

Aspects of string phenomenology in the LHC era

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- High string scale, SUSY and 125 GeV Higgs
- Low scale strings and extra dimensions
- Extra $U(1)$'s
- Tiny string coupling and linear dilaton background

Connect string theory to the real world:

What is the value of the string scale M_s ?

- arbitrary parameter : Planck mass $M_P \longrightarrow$ TeV

- physical motivations \Rightarrow favored energy regions:

- High : $\begin{cases} M_P^* \simeq 10^{18} \text{ GeV} & \text{Heterotic scale} \\ M_{\text{GUT}} \simeq 10^{16} \text{ GeV} & \text{Unification scale} \end{cases}$

- Intermediate : around 10^{11} GeV ($M_s^2/M_P \sim \text{TeV}$)

SUSY breaking, strong CP axion, see-saw scale

- Low : TeV (hierarchy problem)

Beyond the Standard Model of Particle Physics: driven by the mass hierarchy problem

Standard picture: low energy supersymmetry

Natural framework: Heterotic string (or high-scale M/F) theory

Advantages:

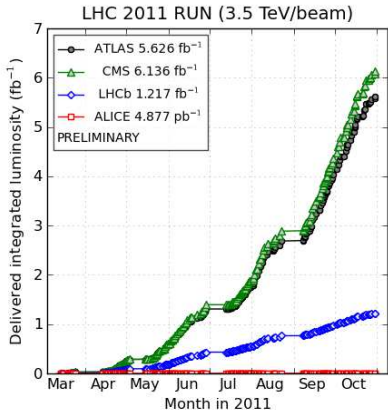
- natural elementary scalars
- gauge coupling unification
- LSP: natural dark matter candidate
- radiative EWSB

Problems:

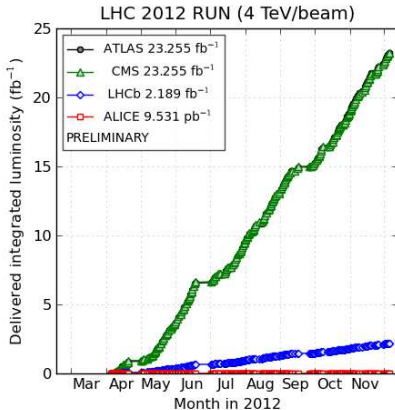
- too many parameters: soft breaking terms
- MSSM : already a % - ‰ fine-tuning 'little' hierarchy problem

Excellent LHC performance

Number of events = Cross section \times Luminosity



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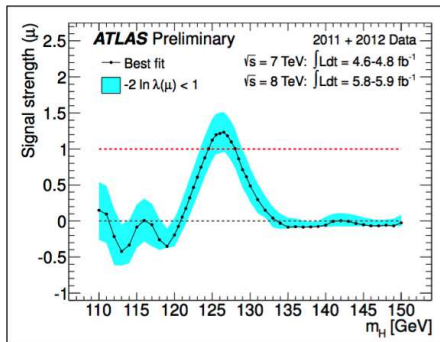


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New boson discovery at LHC:

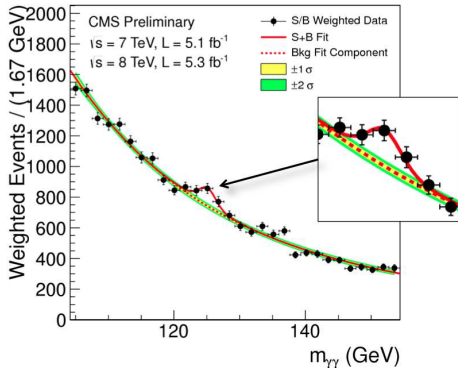
is it the scalar remnant of the Brout-Englert-Higgs mechanism

for breaking the EW symmetry?



$$m_H = 125.5 \pm 0.2 \text{ (stat.)} \pm 0.5 \text{ (syst.)}$$

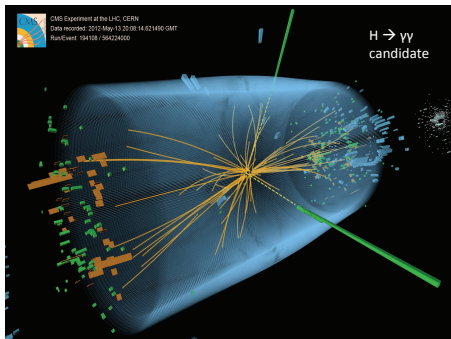
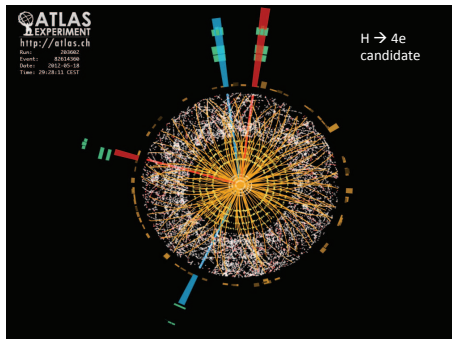
$\gtrsim 6\sigma$



$$m_H = 125.8 \pm 0.4 \pm 0.4 \text{ GeV}$$

6.8 σ significance

Possible Higgs boson events



ATLAS SUSY Searches* - 95% CL Lower Limits (Status: ICHEP 2012)

