

Holographic Type II Goldstone Bosons

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I. Amado, D. Arean, A. Jimenez-Alba, K.L., L. Melgar, I. Salazar-Landea
[arXiv:1302.5641, arXiv:1307.8100]

Crete CTP, 07-10-2013

Outline

- Goldstone theorems
- Field theoretical model
- Holographic model(s)
- Landau criterion
- Summary

Goldstone Theorems

- Spontaneously broken continuous symmetry

$$\lim_{k \rightarrow 0} \omega(k) = 0$$

- At least one mode
- No constraint on power: $\omega(k) \propto k^n$
- Lorentz symmetry:
 - $\omega(k) = ck$
 - One mode for every broken generator

Goldstone Theorems

- No Lorentz symmetry
 - State: temperature T , density μ
 - Principally: non-relativistic, Lifshitz, ...
- Classification : [Nielsen-Chadha '74]
 - Type I : $\omega \propto k^{2n+1}$
 - Type II: $\omega \propto k^{2n}$

Goldstone Theorems

- Chadha-Nielsen

$$N_I + 2N_{II} \geq N_{BG}$$

- Brauner-Watanabe-Murayama

$$\langle [Q_a, Q_b] \rangle = B_{ab}$$
$$N_I + N_{II} = N_{BG} - \frac{1}{2} \text{rank}(B_{ab})$$

- Brauner-Murayama-Watanabe,
Nicolis-Piazza, Kapustin
("massive" Goldstone)

$$\omega = q\mu$$

Field Theory Model

T. Schafer, D. T. Son, M. A. Stephanov, D. Toublan and J. J. M. Verbaarschot, [hep-ph/0108210]

V. A. Miransky and I. A. Shovkovy, [hep-ph/0108178]

$$\mathcal{L} = (\partial_0 - i\mu)\Phi^\dagger (\partial_0 + i\mu)\Phi - \vec{\partial}\Phi^\dagger \vec{\partial}\Phi - M^2\Phi^\dagger\Phi - \lambda^4(\Phi^\dagger\Phi)^2$$

Doublet of U(2) $\phi = (\phi_1, \phi_2)^T$ $\phi = (0, v)^T$

μ inside overall U(1)

$$\omega_1^2 = \frac{\mu^2 - M^2}{3\mu^2 - M^2} p^2 + O(p^4),$$

$$\omega_2^2 = 6\mu^2 - 2M^2 + O(p^2),$$

$$\omega_3^2 = p^2 - 2\mu\omega_3, \quad \omega_3 \approx \frac{p^2}{2\mu}$$

$$\omega_4^2 = p^2 + 2\mu\omega_4. \quad \omega_4 = 2\mu$$

→ Type I Goldstone

→ Type II Goldstone

→ “Massive Goldstone”

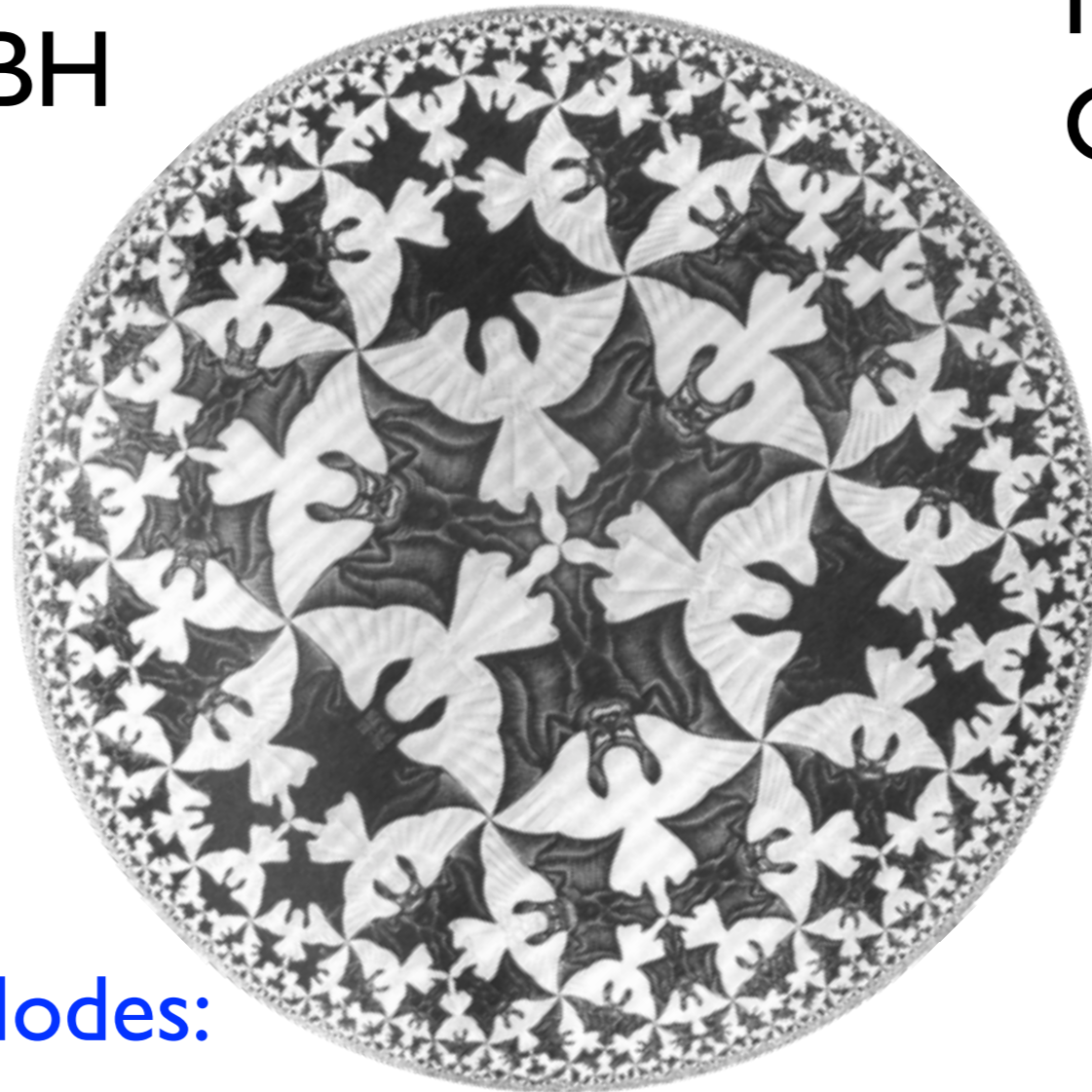
Holography

- Global on boundary = local in Bulk
- $U(2)$ gauge fields + scalar in doublet
- gauged model,
- Chemical potential only in $U(1)$
- Holographic Goldstone modes = Quasinormal Modes
- Simplest: $U(2)$ generalization of HHH
[Hartnoll-Herzog-Horowitz]
- Decoupling limit

Holography

3D AdS
Schwarzschild BH

Maxwell +
Charged Scalar

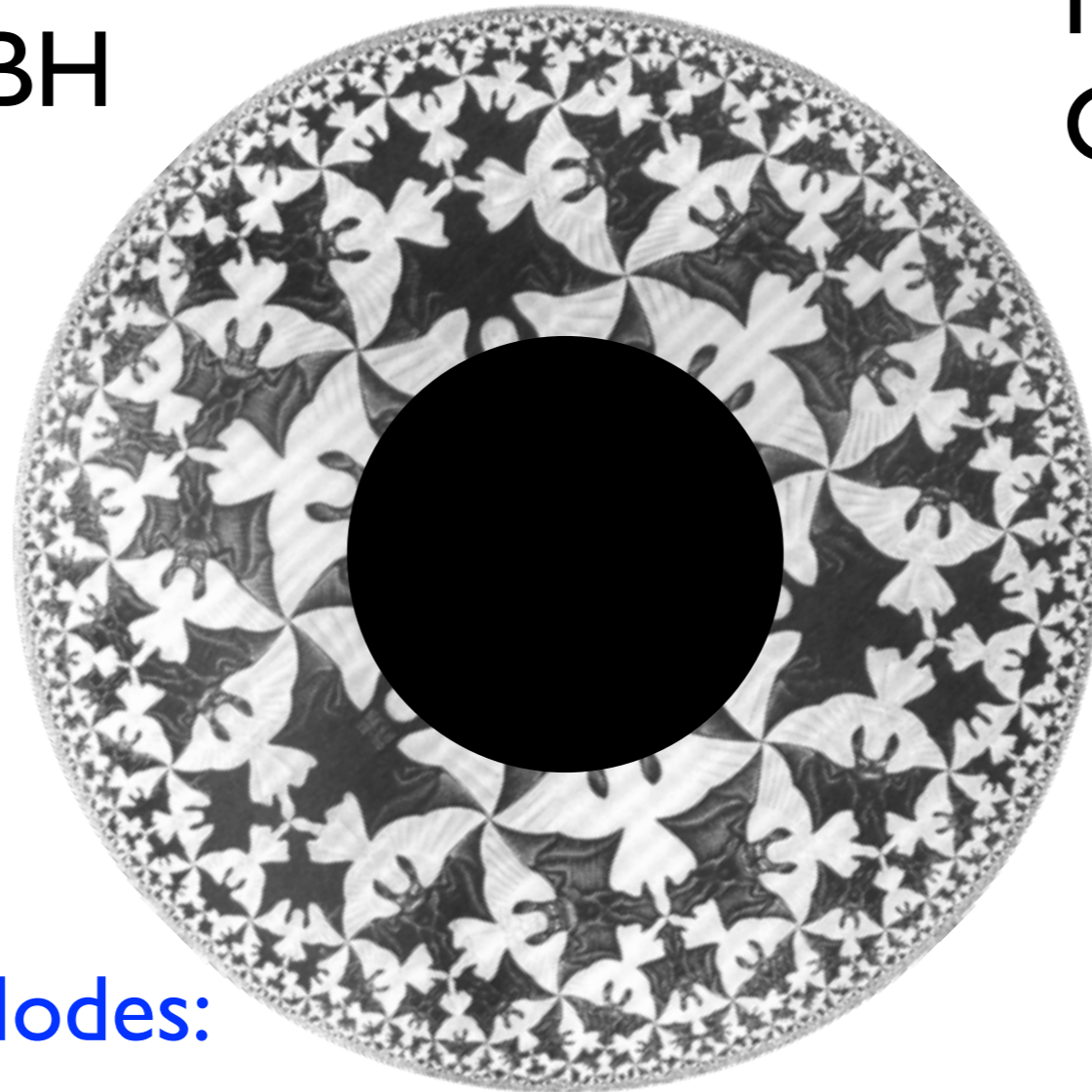


Quasinormal Modes:

