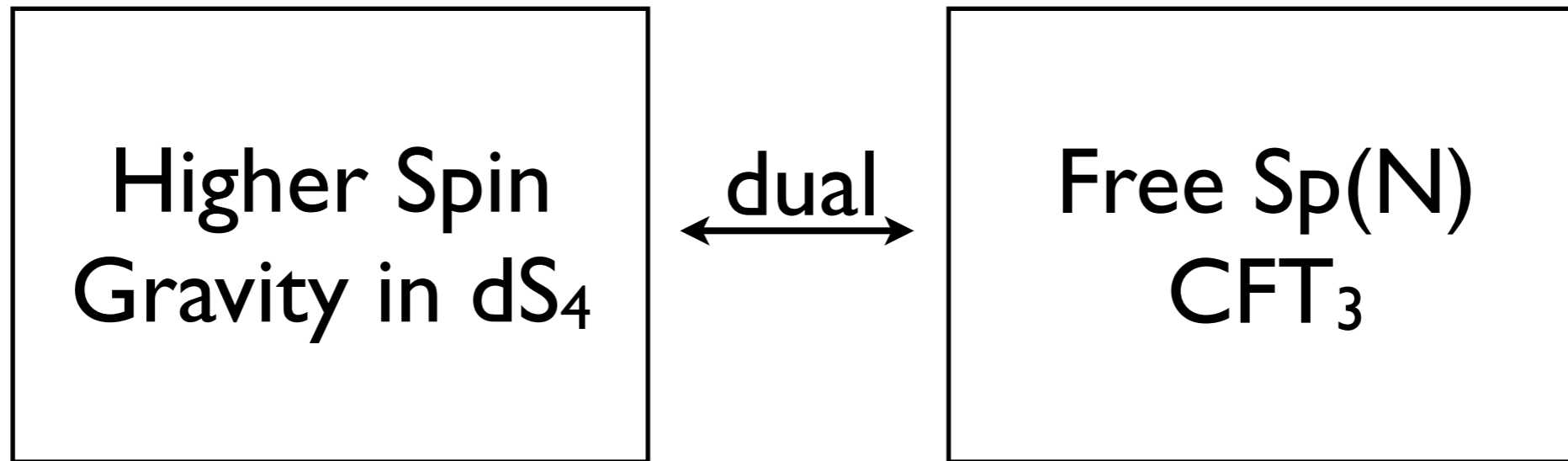
Witten '01; Strominger '01; Maldacena '02.

Introduction



$$\Lambda \sim 1/N$$

Introduction



$$\Lambda \sim 1/N$$

Outline of the talk

- Higher spin gravity
- The Sp(N) CFT
- Duality

I. Higher Spin Gravity

Higher Spin Gravity

Gravity plus large (or infinite) number of massless fields,

$$A_{\mu_1 \cdots \mu_s}$$

with spins

$$s = 0, 1, 2, 3, 4, \cdots ,$$

Higher Spin Gravity

Gravity plus large (or infinite) number of massless fields,

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with spins

$$s = 0, 1, 2, 3, 4, \cdots ,$$

Spin-2 = graviton. Massless higher spin fields mean very large gauge symmetry extending diffeomorphism invariance.

Consistent interacting theory exists for $\Lambda \neq 0$

Fradkin and Vasiliev, 1987

Vasiliev, 1990

Equations of HS Gravity

The equations of motion are known, but very complicated.

The action is unknown, but has higher-derivative interactions of the schematic form:

$$S \sim \int \frac{1}{\Lambda^\#} \partial^n X \partial^m Y \partial^k Z \dots$$

$\Lambda \ell_P^2$ is the only parameter.

Reasons for Higher Spin

- This is a limit of string theory, where the string tension is taken to zero (i.e. on very small geometry). So it is a testing ground for quantum gravity in the stringy regime.
- It has a simple holographic description: Free fields!
 - ▶ the higher spin gauge fields have a dual description as the conserved currents in free field theory,

$$A_{\mu_1 \cdots \mu_s} \quad \leftrightarrow \quad \phi \partial_{(\mu_1} \cdots \partial_{\mu_s)} \phi$$

II. The $Sp(N)$ CFT

The 3d Sp(N) Model

N free anticommuting scalars (“ghosts”)

$$S_{cft} = \frac{1}{2} \int d^3x \Omega_{ab} \partial\chi^a \cdot \partial\chi^b$$