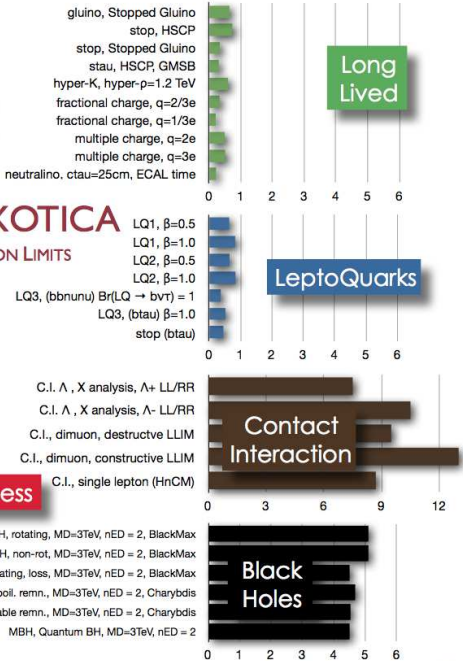
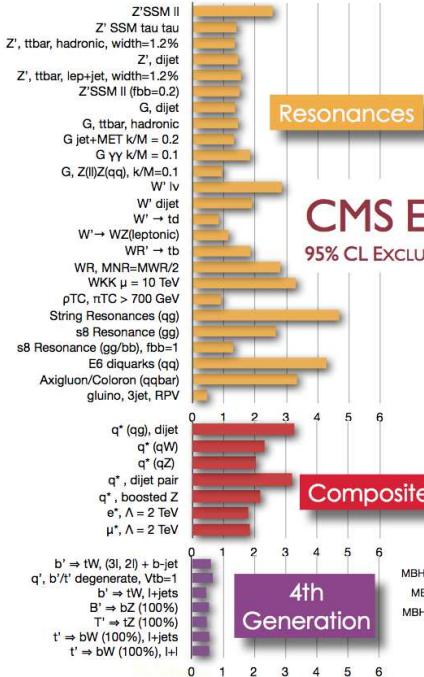
		1.0 TeV	$\tilde{g} = \tilde{g}$ mass
Inclusive searches	MSUGRA/CMSSM : 0 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-033]	1.40 TeV $\tilde{g} = \tilde{g}$ mass
	MSUGRA/CMSSM : 1 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-041]	1.20 TeV $\tilde{g} = \tilde{g}$ mass
	MSUGRA/CMSSM : 0 lep + multijets + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1206.1760]	840 GeV \tilde{g} mass (large m_0)
	Pheno model : 0 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-033]	1.38 TeV \tilde{g} mass ($m_{\tilde{G}} < 2$ TeV, light $\tilde{\chi}_1^0$)
	Pheno model : 0 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-033]	940 GeV \tilde{g} mass ($m_{\tilde{G}} < 2$ TeV, light $\tilde{\chi}_1^0$)
	Glauino med. $\tilde{\chi}^{\pm} (\tilde{g} \rightarrow \tilde{g}\tilde{\chi}^{\pm})$: 1 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-041]	900 GeV \tilde{g} mass ($m_{\tilde{G}} < 200$ GeV, $m_{\tilde{G}} = \frac{1}{2}(m_{\tilde{G}}^{\pm} + m_{\tilde{G}}^0)$)
	GMSB : 2 lep OSSF + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2011-156]	810 GeV \tilde{g} mass ($\tan\beta < 35$)
	GMSB : 1 τ + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1204.3952]	920 GeV \tilde{g} mass ($\tan\beta > 20$)
	GMSB : 2 τ + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1203.6580]	990 GeV \tilde{g} mass ($\tan\beta > 20$)
	GGM : $\gamma\gamma$ + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-072]	1.07 TeV \tilde{g} mass ($m_{\tilde{G}} > 50$ GeV)
3rd gen. squarks gluino mediated	$\tilde{g} \rightarrow b\tilde{b}^*$ (virtual \tilde{b}) : 0 lep + 1/2 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1203.6193]	900 GeV \tilde{g} mass ($m_{\tilde{G}} < 300$ GeV)
	$\tilde{g} \rightarrow b\tilde{b}^*$ (virtual \tilde{b}) : 0 lep + 3 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-058]	1.02 TeV \tilde{g} mass ($m_{\tilde{G}} < 400$ GeV)
	$\tilde{g} \rightarrow b\tilde{b}^*$ (real \tilde{b}) : 0 lep + 3 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-058]	1.00 TeV \tilde{g} mass ($m_{\tilde{G}} = 60$ GeV)
	$\tilde{g} \rightarrow t\tilde{t}^*$ (virtual \tilde{t}) : 1 lep + 1/2 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1203.6193]	710 GeV \tilde{g} mass ($m_{\tilde{G}} < 150$ GeV)
	$\tilde{g} \rightarrow t\tilde{t}^*$ (virtual \tilde{t}) : 2 lep (SS) + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1203.5763]	650 GeV \tilde{g} mass ($m_{\tilde{G}} < 210$ GeV)
	$\tilde{g} \rightarrow t\tilde{t}^*$ (virtual \tilde{t}) : 0 lep + multi-j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1206.1760]	870 GeV \tilde{g} mass ($m_{\tilde{G}} < 100$ GeV)
	$\tilde{g} \rightarrow t\tilde{t}^*$ (virtual \tilde{t}) : 0 lep + 3 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-058]	940 GeV \tilde{g} mass ($m_{\tilde{G}} < 50$ GeV)
	$\tilde{g} \rightarrow t\tilde{t}^*$ (real \tilde{t}) : 0 lep + 3 b -j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-058]	820 GeV \tilde{g} mass ($m_{\tilde{G}} = 60$ GeV)
	$bb, b \rightarrow b\tilde{b}^*$: 0 lep + 2 b -jets + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1112.3832]	390 GeV \tilde{b} mass ($m_{\tilde{G}} < 60$ GeV)
	$\tilde{t}\tilde{t}$ (very light), $\tilde{t} \rightarrow b\tilde{b}^*$: 2 lep + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-059]	135 GeV \tilde{t} mass ($m_{\tilde{G}} = 45$ GeV)
3rd gen. squarks direct production	$\tilde{t}\tilde{t}$ (light), $\tilde{t} \rightarrow b\tilde{b}^*$: 1/2 lep + b-jet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-070]	120-173 GeV \tilde{t} mass ($m_{\tilde{G}} = 45$ GeV)
	$\tilde{t}\tilde{t}$ (heavy), $\tilde{t} \rightarrow b\tilde{b}^*$: 0 lep + b-jet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-074]	380-465 GeV \tilde{t} mass ($m_{\tilde{G}} = 0$)
	$\tilde{t}\tilde{t}$ (heavy), $\tilde{t} \rightarrow t\tilde{b}^*$: 1 lep + b-jet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-073]	230-440 GeV \tilde{t} mass ($m_{\tilde{G}} = 0$)
	$\tilde{t}\tilde{t}$ (heavy), $\tilde{t} \rightarrow t\tilde{b}^*$: 2 lep + b-jet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-071]	298-305 GeV \tilde{t} mass ($m_{\tilde{G}} = 0$)
	$\tilde{t}\tilde{t}$ (GMSB) : $Z \rightarrow \ell\ell$ + b-jet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1204.6736]	310 GeV \tilde{t} mass ($115 < m_{\tilde{G}} < 230$ GeV)
	$\tilde{t}\tilde{t}$ (GMSB) : $Z \rightarrow \ell\ell$ + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-076]	93-180 GeV \tilde{t} mass ($m_{\tilde{G}} = 0$)
	$\tilde{t}\tilde{t}$ (GMSB) : $Z \rightarrow \ell\ell$ + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-076]	120-330 GeV \tilde{t} mass ($m_{\tilde{G}} = 0, m_{\tilde{G}} = \frac{1}{2}(m_{\tilde{G}}^{\pm} + m_{\tilde{G}}^0)$)
	$\tilde{t}\tilde{t}$ (GMSB) : $Z \rightarrow \ell\ell$ + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [CONF-2012-077]	60-500 GeV \tilde{t} mass ($m_{\tilde{G}} = m_{\tilde{G}}^{\pm}, m_{\tilde{G}}^0 = 0, m_{\tilde{G}}^{\pm}$ as above)
	AMSb : long-lived $\tilde{\chi}_1^0$	$\sqrt{s} = 7$ TeV [CONF-2012-034]	518 GeV $\tilde{\chi}_1^0$ mass ($1 < \tau(\tilde{\chi}_1^0) < 2$ ns, 90 GeV limit in $[0.2, 0.90]$ ns)
	Long-lived particles	Stable \tilde{g} R-hadrons : Full detector	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-075]
Stable \tilde{b} R-hadrons : Full detector		$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-075]	612 GeV \tilde{b} mass
Stable \tilde{t} R-hadrons : Full detector		$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-075]	683 GeV \tilde{t} mass
Metastable \tilde{g} R-hadrons : Pixel det. only		$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-075]	910 GeV \tilde{g} mass ($\tau(\tilde{g}) > 10$ ns)
GMSB ; stable $\tilde{\tau}$		$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-075]	310 GeV $\tilde{\tau}$ mass ($5 < \tau(\tilde{\tau}) < 20$)
RPV	RPV : high-mass $e\mu$	$\sqrt{s} = 1.1$ TeV [1109.3098]	1.32 TeV $\tilde{\nu}_\tau$ mass ($\lambda_{112}^e = 0.10, \lambda_{312}^e = 0.05$)
	Bilinear RPV : 1 lep + j's + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [1109.6656]	760 GeV $\tilde{g} = \tilde{g}$ mass ($c_{\tau,bb} < 15$ mm)
	BC1 RPV : 4 lep + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-036]	1.77 TeV \tilde{g} mass
Other	Hypercolour scalar gluons : 4 jets, $m_s = m_{\tilde{g}}$	$\sqrt{s} = 34$ pb $^{-1}$, 7 TeV [1110.2693]	100-185 GeV sgluon mass (not excluded: $m_{\tilde{g}} = 140 \pm 3$ GeV)
	Spin dep. WIMP interaction : monojet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-064]	709 GeV M^* scale ($m_\chi < 100$ GeV, vector D5, Dirac χ)
	Spin indep. WIMP interaction : monojet + $E_{T,miss}$	$\sqrt{s} = 7$ TeV [ATLAS-CONF-2012-064]	548 GeV M^* scale ($m_\chi < 100$ GeV, tensor D9, Dirac χ)