Holography

3D AdS
Schwarzschild BH

Maxwell +
Charged Scalar

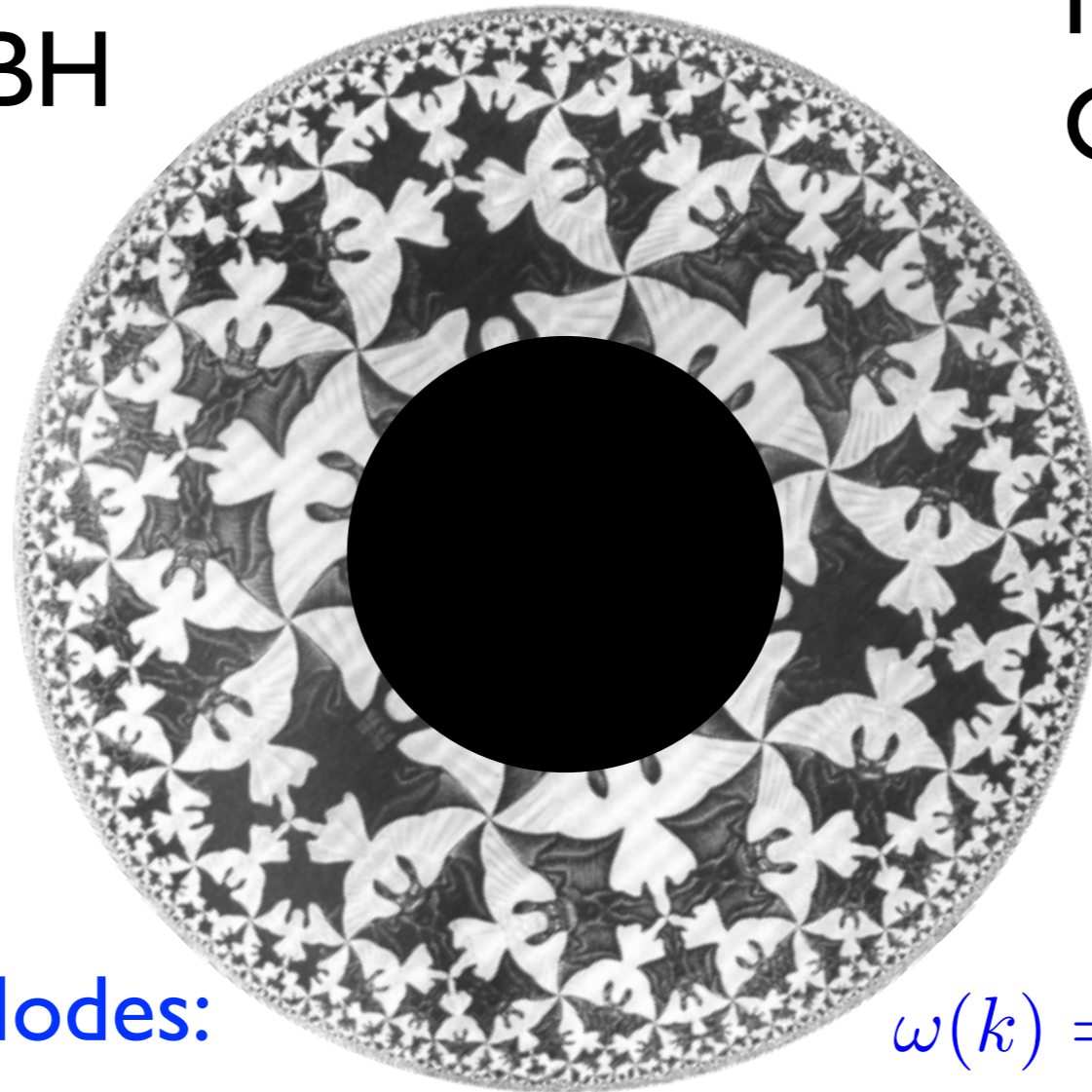


Quasinormal Modes:

Holography

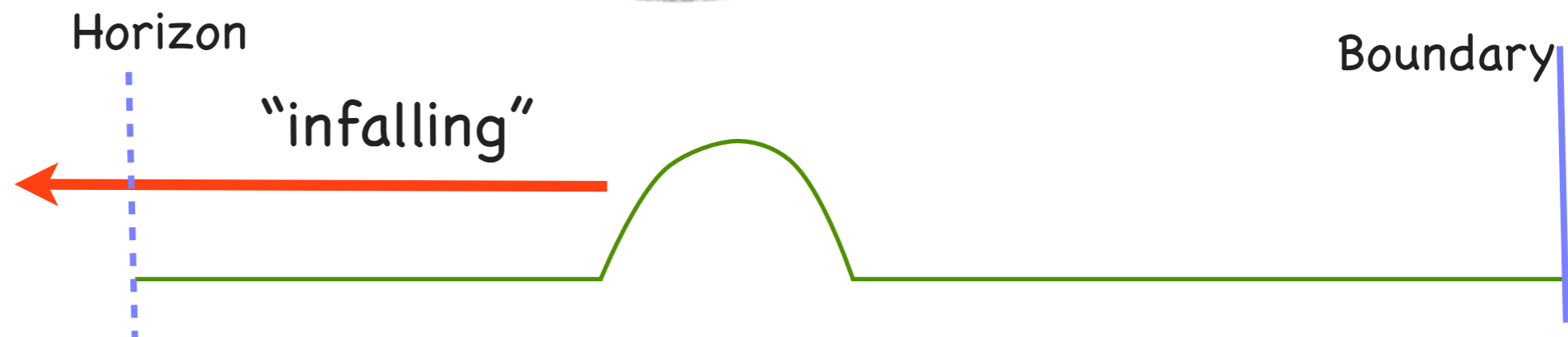
3D AdS
Schwarzschild BH

Maxwell +
Charged Scalar



Quasinormal Modes:

$$\omega(k) = \pm\Omega(k) - i\Gamma(k)$$



Holography

U(2) decomposition:

$$\sigma_- = \frac{1}{2}(\sigma_0 - \sigma_3) = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

broken:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_2 = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}$$

unbroken:

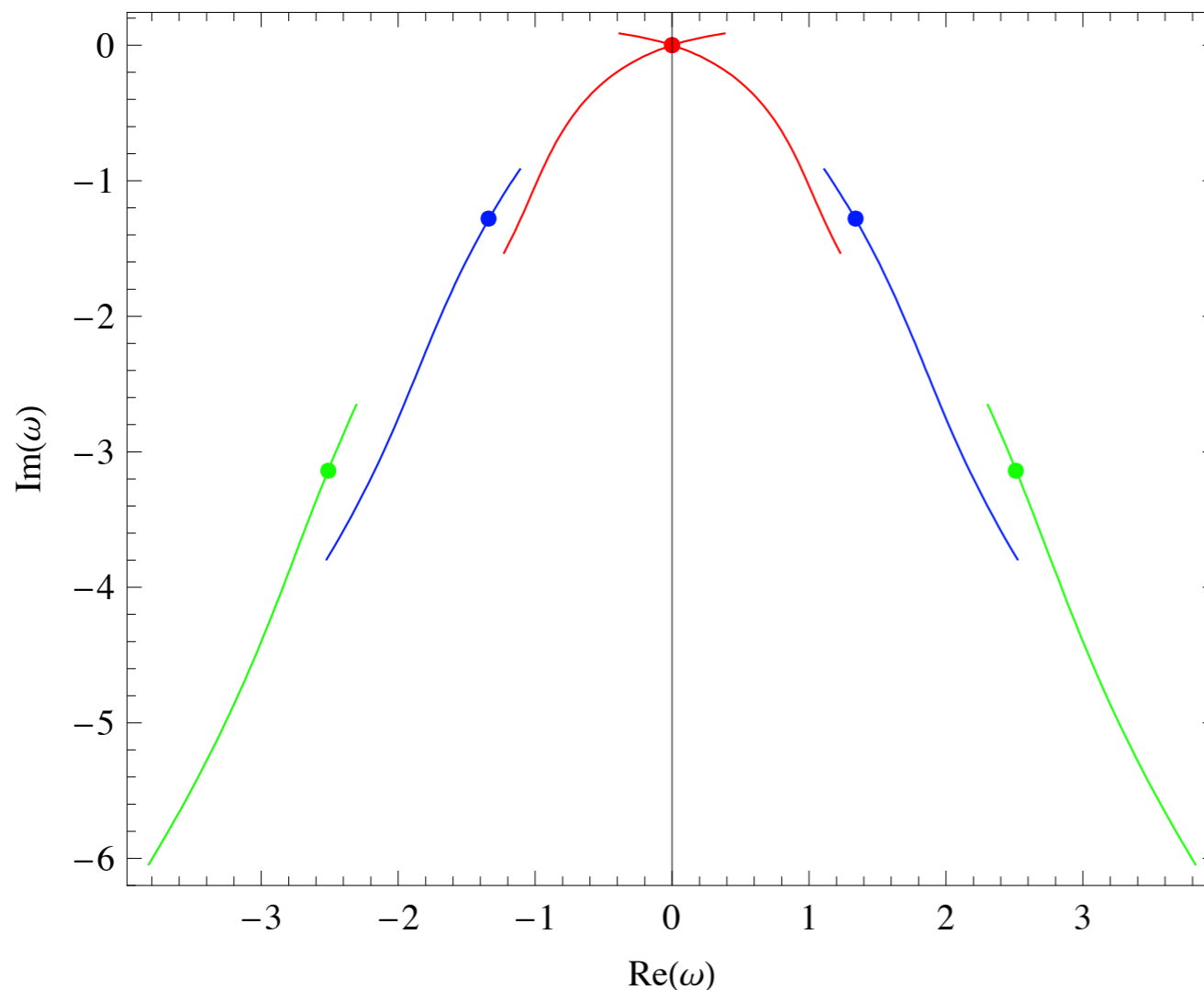
$$\sigma_+ = \frac{1}{2}(\sigma_0 + \sigma_3) = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Holography

Normal phase (high T)

I. Amado, M. Kaminski and K. L., [arXiv:0903.2209 [hep-th]]

O_2 Theory



Scalar QNMs
doubly degenerate

There are also 4 Diffusion modes $\omega = -iDk^2$ $D = \frac{3}{4\pi T}$

Holography

$$\Psi'' + \left(\frac{f'}{f} + \frac{2}{\rho} \right) \Psi' + \frac{\chi^2}{f^2} \Psi - \frac{m^2}{f} \Psi = 0,$$

$$\chi'' + \frac{2}{\rho} \chi' - \frac{2\Psi^2}{f} \chi = 0,$$

$$\xi'' + \frac{2}{\rho} \xi' = 0,$$

$$\chi = \bar{\mu}_\chi - \frac{\bar{n}_\chi}{\rho} + O\left(\frac{1}{\rho^2}\right)$$

$$\xi = \bar{\mu}_\xi - \frac{\bar{n}_\xi}{\rho} + O\left(\frac{1}{\rho^2}\right)$$

$$\Psi = \frac{\psi_1}{\rho} + \frac{\psi_2}{\rho^2} + O\left(\frac{1}{\rho^3}\right)$$

b.c.s: $\bar{\mu}_\chi = \bar{\mu}_\xi = \bar{\mu}.$

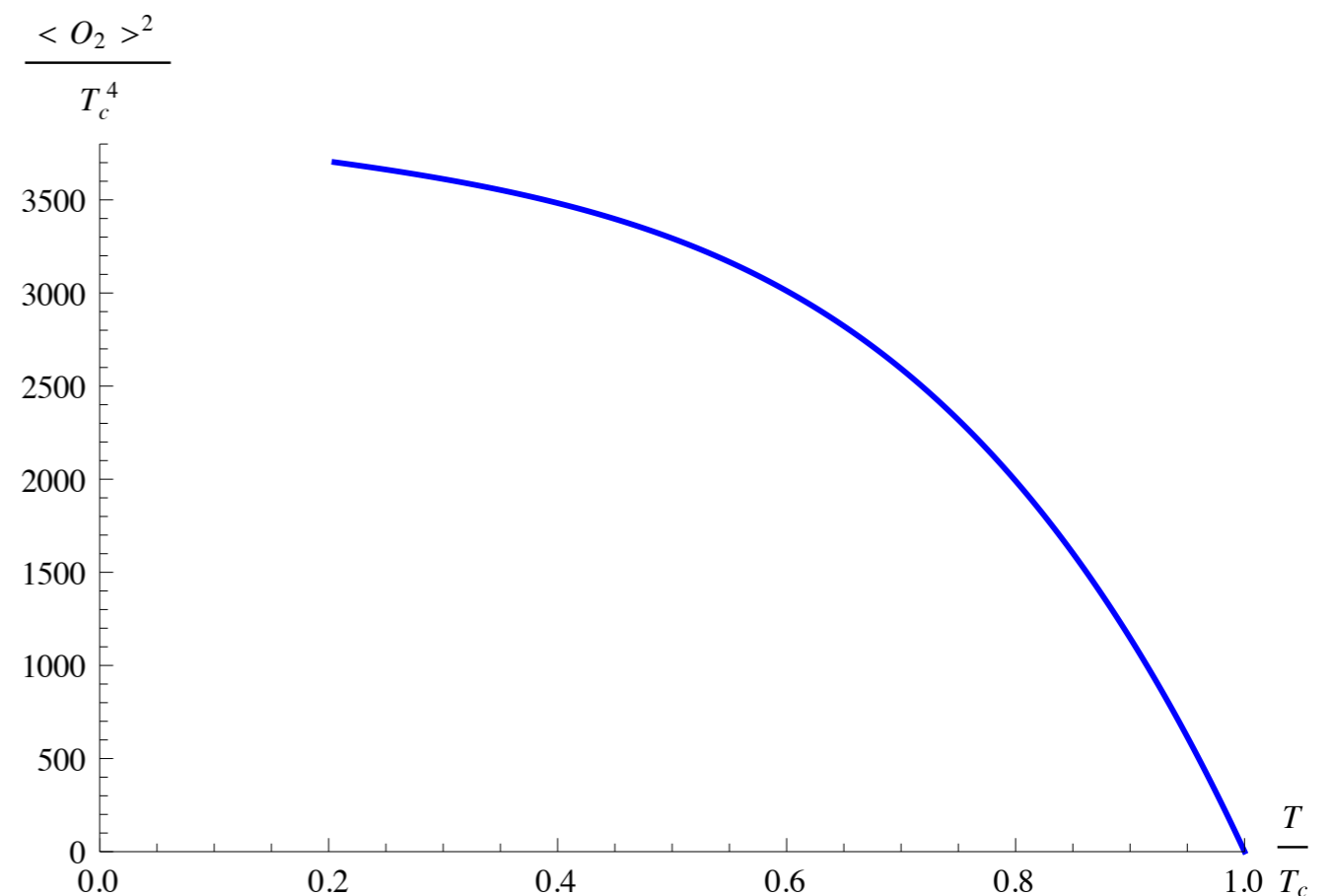
$$\psi_1 = 0.$$

$$\psi_2 \propto \langle O \rangle$$

Scalar

σ_- Sector

σ_+ Sector



Holography

$$\Psi'' + \left(\frac{f'}{f} + \frac{2}{\rho} \right) \Psi' + \frac{\chi^2}{f^2} \Psi - \frac{m^2}{f} \Psi = 0,$$