$\Omega \equiv$ antisym. symplectic form

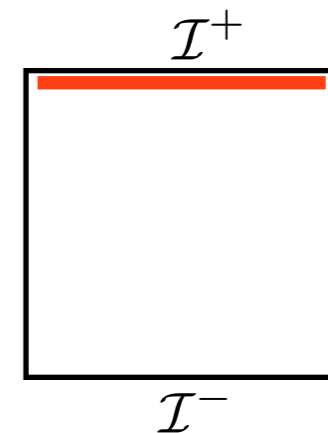
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- In Lorentzian signature, this would be non-unitary.
- Our CFT is Euclidean; unitarity is not an issue.
- Nonetheless unitarity of the gravitational theory should be encoded somehow in the CFT.



Currents

The $\text{Sp}(N)$ model has a conserved current of every even spin:

$$J_{\mu_1 \cdots \mu_s} = \Omega_{ab} \chi^a \partial_{(\mu_1} \cdots \partial_{\mu_s)} \chi^b + \cdots$$

We restrict to the singlet sector (by gauging and sending the gauge coupling to zero).

Then at large N , these currents are the only local operators of small dimension. This is precisely the spectrum of higher spin gravity.

The Claim

Higher spin gravity in $dS_4 = Sp(N)$ CFT

$$\Lambda \sim \frac{1}{N}$$

- Conjecture was based on 3pt correlators

Anninos, TH, Strominger 2011

- Now proven for all n-pt functions to all orders in perturbation theory

Maldacena, Zhiboedov 2011

O(N) duality

An AdS/CFT Example

- Higher spin gravity in AdS is dual to free bosons, ie the O(N) model

$$\phi^a \partial \dots \partial \phi^a$$

- 3pt correlators were computed by Giombi and Yin

- The dS/CFT example in this talk is constructed by modifying this duality for use in de Sitter. This procedure does not work in string theory examples of AdS/CFT.

Fronsdal '79

Witten

Sundborg

Mikhailov

Sezgin & Sundell

Klebanov & Polyakov

Giombi and Yin

...

Petkou; Giombi and Yin

III. dS/CFT Duality

Statement of dS/CFT

In words:

The CFT computes the Hartle-Hawking/Bunch-Davies wavefunction of the gravity theory.

Statement of the Duality

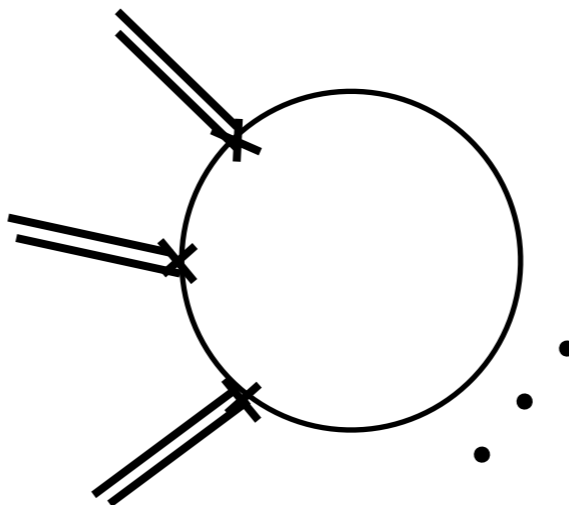
The statement of dS/CFT is:

Witten '01
Strominger '01
Maldacena '02

$$\boxed{Z_{cft}[X_0] = \Psi_{grav}[X_0]}$$
$$= \int_{X_0} DX e^{iS_{grav}[X]}$$

- X is a field in de Sitter
- X_0 is its value at the future boundary \mathcal{I}^+
$$X \sim X_0(x)\eta^{3-\Delta}$$
- In the CFT, X_0 is the source for the dual higher-spin current