$\int L dt = (0.03 - 4.8) \text{ fb}^{-1}$
 $\sqrt{s} = 7 \text{ TeV}$

ATLAS Preliminary



* Only a selection of the available mass limits on new states or phenomena shown

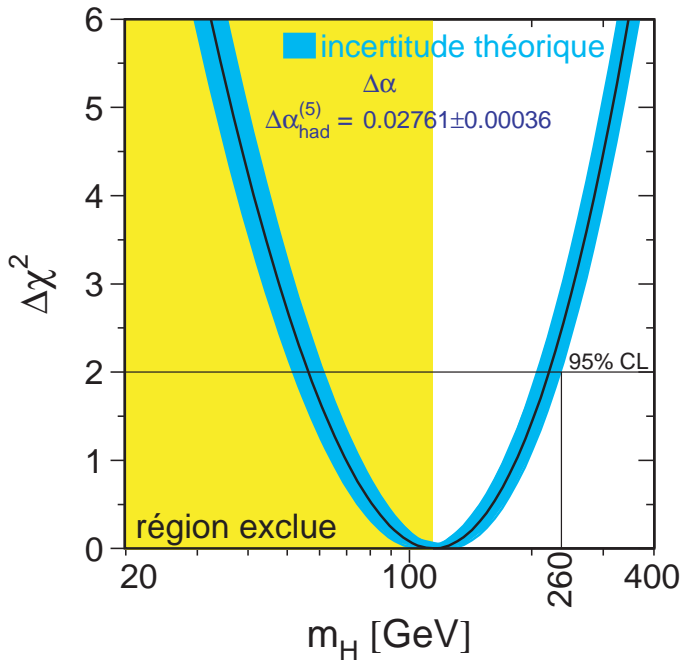


Remarks on the value ~ 125 GeV

- consistent with expectation from precision tests of the SM
- favors perturbative physics quartic coupling $\lambda = m_H^2/v^2 \simeq 1/8$
- window to new physics

If confirmed :

- compatible with supersymmetry
but appears fine-tuned in its minimal version [11]
early to draw a general conclusion before LHC13/14
e.g. an extra singlet or split families can alleviate the fine tuning [12]
- very important to measure its properties and couplings [16]
any deviation of its couplings to top, bottom and EW gauge bosons
implies new light states involved in the EWSB altering the fine-tuning



Fine-tuning in MSSM

Upper bound on the lightest scalar mass:

$$m_h^2 \lesssim m_Z^2 \cos^2 2\beta + \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{A_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{A_t^2}{12m_{\tilde{t}}^2} \right) \right] \lesssim (130 \text{ GeV})^2$$

$$m_h \simeq 126 \text{ GeV} \Rightarrow m_{\tilde{t}} \simeq 3 \text{ TeV or } A_t \simeq 3m_{\tilde{t}} \simeq 1.5 \text{ TeV}$$