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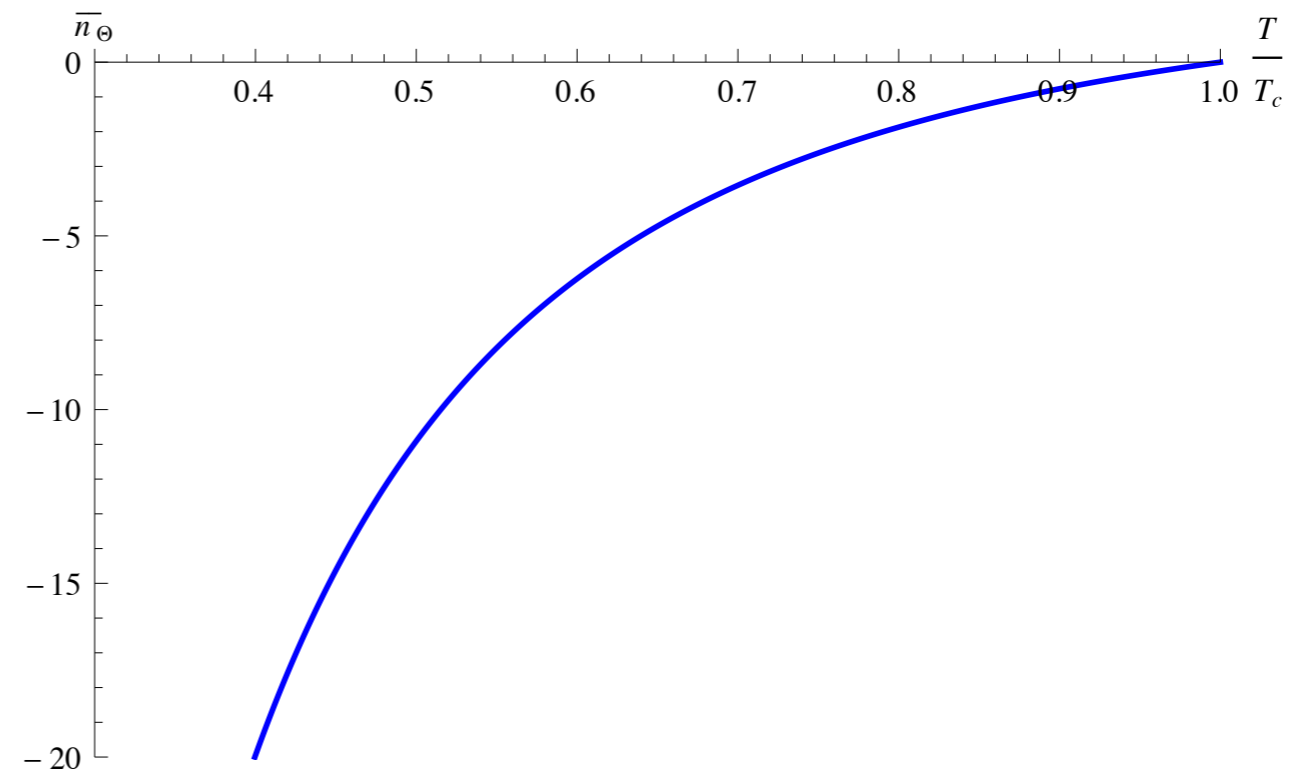
Scalar

σ_- Sector

σ_+ Sector

Charge in σ_3 Sector:

(Charge in σ_0 sector always >0)



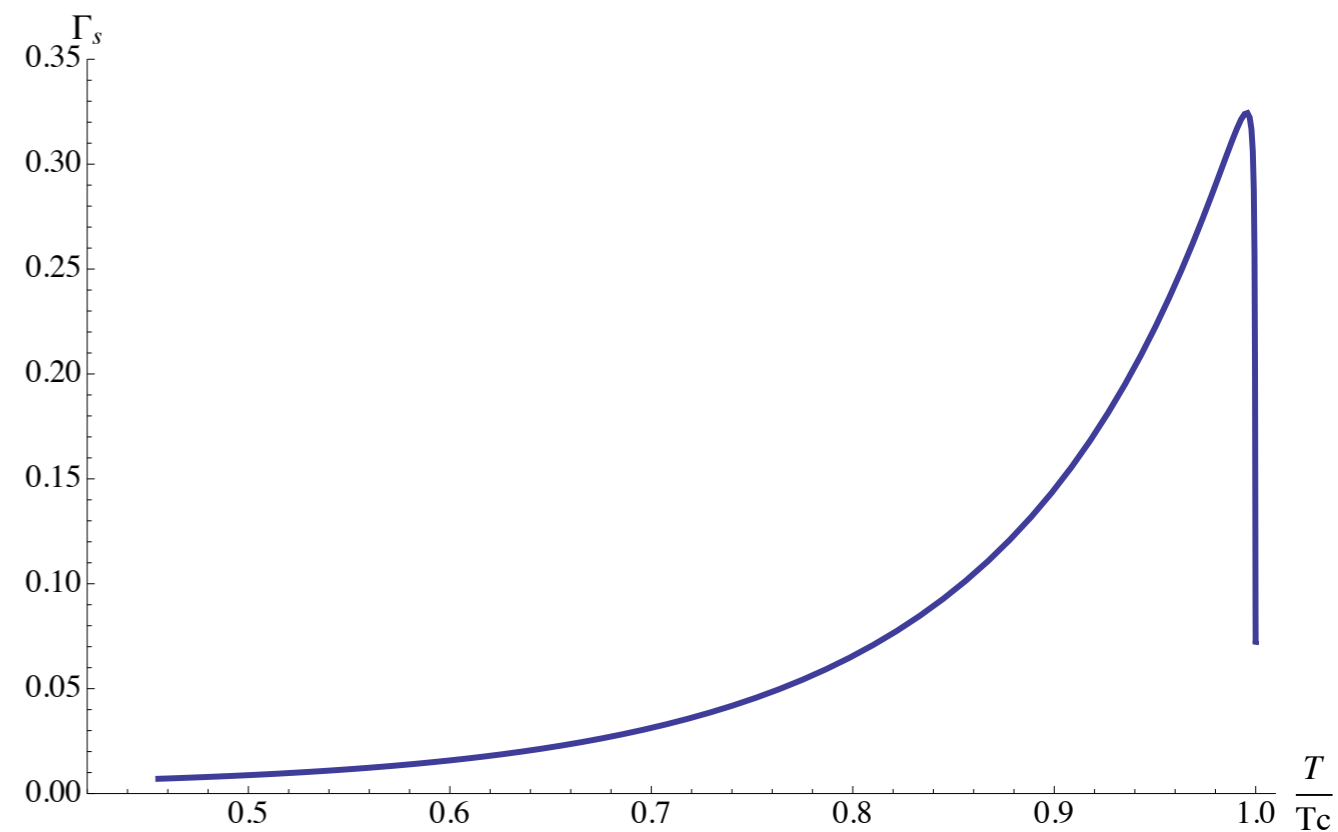
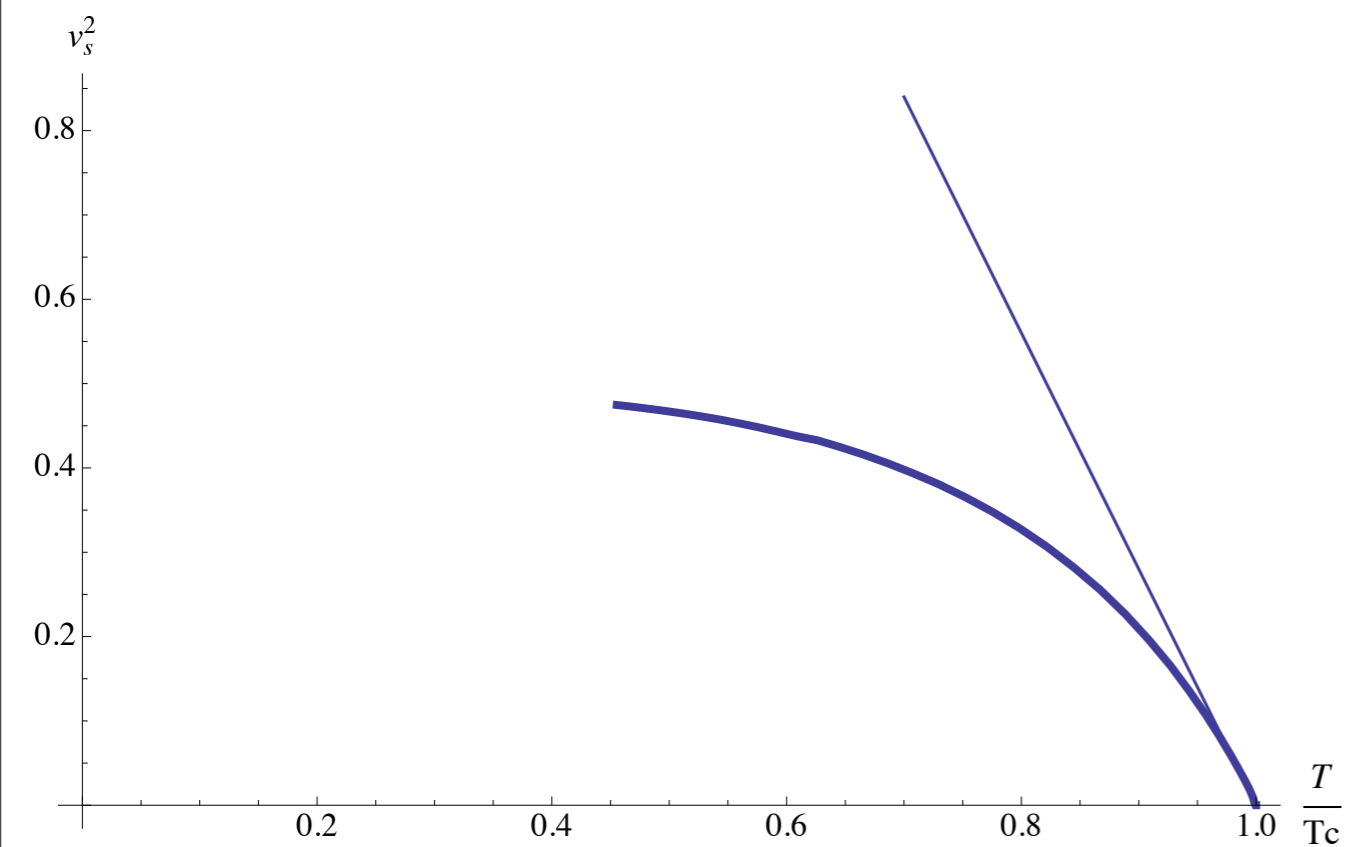
triggers 2nd phase transition to p-wave