Computing CFT Correlators

$$\langle J^{(s_1)} J^{(s_2)} \dots J^{(s_n)} \rangle_{cft} = \text{Diagram}$$


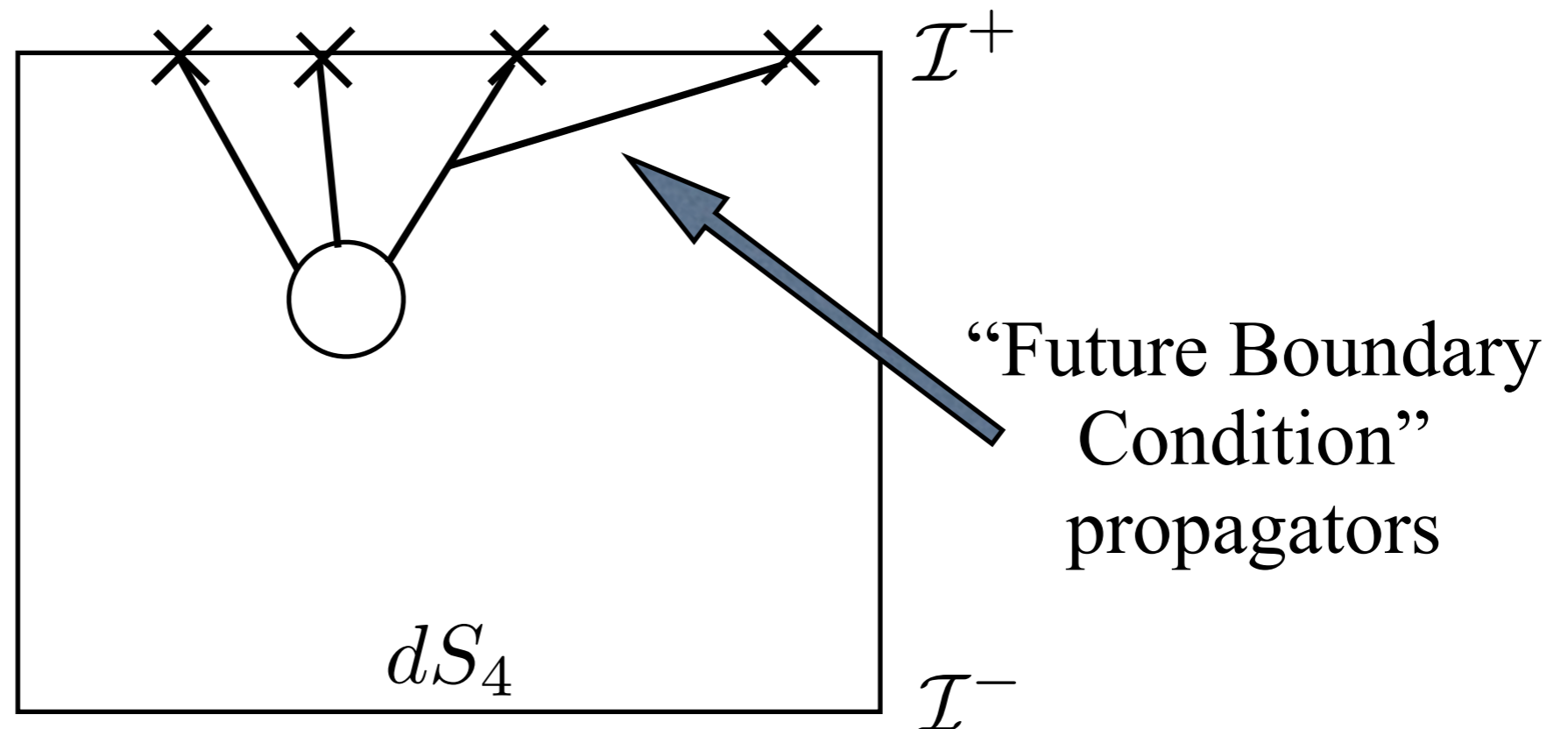
This is a trivial free-field Feynman diagram.

Computing the Bulk Wavefunction

Semiclassically, the HH wavefunction is

$$\Psi[X_0] \approx e^{iS_{bulk}[X_0]}$$

This can also be computed diagrammatically:



Aside: In-In

These are not In-In correlators.