\Rightarrow % to a few ‰ fine-tuning

$$\text{minimum of the potential: } m_Z^2 = 2 \frac{m_1^1 - m_2^2 \tan^2 \beta}{\tan^2 \beta - 1} \sim -2m_2^2 + \dots$$

$$\text{RG evolution: } m_2^2 = m_2^2(M_{\text{GUT}}) - \frac{3\lambda_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln \frac{M_{\text{GUT}}}{m_{\tilde{t}}} + \dots$$

$$\sim m_2^2(M_{\text{GUT}}) - \mathcal{O}(1)m_{\tilde{t}}^2 + \dots \quad [9]$$

MSSM with dim-5 and 6 operators

I.A.-Dudas-Ghilencea-Tziveloglou '08, '09, '10

parametrize new physics above MSSM by higher-dim effective operators

relevant super potential operators of dimension-5:

$$\mathcal{L}^{(5)} = \frac{1}{M} \int d^2\theta (\eta_1 + \eta_2 S) (H_1 H_2)^2$$

η_1 : generated for instance by a singlet

$$W = \lambda \sigma H_1 H_2 + M \sigma^2 \quad \rightarrow \quad W_{\text{eff}} = \frac{\lambda^2}{M} (H_1 H_2)^2$$

Strumia '99 ; Brignole-Casas-Espinosa-Navarro '03

Dine-Seiberg-Thomas '07

η_1 : corresponding soft breaking term spurion $S \equiv m_S \theta^2$

Physical consequences of MSSM₅: Scalar potential

$$\mathcal{V} = m_1^2 |h_1|^2 + m_2^2 |h_2|^2 + B\mu(h_1 h_2 + \text{h.c.}) + \frac{g_2^2 + g_Y^2}{8} (|h_1|^2 - |h_2|^2)^2 \\ + (|h_1|^2 + |h_2|^2) (\eta_1 h_1 h_2 + \text{h.c.}) + \frac{1}{2} [\eta_2 (h_1 h_2)^2 + \text{h.c.}]$$

- $\eta_{1,2} \Rightarrow$ quartic terms along the D-flat direction $|h_1| = |h_2|$
- potential stability $\Rightarrow \eta_2 \geq 4|\eta_1|$

requiring η -corrections to be smaller than MSSM mass matrix elements \Rightarrow

only η_2 can change the tree-level bound $m_h \leq m_Z$ but marginally

Relevance of dim-6 operators

Relaxing the condition on potential positivity: guaranteed by dim-6 ops

only one dim-6 along the D-flat direction induced by dim-5: $\propto \eta_1^2$

$$W = \eta_1 (H_1 H_2)^2 \longrightarrow V = \left| \frac{\partial W}{\partial H_i} \right|^2 \sim \eta_1^2 |H_1 H_2|^2 (|H_1|^2 + |H_2|^2)$$

- tree-level mass can increase significantly
- bigger parameter space for LSP being dark matter

Bernal-Blum-Nir-Losada '09

MSSM Higgs with dim-6 operators

dim-6 operators can have an independent scale from dim-5

Classification of all dim-6 contributing to the scalar potential

(without SUSY) \Rightarrow

large $\tan \beta$ expansion: $\delta_6 m_h^2 = f v^2 + \dots$

constant receiving contributions from several operators

$f \sim f_0 \times (\mu^2/M^2, m_S^2/M^2, \mu m_S/M^2, v^2/M^2)$

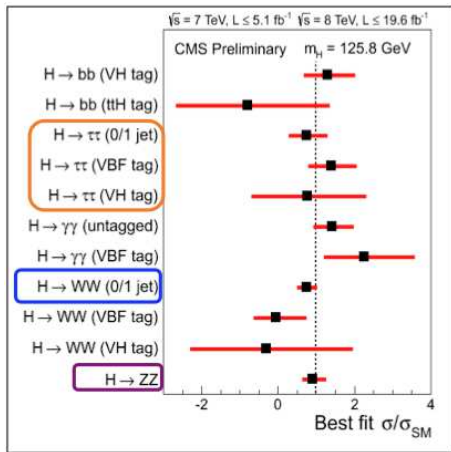
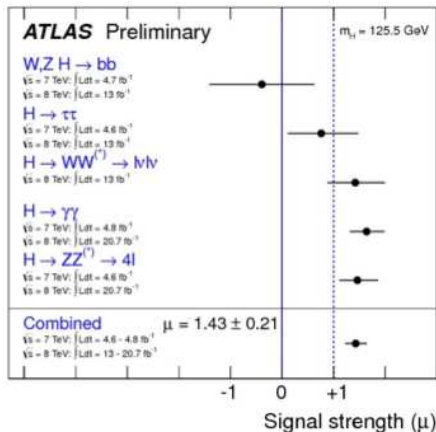
$m_S = 1 \text{ TeV}, M = 10 \text{ TeV}, f_0 \sim 1 - 2.5$ for each operator

$\Rightarrow m_h \simeq 103 - 119 \text{ GeV}$

\Rightarrow MSSM with dim-5 and dim-6 operators:

possible resolution of the MSSM fine-tuning problem [9]

Couplings of the new boson vs SM

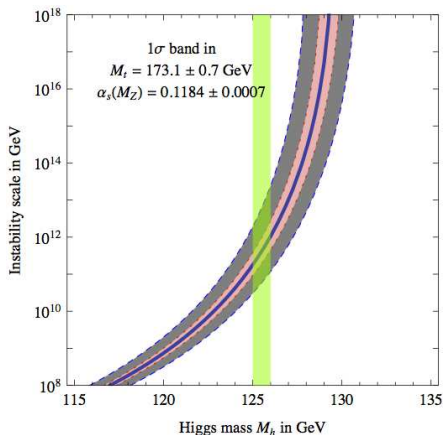
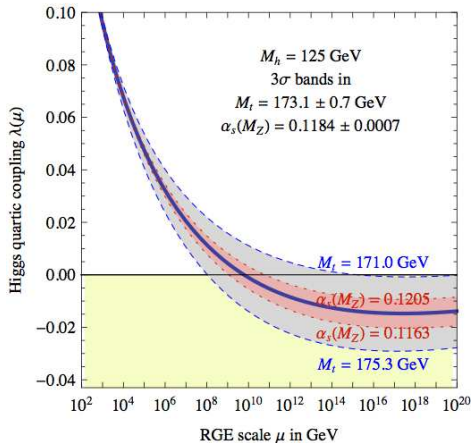


exclusion : spin 2 and pseudoscalar at 95% CL

Agreement with Standard Model expectation at $\sim 2\sigma$

Can the SM be valid at high energies?