[I. Amado, D. Arean, A. Jimenez-Alba, L. Melgar, I. Salazar-Landea, arXiv:1309.5986]

Holography

“ σ_- ” Sector:

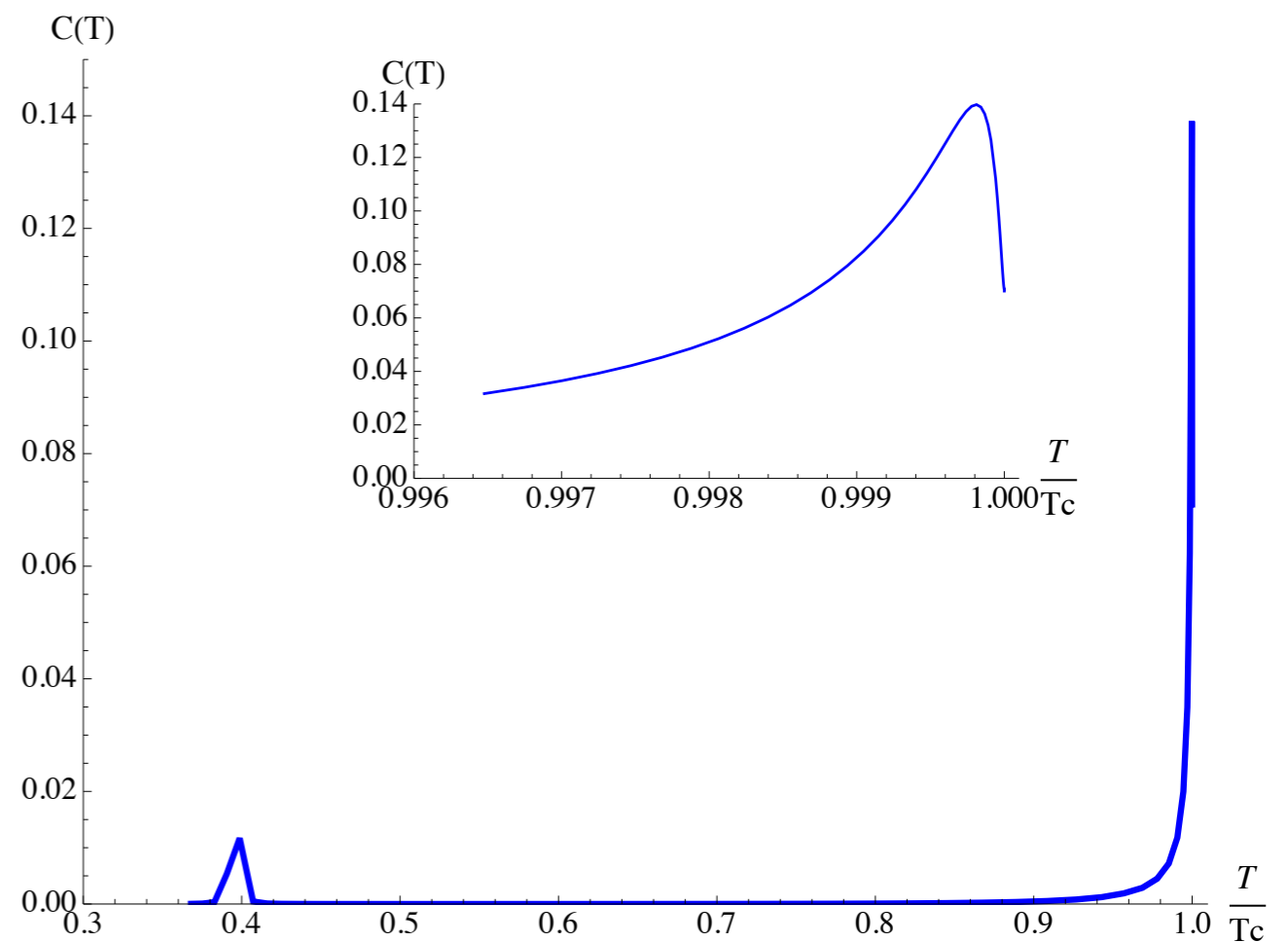
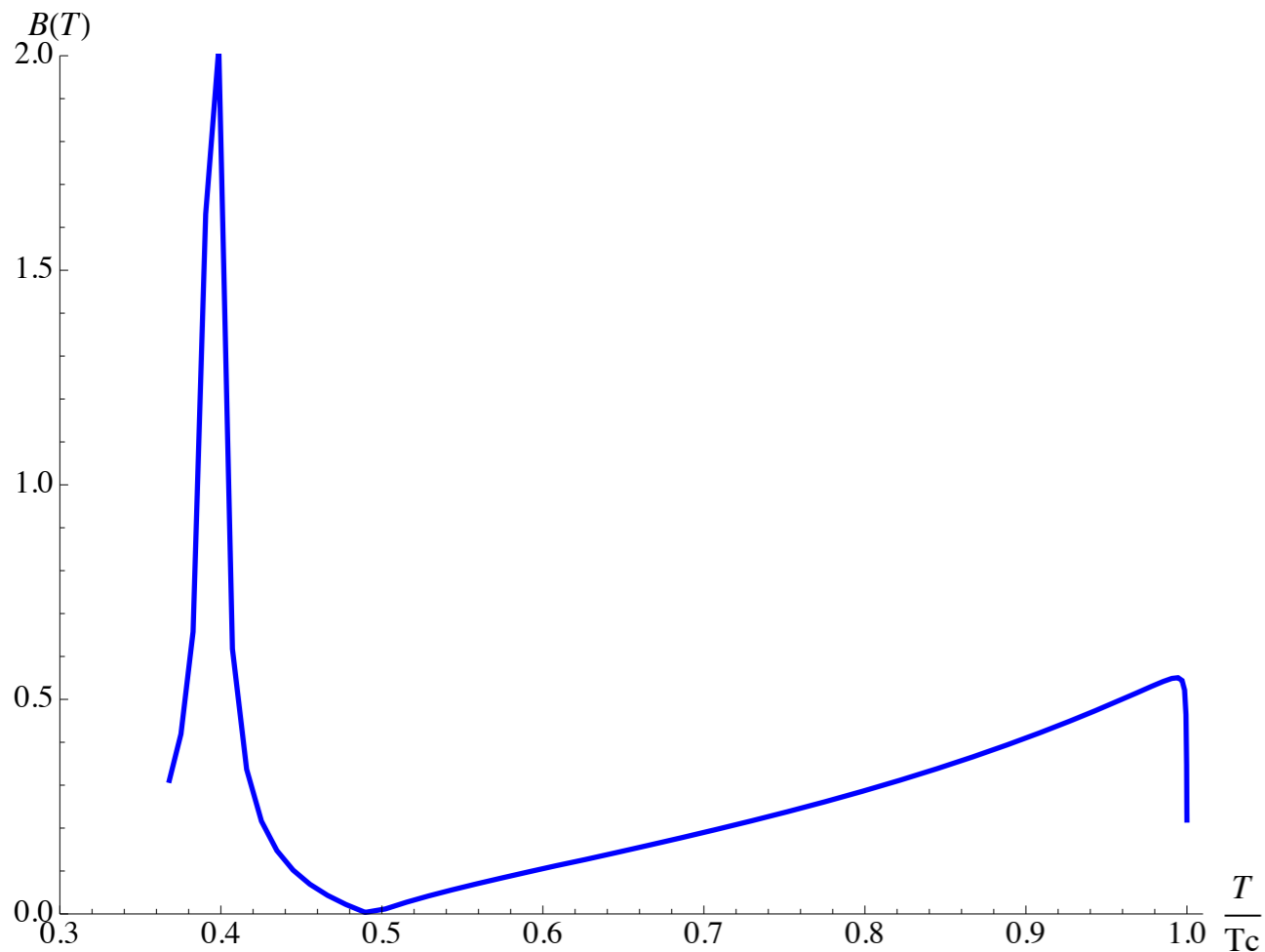
Broken phase, Type I (4th sound) $\omega = v_s k + (b - i\Gamma_s)k^2$



Holography

“ $\sigma_{1,2}$ ” Sector:

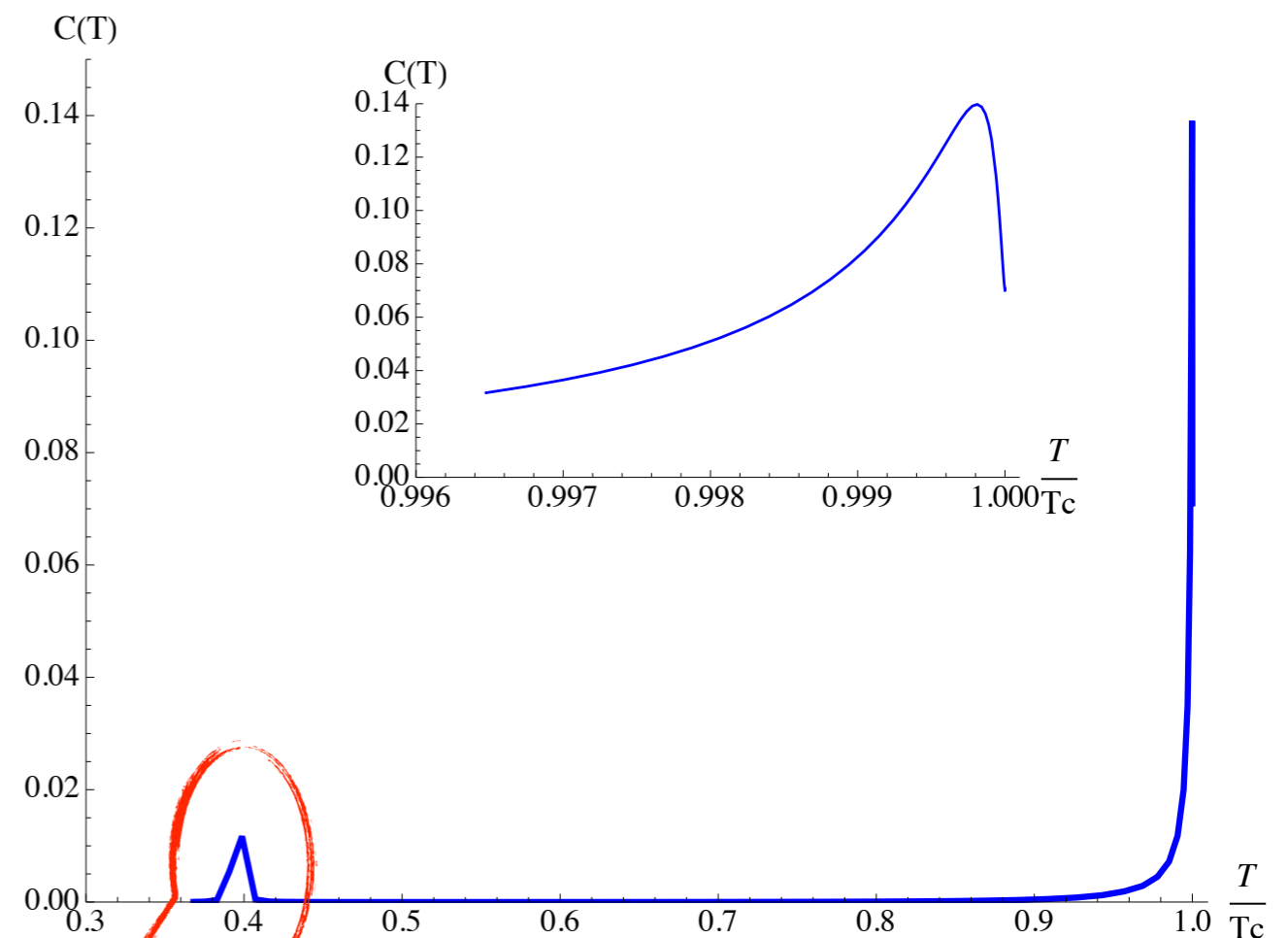
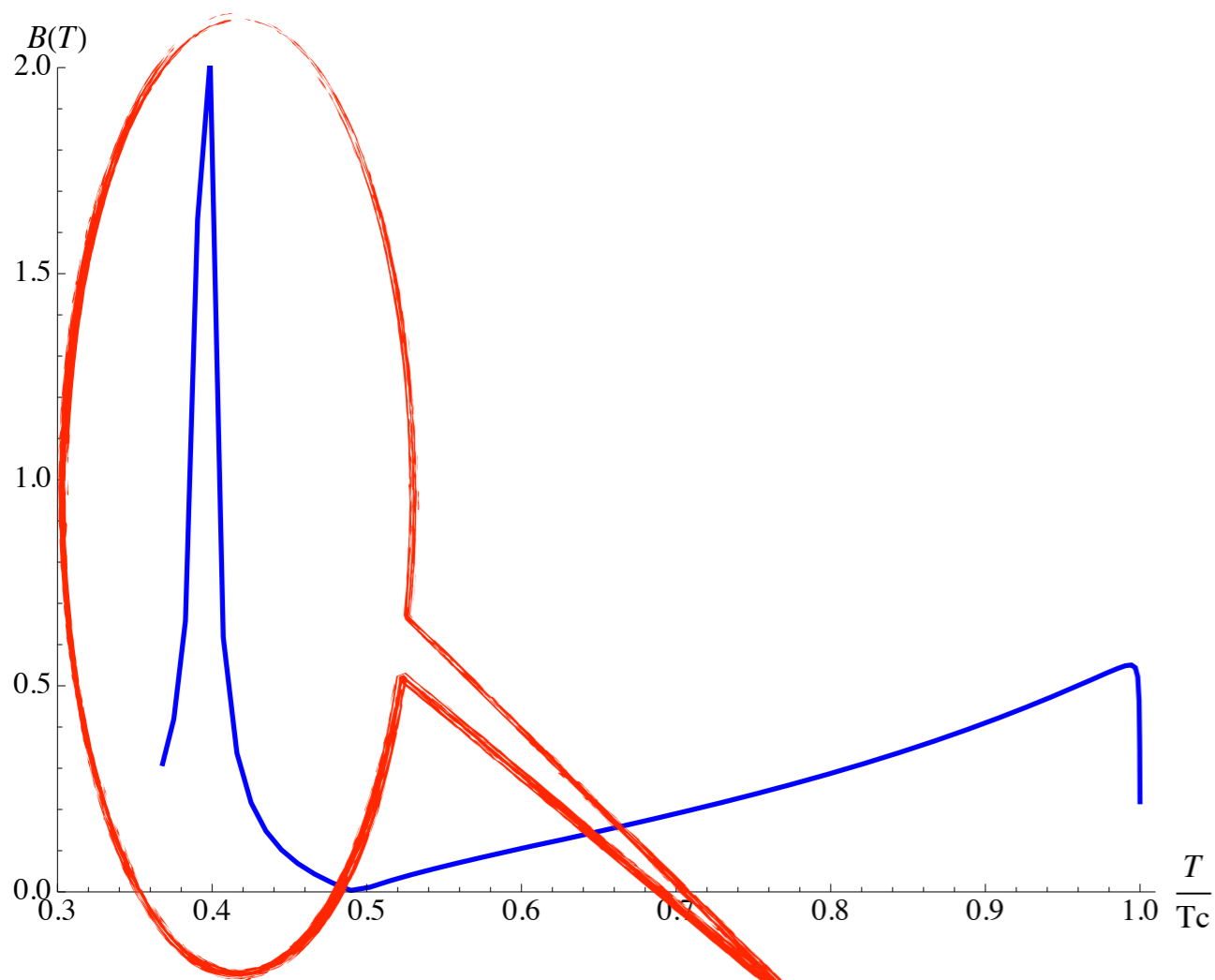
Broken phase, Type II $\omega = (B - iC)k^2$



Holography

“ $\sigma_{1,2}$ ” Sector:

Broken phase, Type II $\omega = (B - iC)k^2$

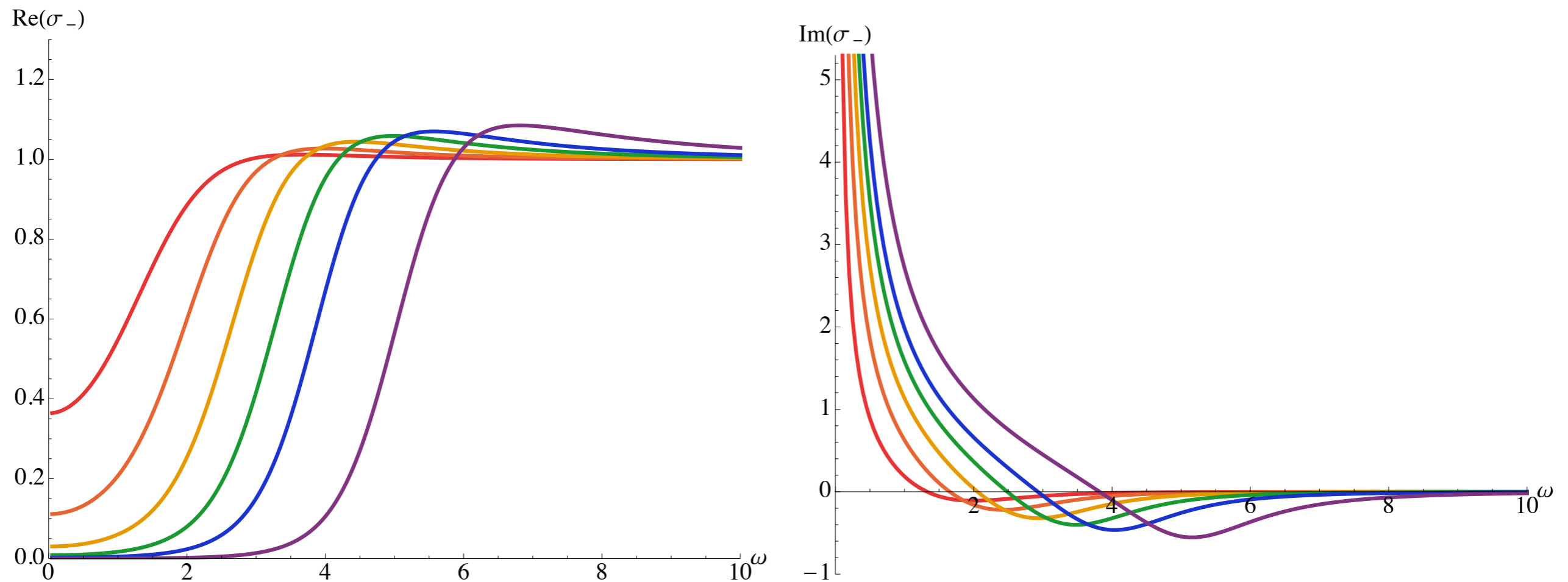


additional unstable mode: p-wave condensate

Holography

“ σ_- ” Sector:

Conductivities related to type I

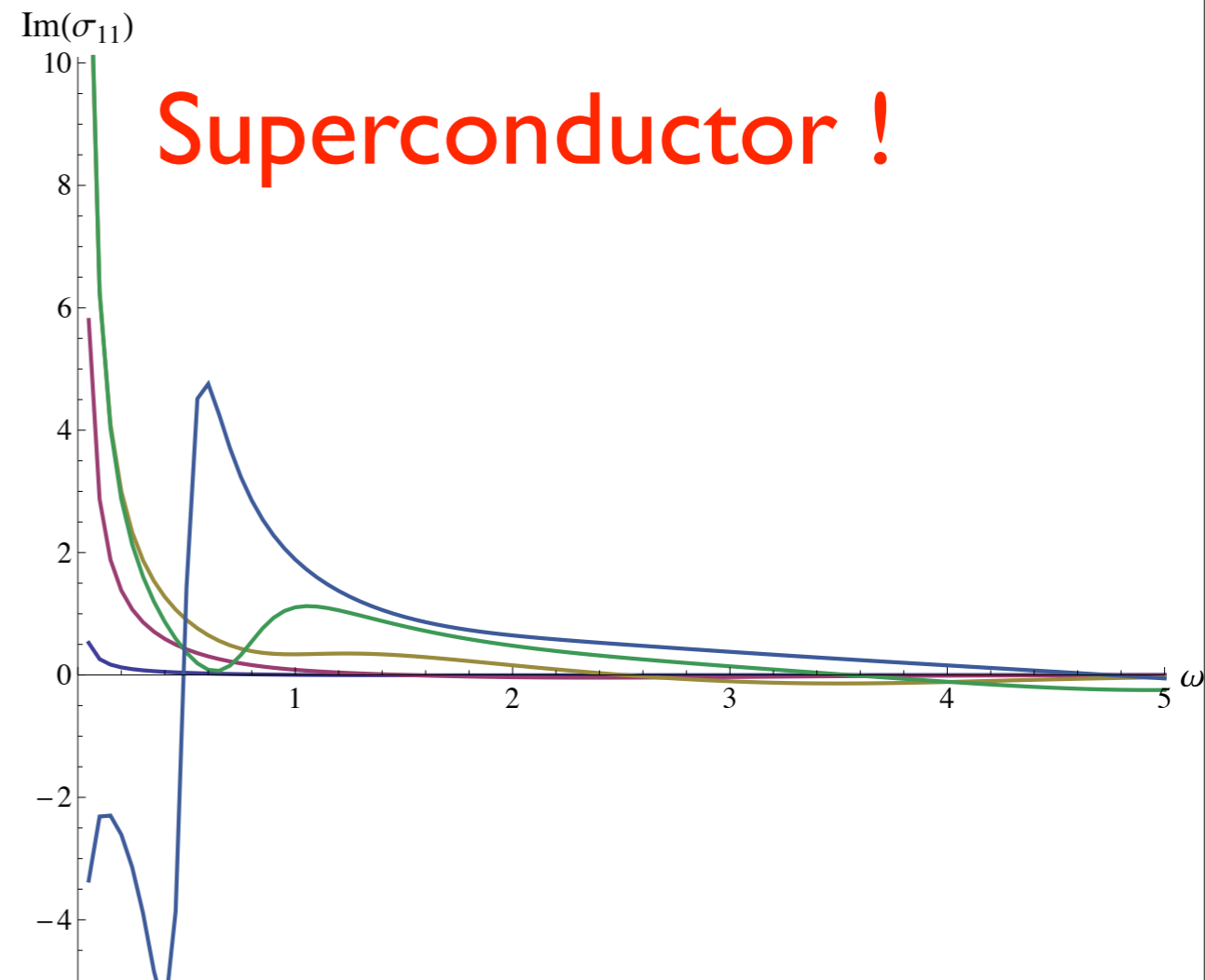
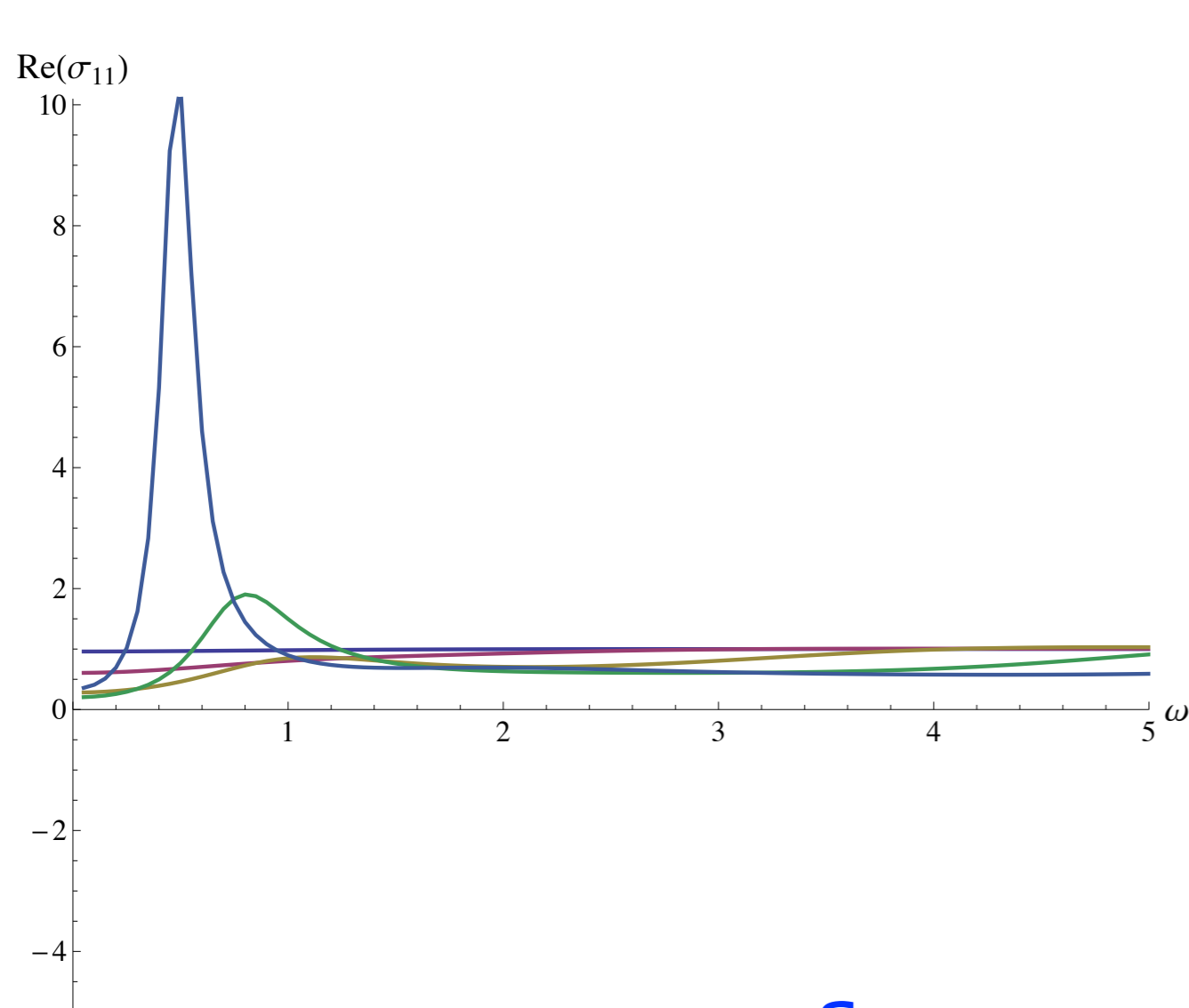