Maldacena '02

The wavefunction is only the first half of an In-In computation:

$$\langle X_1 X_2 \cdots \rangle_{in-in} = \int D X_0 |\Psi_{grav}|^2 X_1 X_2 \cdots$$

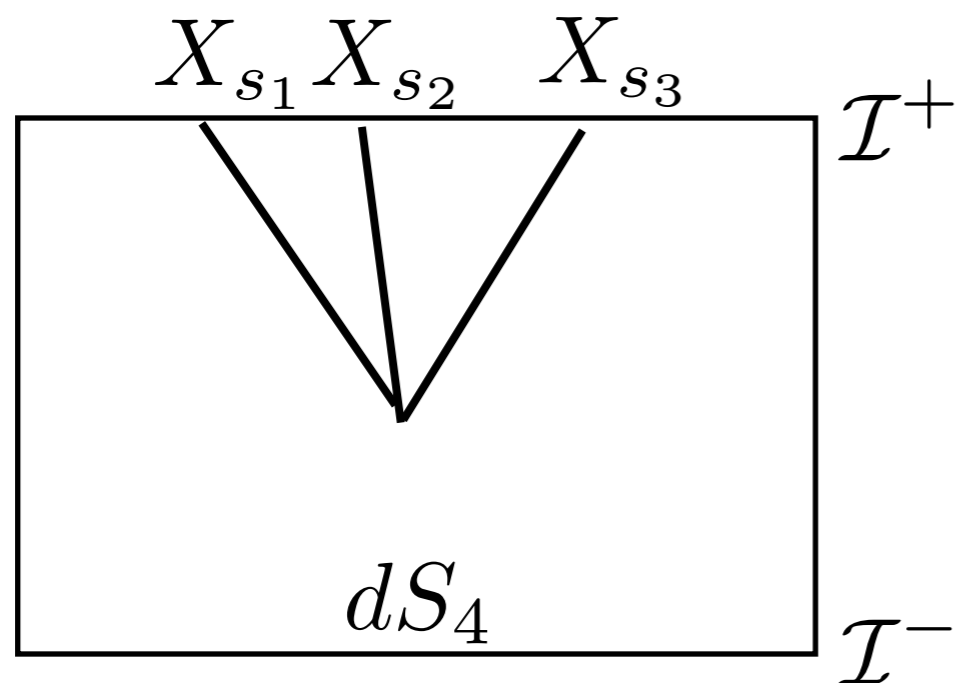
Therefore

CFT correlators \neq gravity correlators

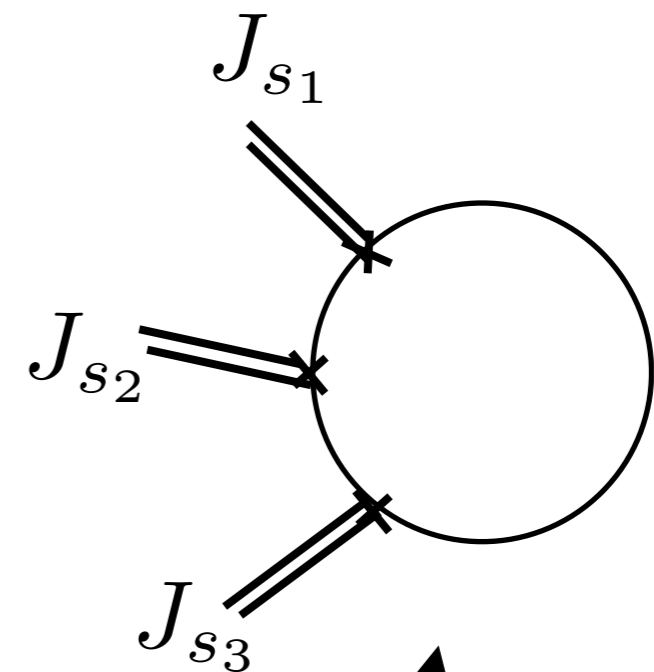
To compute gravity correlators, you must allow fields to fluctuate at \mathcal{I}^+

Results

By explicit computation,



=



↑
computed in
higher spin gravity

↑
computed in
free $Sp(N)$ CFT

Comments

- This is the cubic part of

$$Z_{cft}[X_0] = \Psi_{grav}[X_0]$$

- The most difficult part is finding the propagators and interaction vertices of higher spin gravity. This was done by Giombi and Yin in AdS; we analytically continued to produce de Sitter results.

$$N \rightarrow -N , \quad \Lambda \rightarrow -\Lambda$$

- Under some assumptions, Maldacena and Zhiboedov used constraints from higher spin symmetry to prove equality for all n-point correlators.

Conclusion

Higher spin gravity is dual to a free CFT and may allow a UV-complete holographic duality in de Sitter.

Things you might hope to compute:

- RG flows corresponding to universe production
- non-perturbative wavefunction of the universe
- de Sitter entropy
- static patch observables

Anninos, Denef, Harlow '12

Ultimately, seeking general rules for de Sitter holography.