Degrassi-Di Vita-Elias Miró-Espinosa-Giudice-Isidori-Strumia '12



Instability of the SM Higgs potential \Rightarrow metastability of the EW vacuum

If the weak scale is tuned \Rightarrow split supersymmetry is a possibility

Arkani Hamed-Dimopoulos '04, Giudice-Romanino '04

- natural splitting: gauginos, higgsinos carry R-symmetry, scalars do not
- main good properties of SUSY are maintained
 - gauge coupling unification and dark matter candidate
- also no dangerous FCNC, CP violation, ...
- experimentally allowed Higgs mass \Rightarrow 'moderate' split

$m_S \sim \text{few} - \text{thousands TeV}$

gauginos: a loop factor lighter than scalars ($\sim m_{3/2}$)

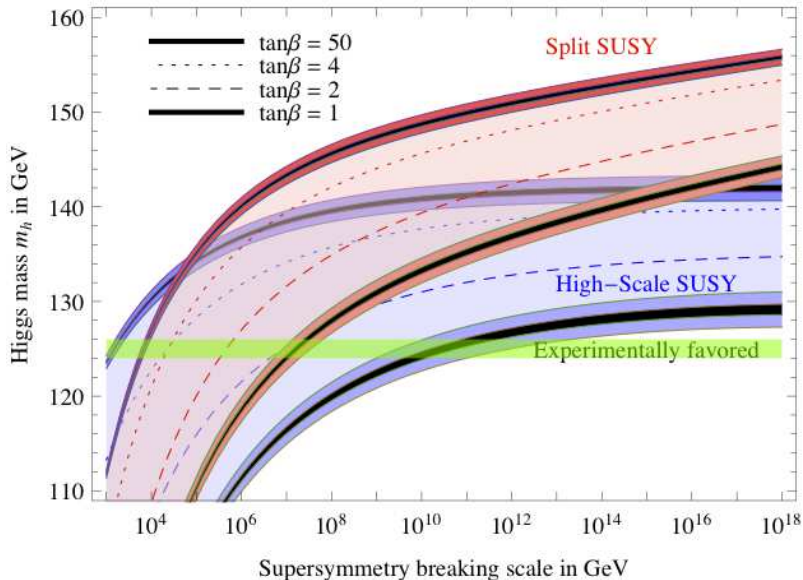
- natural string framework: intersecting (or magnetized) branes

IA-Dimopoulos '04

D-brane stacks are supersymmetric with massless gauginos

intersections have chiral fermions with broken SUSY & massive scalars

Predicted range for the Higgs mass

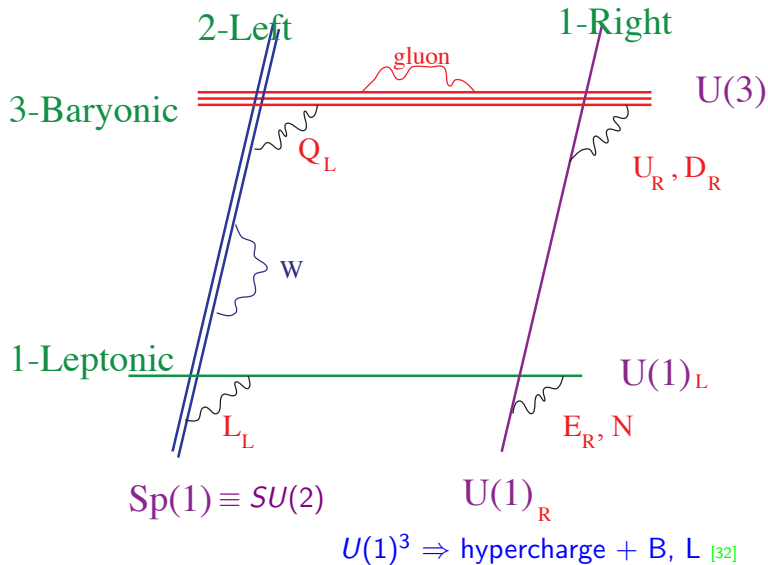


An extra $U(1)$ can also cure the instability problem

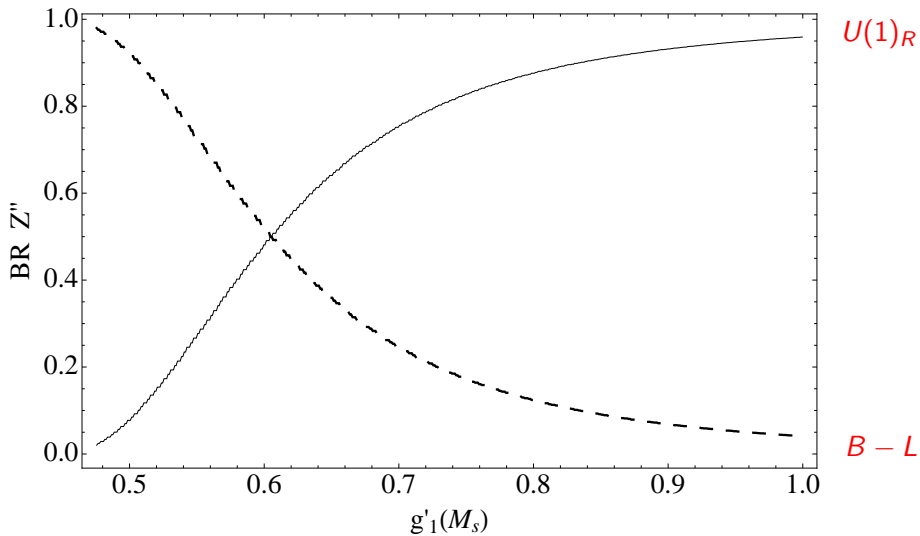
Anchordoqui-IA-Goldberg-Huang-Lüst-Taylor-Vlcek '12