(just HHH model)

Holography

“ $\sigma_{1,2}$ ” Sector:

Conductivities related to type II, diagonal

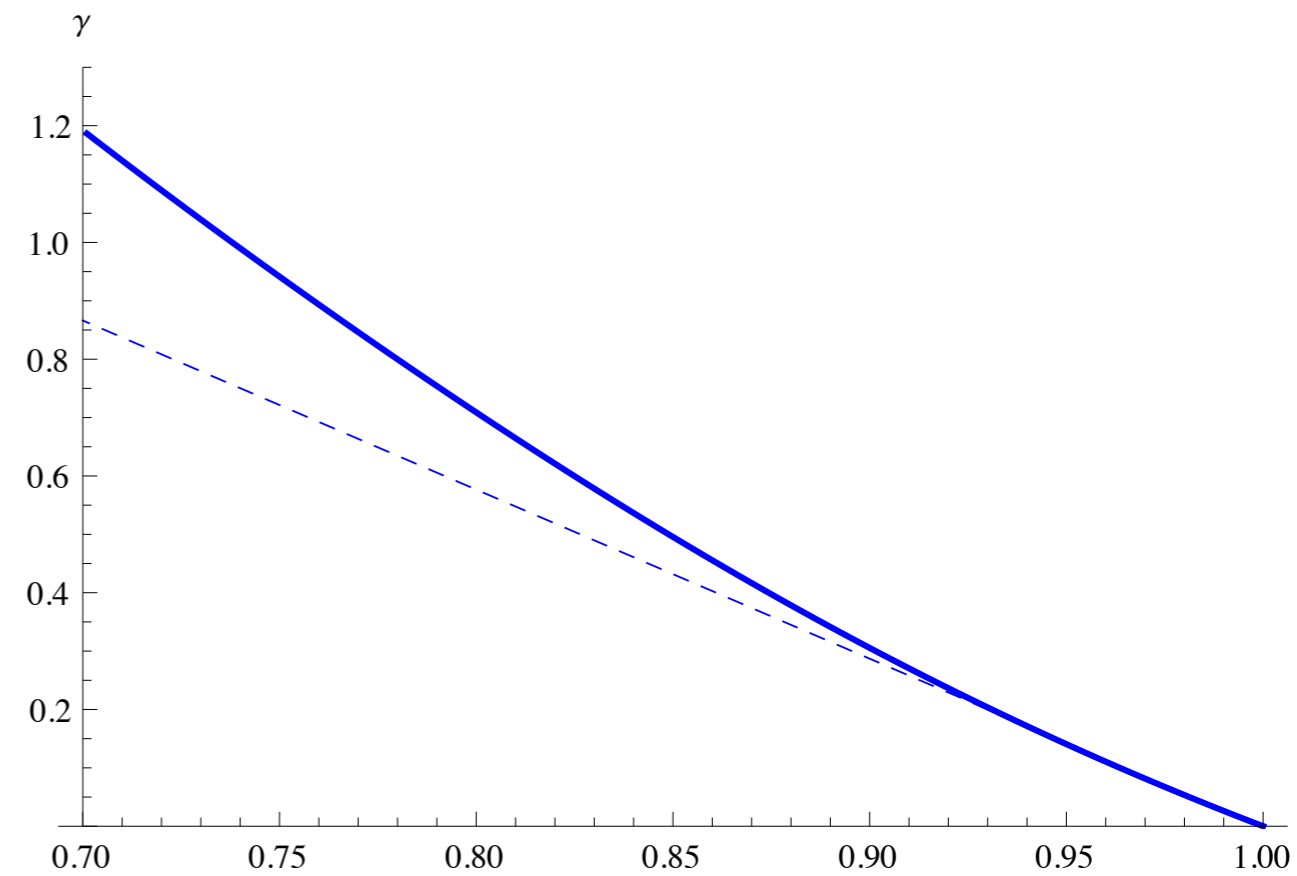
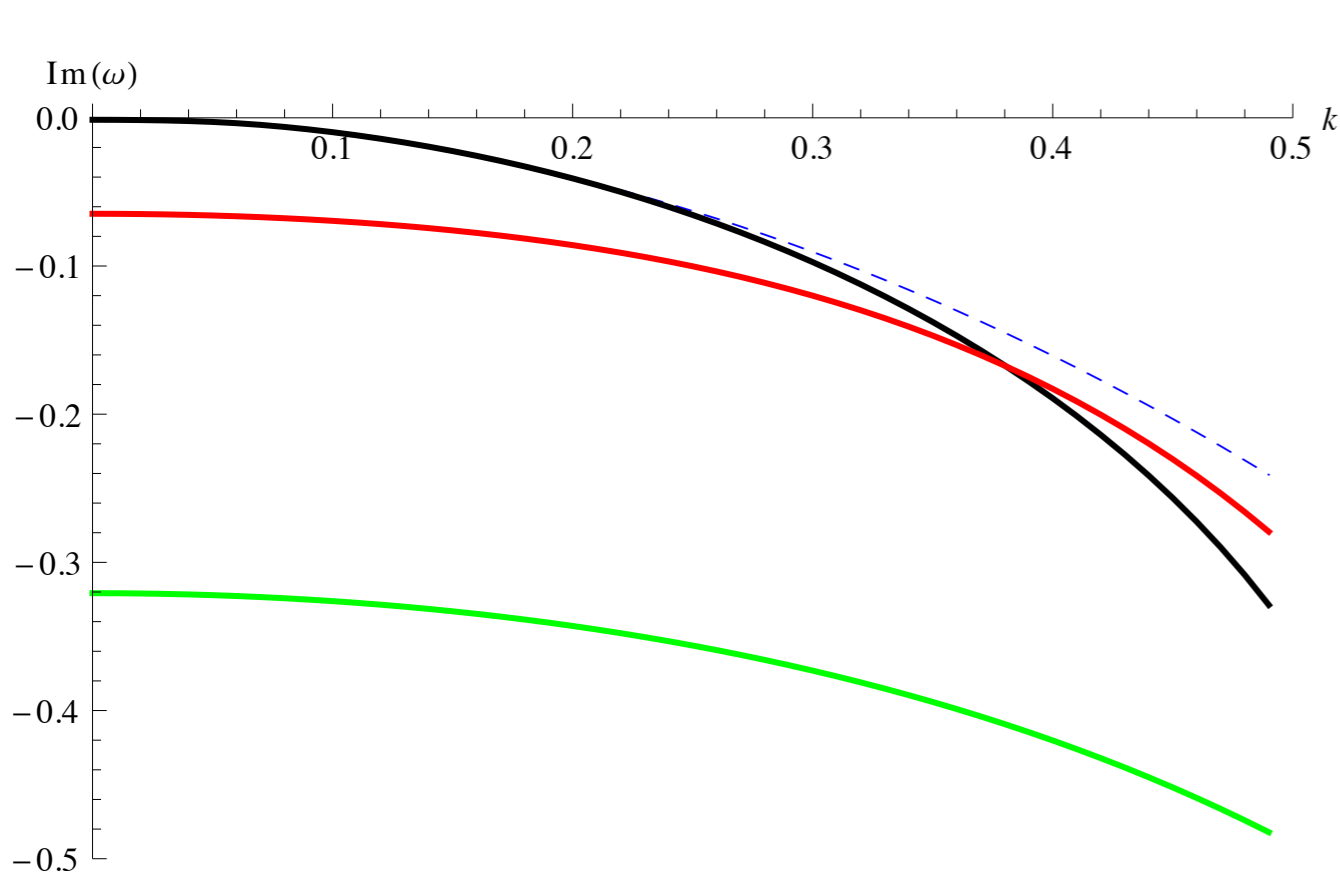


Off-diagonal σ 's: no δ -function pole !