- B anomalous and superheavy
- $B - L$ massless at the string scale (no associated 6d anomaly)
but broken at TeV by a scalar VEV with the quantum numbers of N_R
- L -violation from higher-dim operators suppressed by the string scale
- $U(3)$ unification, Y combination \Rightarrow 2 parameters: 1 coupling + $m_{Z''}$
- perturbativity $\Rightarrow 0.5 \lesssim g_{U(1)_R} \lesssim 1$
- present LHC limits: $m_{Z''} \gtrsim 3 - 4$ TeV (for $Z'' \simeq B - L$ or $U(1)_R$)
- interesting LHC phenomenology and cosmology [22]

Standard Model on D-branes : SM⁺⁺



- Rotation of $U(1)$'s from the string to low energy basis Z, Z', Z'' :
completely fixed in terms of the couplings
 - Decoupling of anomalous $Z' \simeq B$
 - Z'' linear combination of $B - L$ and $U(1)_R$
- Recent cosmological observations indicate extra relativistic component
dark radiation parametrized by an effective neutrino number close to 4
→ use the 3 ν_R 's interacting with SM fermions via Z''
data: their decoupling during the quark-hadron transition
⇒ $3.5 \lesssim M_{Z''} \lesssim 7 \text{ TeV}$ (within LHC14 discovery potential)



Scalar potential:

$$V(H, H'') = \mu^2 |H|^2 + \mu'^2 |H''|^2 + \lambda_1 |H|^4 + \lambda_2 |H''|^4 + \lambda_3 |H|^2 |H''|^2$$

5 parameters $\Rightarrow v, m_h, v'', m_{h''}$ + a scalar mixing angle α

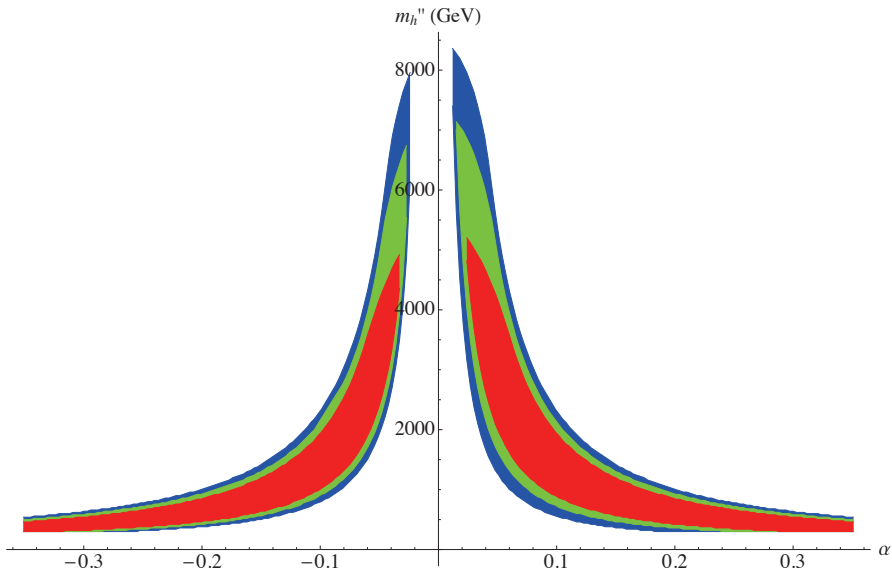
\Rightarrow 3 free parameters : $m_{h''}, \alpha, v'' \leftrightarrow M_{Z''}$

Stability conditions: $\lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_1 \lambda_2 > \frac{1}{4} \lambda_3^2$

RGE analysis up to $M_s \Rightarrow$ stability is possible in SM⁺⁺

for $0.02 \lesssim |\alpha| \lesssim 0.35$ and $500 \text{ GeV} \lesssim m_{h''} \lesssim 5 \text{ TeV}$

$$M_{Z''} = 4.5 \text{ TeV}; \quad M_S = 10^{14}, 10^{16}, 10^{19} \text{ GeV}$$



Alternative answer: Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity \Rightarrow extra dimensions: large flat or warped
- low string scale \Rightarrow low scale gravity, ultra weak string coupling

$M_s \sim 1 \text{ TeV} \Rightarrow$ volume $R_{\perp}^n = 10^{32} l_s^n$ [34] ($R_{\perp} \sim .1 - 10^{-13} \text{ mm}$ for $n = 2 - 6$)

- spectacular model independent predictions
- radical change of high energy physics at the TeV scale

Moreover no little hierarchy problem:

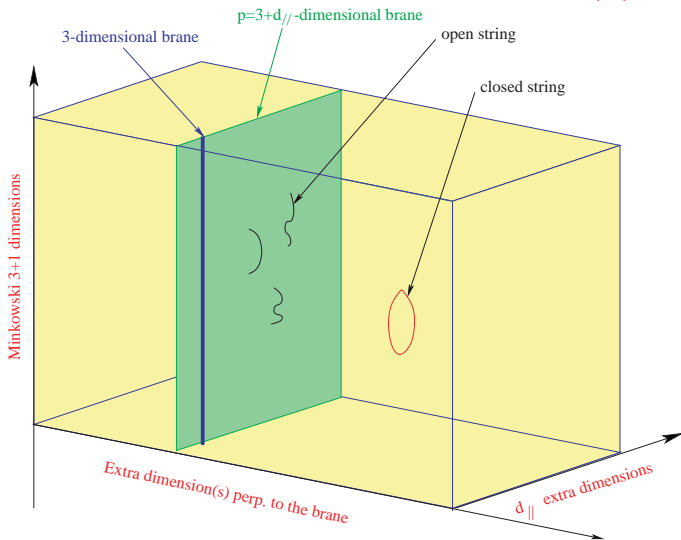
radiative electroweak symmetry breaking with no logs

$\Lambda \sim$ a few TeV and $m_H^2 =$ a loop factor $\times \Lambda^2$ [28]

But unification has to be probably dropped

New Dark Matter candidates e.g. in the extra dims

- 2 types of compact extra dimensions:
- parallel (d_{\parallel}): $\lesssim 10^{-16}$ cm (TeV)
 - transverse (\perp): $\lesssim 0.1$ mm (meV) [??]



Origin of EW symmetry breaking?

possible answer: radiative breaking

I.A.-Benakli-Quiros '00

$$V = \mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

$\mu^2 = 0$ at tree but becomes < 0 at one loop

non-susy vacuum

simplest case: one scalar doublet from the same brane

\Rightarrow tree-level V same as susy: $\lambda = \frac{1}{8}(g_2^2 + g'^2)$

D-terms

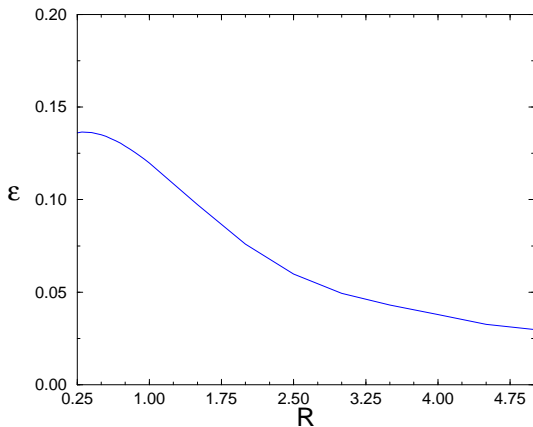
$\mu^2 = -g^2 \epsilon^2 M_5^2 \leftarrow$ effective UV cutoff

$$\epsilon^2(R) = \frac{R^3}{2\pi^2} \int_0^\infty dl l^{3/2} \frac{\theta_2^4}{16l^4 \eta^{12}} \left(il + \frac{1}{2} \right) \sum_n n^2 e^{-2\pi n^2 R^2 l}$$

UV \swarrow IR \nearrow

1 \nwarrow 1 \nearrow

e^{-πl}



$R \rightarrow 0$: $\varepsilon(R) \simeq 0.14$ large transverse dim $R_{\perp} = l_s^2/R \rightarrow \infty$

$R \rightarrow \infty$: $\varepsilon(R)M_s \sim \varepsilon_{\infty}/R$ $\varepsilon_{\infty} \simeq 0.008$ UV cutoff: $M_s \rightarrow 1/R$

Higgs scalar = component of a higher dimensional gauge field

$\Rightarrow \varepsilon_{\infty}$ calculable in the effective field theory

Quartic coupling \Rightarrow mass prediction:

- tree level : $M_H = M_Z$

- low-energy SM radiative corrections (from top quark) : $M_H \sim 120$ GeV

Casas-Espinosa-Quiros-Riotto, Carena-Espinosa-Quiros-Wagner '95

Increasing $\lambda \rightarrow g^2/4 \sim 1/8 \Rightarrow M_H \simeq v/2 = 125$ GeV

Also M_5 or $1/R \sim$ a few or several TeV

Accelerator signatures: 4 different scales

- Gravitational radiation in the bulk \Rightarrow missing energy

present LHC bounds: $M_* \gtrsim 2.5 - 4$ TeV

- Massive string vibrations \Rightarrow e.g. resonances in dijet distribution

$$M_j^2 = M_0^2 + M_s^2 j \quad ; \quad \text{maximal spin : } j + 1$$

higher spin excitations of quarks and gluons with strong interactions

present LHC limits: $M_s \gtrsim 4.5$ TeV

- Large TeV dimensions \Rightarrow KK resonances of SM gauge bosons I.A. '90

$$M_k^2 = M_0^2 + k^2/R^2 \quad ; \quad k = \pm 1, \pm 2, \dots$$

experimental limits: $R^{-1} \gtrsim 0.5 - 4$ TeV (UED - localized fermions)

- extra $U(1)$'s and anomaly induced terms

masses suppressed by a loop factor from M_s [32]

Extra $U(1)$'s and anomaly induced terms

masses suppressed by a loop factor

usually associated to known global symmetries of the SM

(anomalous or not) such as (combinations of)

Baryon and Lepton number, or PQ symmetry

Two kinds of massive $U(1)$'s:

I.A.-Kiritsis-Rizos '02

- 4d anomalous $U(1)$'s: $M_A \simeq g_A M_s$

- 4d non-anomalous $U(1)$'s: (but masses related to 6d anomalies)

$$M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d) \text{ internal space} \Rightarrow M_{NA} \geq M_A$$

or massless in the absence of such anomalies [21]


- B and L become massive due to anomalies

Green-Schwarz terms

- the global symmetries remain in perturbation

- Baryon number \Rightarrow proton stability

- Lepton number \Rightarrow protect small neutrino masses

no Lepton number $\Rightarrow \frac{1}{M_s} LLHH \rightarrow$ Majorana mass: $\frac{\langle H \rangle^2}{M_s} LL$


- $B, L \Rightarrow$ extra Z' 's

with possible leptophobic couplings leading to CDF-type Wjj events

$Z' \simeq B$ lighter than 4d anomaly free $Z'' \simeq B - L$

More general framework: large number of species

N particle species \Rightarrow lower quantum gravity scale : $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10

derivation from: black hole evaporation or quantum information storage

$$M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32} \text{ particle species !}$$

2 ways to realize it lowering the string scale

① Large extra dimensions SM on D-branes [26]

$N = R_{\perp}^n l_s^n$: number of KK modes up to energies of order $M_* \simeq M_s$

② Effective number of string modes contributing to the BH bound

$N = \frac{1}{g_s^2}$ with $g_s \simeq 10^{-16}$ SM on NS5-branes

I.A.-Pioline '99, I.A.-Dimopoulos-Giveon '01

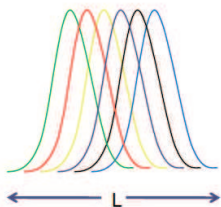
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derivation from: black hole evaporation or quantum information storage

Pixel of size L containing N species storing information:



localization energy $E \gtrsim N/L \rightarrow$

Schwarzschild radius $R_s = N/(LM_p^2)$

no collapse to a black hole : $L \gtrsim R_s \Rightarrow L \gtrsim \sqrt{N}/M_p = 1/M_*$

$M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32}$ particle species !