Holography

Fate of the diffusion modes in broken phase
“ σ_- ” Sector: gapped (pseudo) diffusion

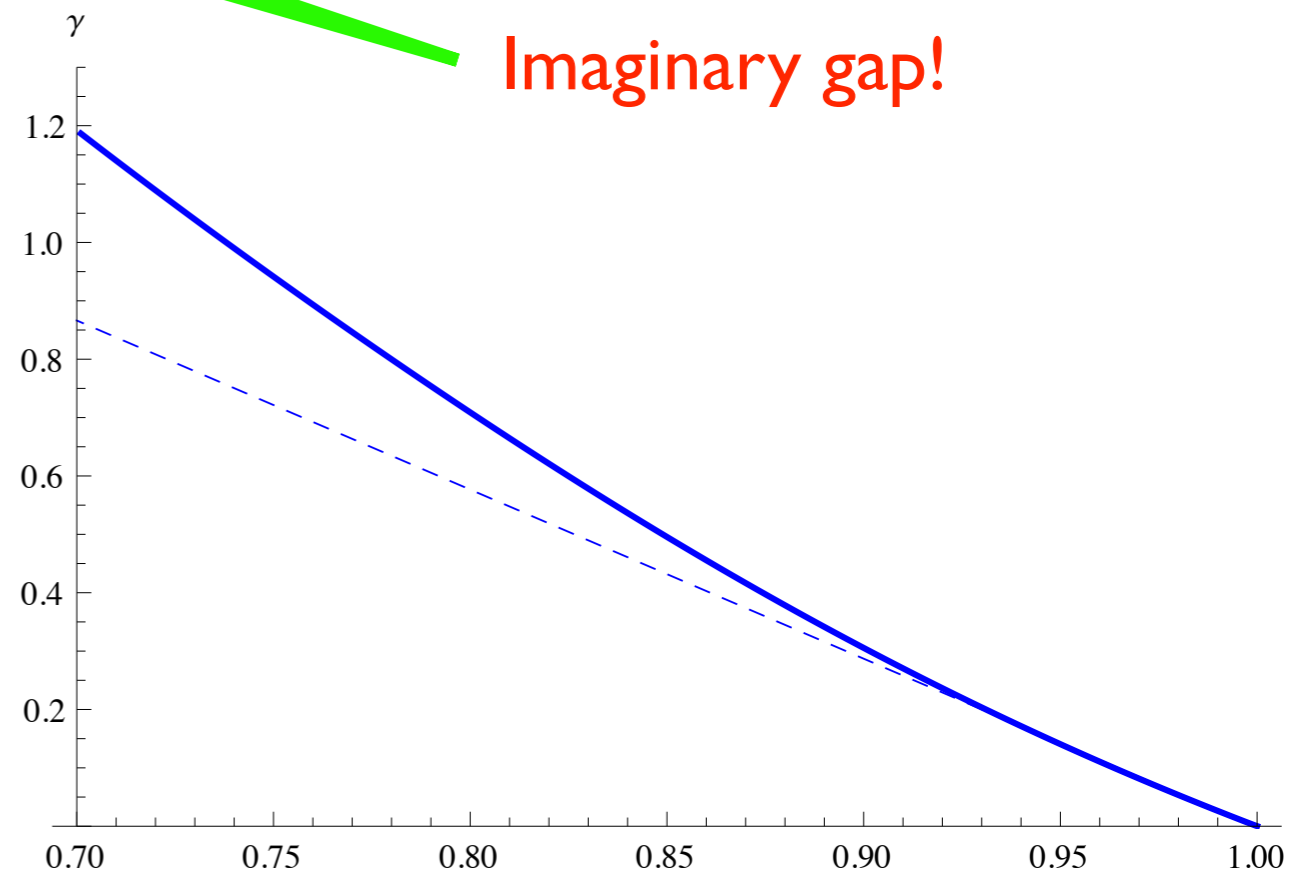
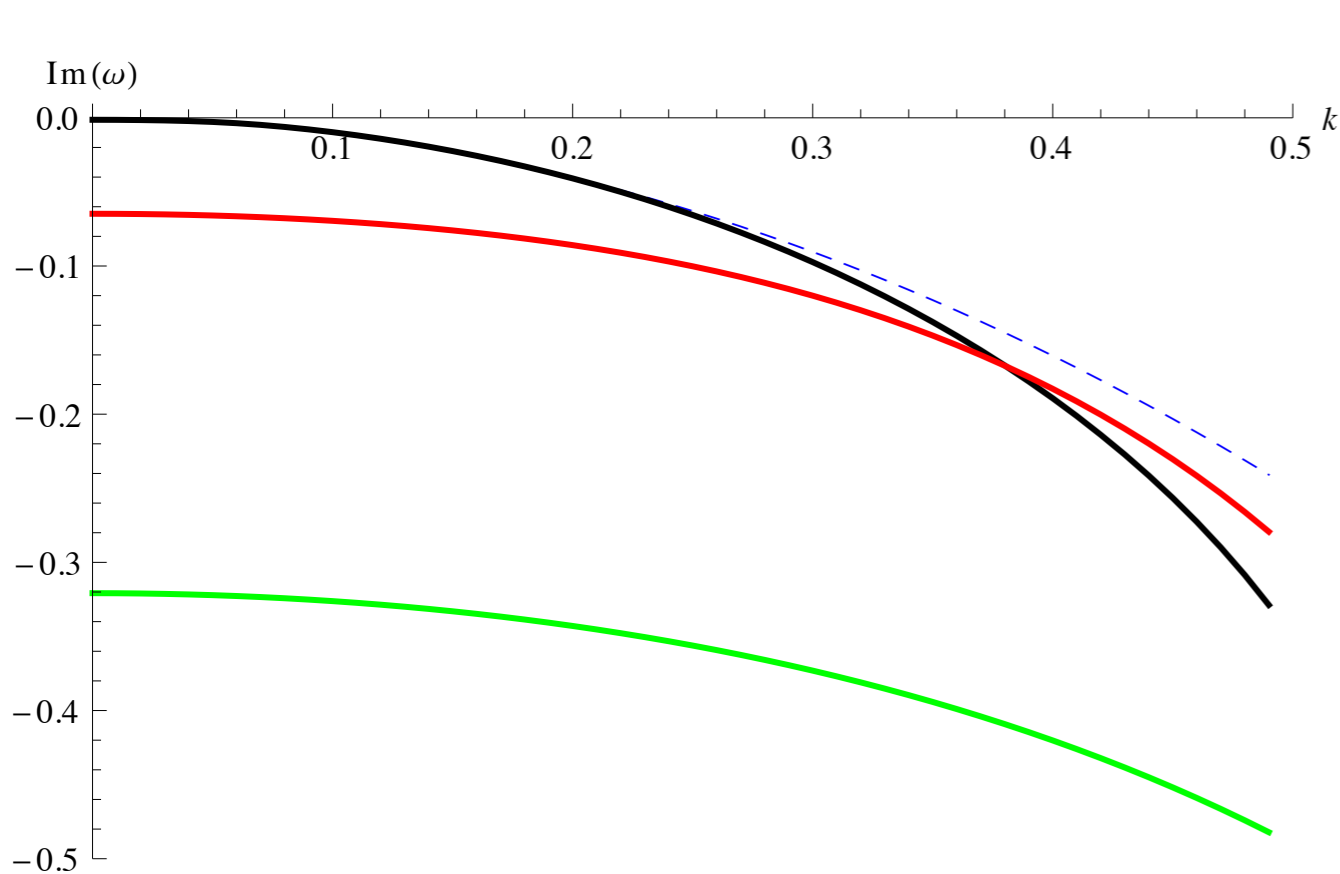
$$\omega = -i\gamma - iDk^2$$



Holography

Fate of the diffusion modes in broken phase
“ σ_- ” Sector: gapped (pseudo) diffusion

$$\omega = -i\gamma - iDk^2$$

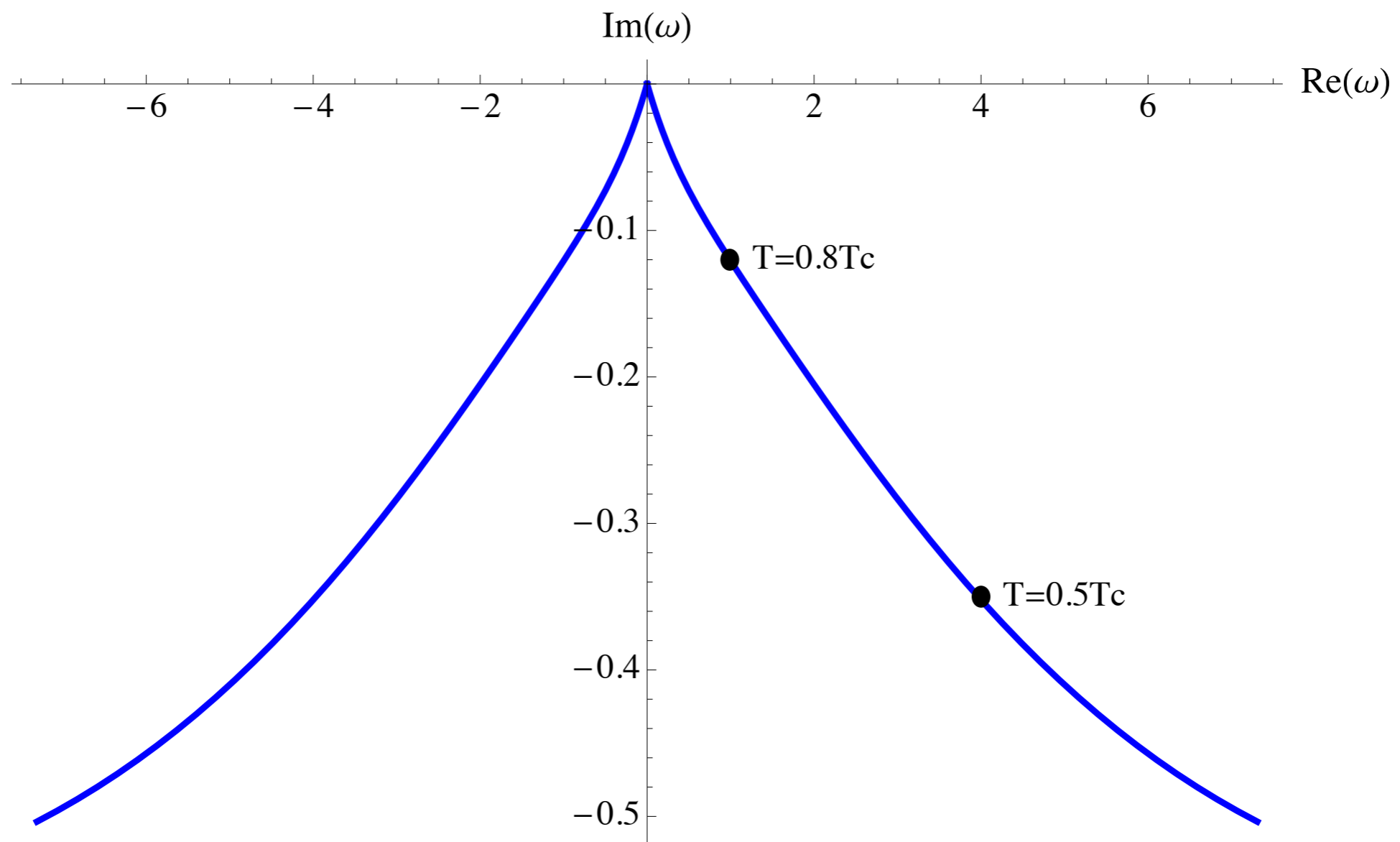


Imaginary gap!

Holography

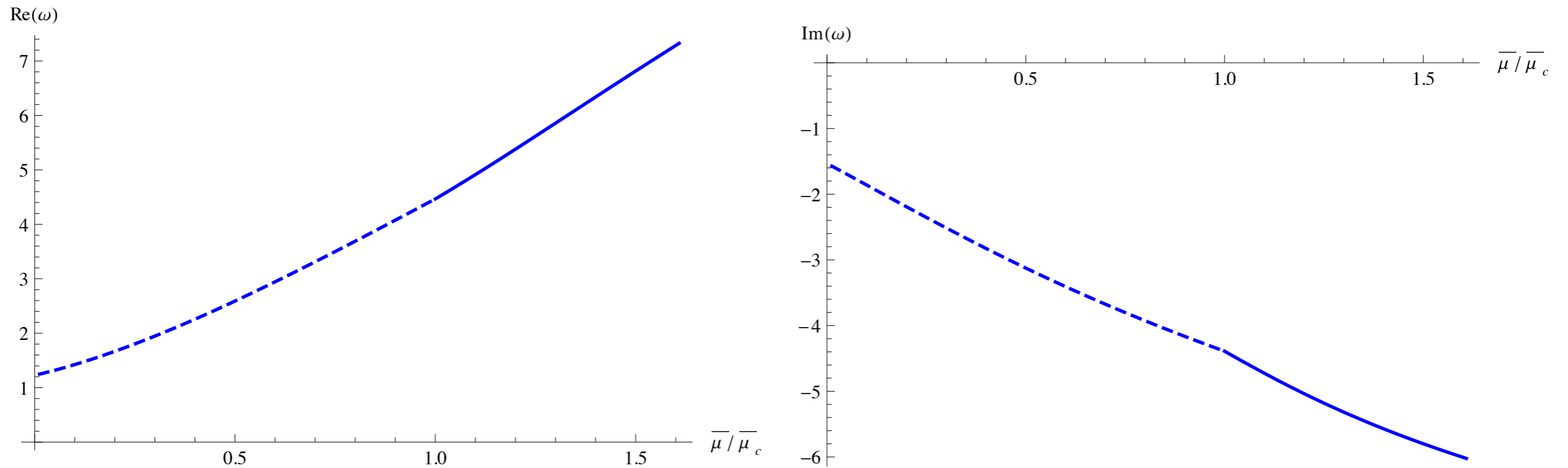
Fate of the diffusion modes in broken phase