What is LST ? **Decouple gravity from NS5-branes**

Analogy from D3-branes : decouple gravity $\Rightarrow M_s \rightarrow \infty$, g_s fixed
 \rightarrow (conformal) Field Theory (CFT)

simplest case: 4d $\mathcal{N} = 4$ super Yang Mills **$SU(N)$**

parameters: number of branes N , gauge coupling g_{YM}

NS-5 branes: M_s finite, $g_s \rightarrow 0 \rightarrow$ (little) String Theory without gravity

simplest case: 6d LST (**chiral IIA or non-chiral IIB**)

massless sector: 6d $SU(N)$ of tensors (IIA) or vectors (IIB)

at a non-trivial fixed point

parameters: number of branes N , string scale M_s

How to study LST ? Using gauge/gravity duality

Gravity background : near horizon geometry (holography) Maldacena '98

Analogy from D3-branes : $AdS_5 \times S^5$

parameters: AdS radius $r_{AdS} M_s$, $g_s \leftrightarrow N, g_{YM}$

supergravity validity: $r_{AdS} M_s \gg 1$, $g_s \ll 1 \Rightarrow$ large N , $g_{YM}^2 N$

\rightarrow model independent part : AdS_5

NS-5 branes : $(\mathcal{M}_6 \otimes R_+) \times SU(2) \equiv S^3$

\uparrow
linear dilaton background in 7d flat string-frame metric $\Phi = -\alpha|y|$

Aharony-Berkooz-Kutasov-Seiberg '98

parameters: M_s , α (or S^3 radius) $\leftrightarrow N$

sugra validity: small $\alpha \Rightarrow$ large N

compactify to $d = 4$ ($\mathcal{M}_6 \rightarrow \mathcal{M}_4$) $\Rightarrow g_{YM} \sim$ 2d volume

\rightarrow model independent part : linear dilaton

Put gravity back **but weakly coupled**

“cut” the space of the extra dimension \Rightarrow gravity on the brane

Toy 5d bulk model

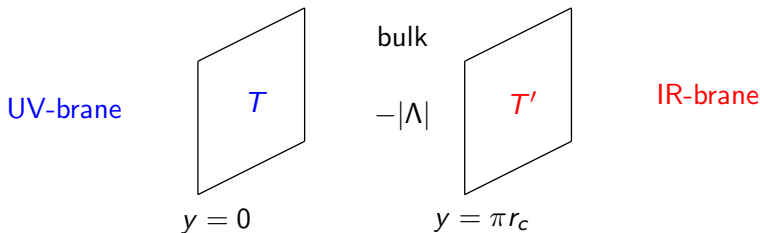
$$S_{bulk} = \int d^4x \int_0^{r_c} dy \sqrt{-g} e^{-\Phi} (M_5^3 R + M_5^3 (\nabla\Phi)^2 - \Lambda)$$

$$S_{vis(hid)} = \int d^4x \sqrt{-g} (e^{-\Phi}) (L_{SM(hid)} - T_{vis(hid)})$$

Tuning conditions: $T_{vis} = -T_{hid} \leftrightarrow \Lambda < 0$ [40]

Constant dilaton and AdS metric : Randal Sundrum model

spacetime = slice of AdS_5 : $ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$ $k^2 \sim \Lambda/M_5^3$



• exponential hierarchy: $M_W = M_P e^{-2kr_c}$ $M_P^2 \sim M_5^3/k$ $M_5 \sim M_{GUT}$

• 4d gravity localized on the UV-brane, but KK gravitons on the IR

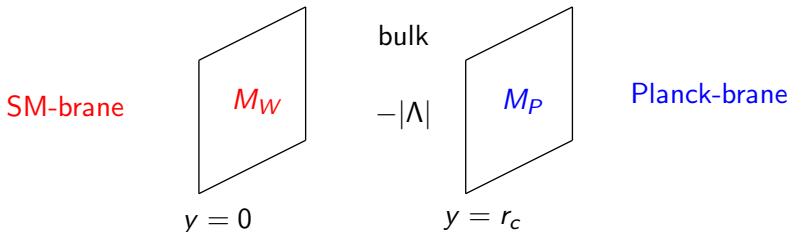
$$m_n = c_n k e^{-2kr_c} \sim \text{TeV} \quad c_n \simeq (n + 1/4) \text{ for large } n$$

\Rightarrow spin-2 TeV resonances in di-lepton or di-jet channels

dilaton $\Phi = -\alpha|y|$ and flat metric \Rightarrow

$$g_s^2 = e^{-\alpha|y|} ; ds^2 = e^{\frac{2}{3}\alpha|y|} (\eta_{\mu\nu} dx^\mu dx^\nu + dy^2) \leftarrow \text{Einstein frame}$$

$z \sim e^{\alpha y/3} \Rightarrow$ polynomial warp factor + log varying dilaton



- exponential hierarchy: $g_s^2 = e^{-\alpha|y|}$ $M_P^2 \sim \frac{M_5^3}{\alpha} e^{\alpha r_c}$ $\alpha \equiv k_{RS}$
- 4d graviton flat, KK gravitons localized near SM

LST KK graviton phenomenology

- KK spectrum : $m_n^2 = \left(\frac{n\pi}{r_c}\right)^2 + \frac{\alpha^2}{4}$; $n = 1, 2, \dots$

⇒ mass gap + dense KK modes $\alpha \sim 1 \text{ TeV}$ $r_c^{-1} \sim 30 \text{ GeV}$

- couplings : $\frac{1}{\Lambda_n} \sim \frac{1}{(\alpha r_c) M_5}$

⇒ extra suppression by a factor $(\alpha r_c) \simeq 30$

- width : $1/(\alpha r_c)^2$ suppression $\sim 1 \text{ GeV}$

⇒ narrow resonant peaks in di-lepton or di-jet channels

- extrapolates between RS and flat extra dims ($n = 1$)

⇒ distinct experimental signals

Similar to RS using the dilaton as the Goldberger-Wise scalar

add dilaton boundary potentials \Rightarrow

radion stabilization with the desired hierarchy

Radion phenomenology different from RS:

- mass spectrum: similar to the graviton KK modes
 - with possible lower parametrically mass gap
- new radion couplings to SM fields besides to the trace of $T_{\mu\nu}$
- larger coupling to the radion 0-mode relative to KK excitations
- Higgs-radion mixing \Rightarrow
 - branching fraction to $\gamma\gamma$ can be significantly enhanced

Conclusions

- Confirmation of the EWSB scalar at the LHC:
important milestone of the LHC research program
- Precise measurement of its couplings is of primary importance
- Hint on the origin of mass hierarchy and of BSM physics
 - natural or unnatural SUSY?
 - low string scale in some realization?
 - something new and unexpected?
- LHC enters a new era with possible new discoveries