$\sigma_{1,2}$ - Sector: 2 gapped modes $\omega = \pm\Omega - i\Gamma$



Holography

“Massive” Goldstone $Re(\omega) \approx 1.1q\mu$



Landau criterion

- Superflow = spatial gauge field on boundary
- Critical superfluid velocity

- T=0 via boosts $\omega(p) + \vec{p} \cdot \vec{S} \leq 0 \Rightarrow S_c = \min \frac{\omega(p)}{p}$

- T>0 more complicated

- basic idea: negative energy , instability

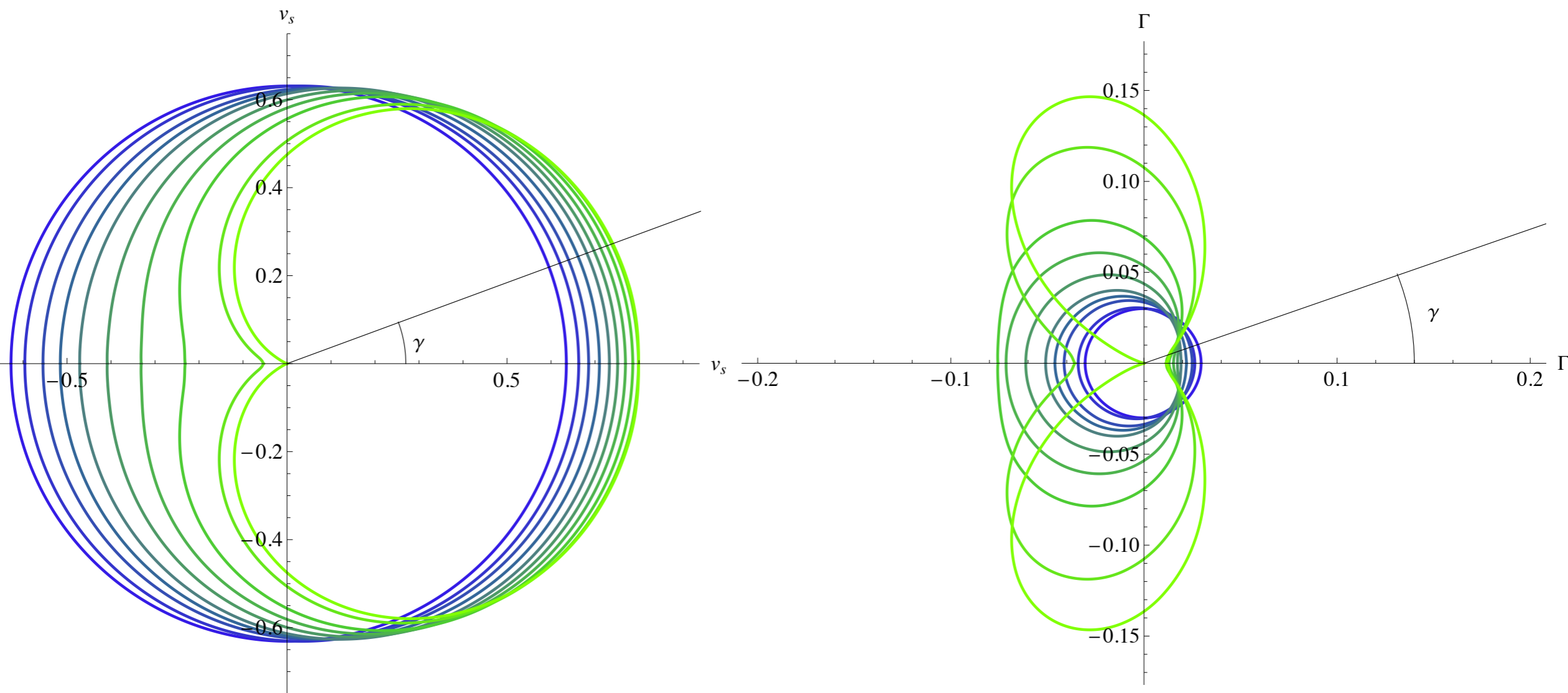
- QNMs:

$$\Re(\omega(p, S, T)) \leq 0$$

$$\Im(\omega(p, S, T)) \geq 0$$

Landau criterion

Type I Goldstone

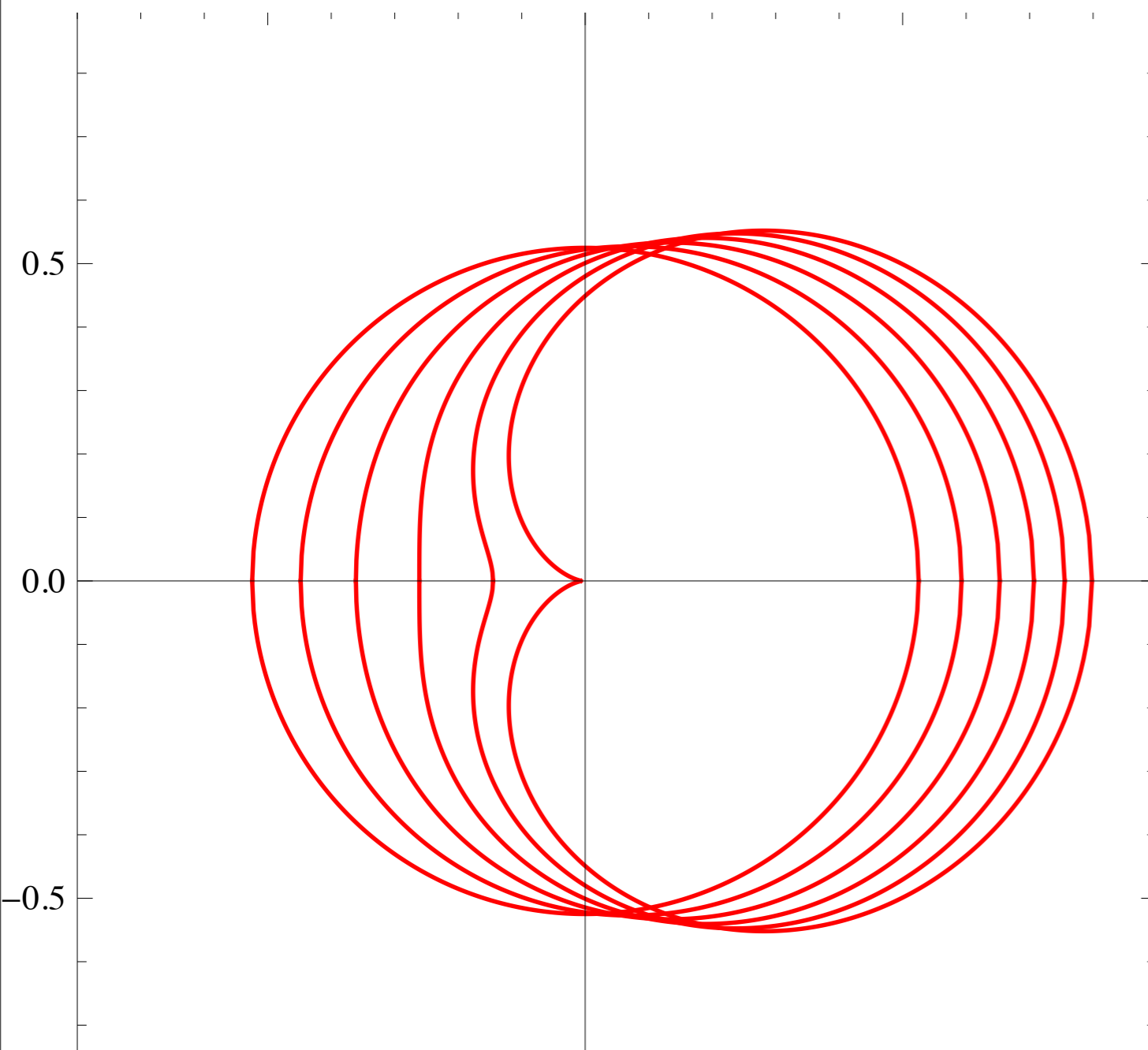


Landau criterion

Weak Coupling

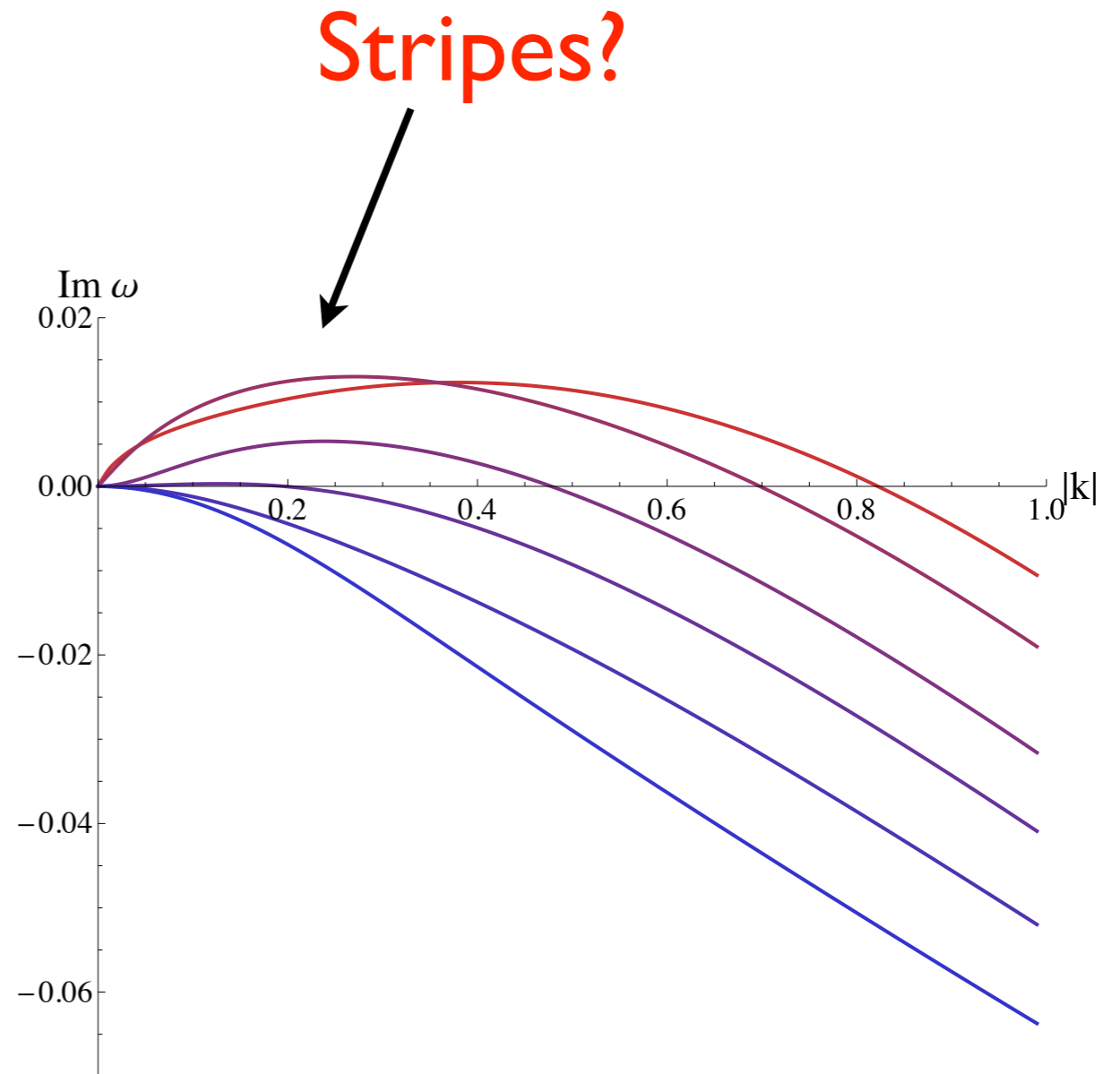
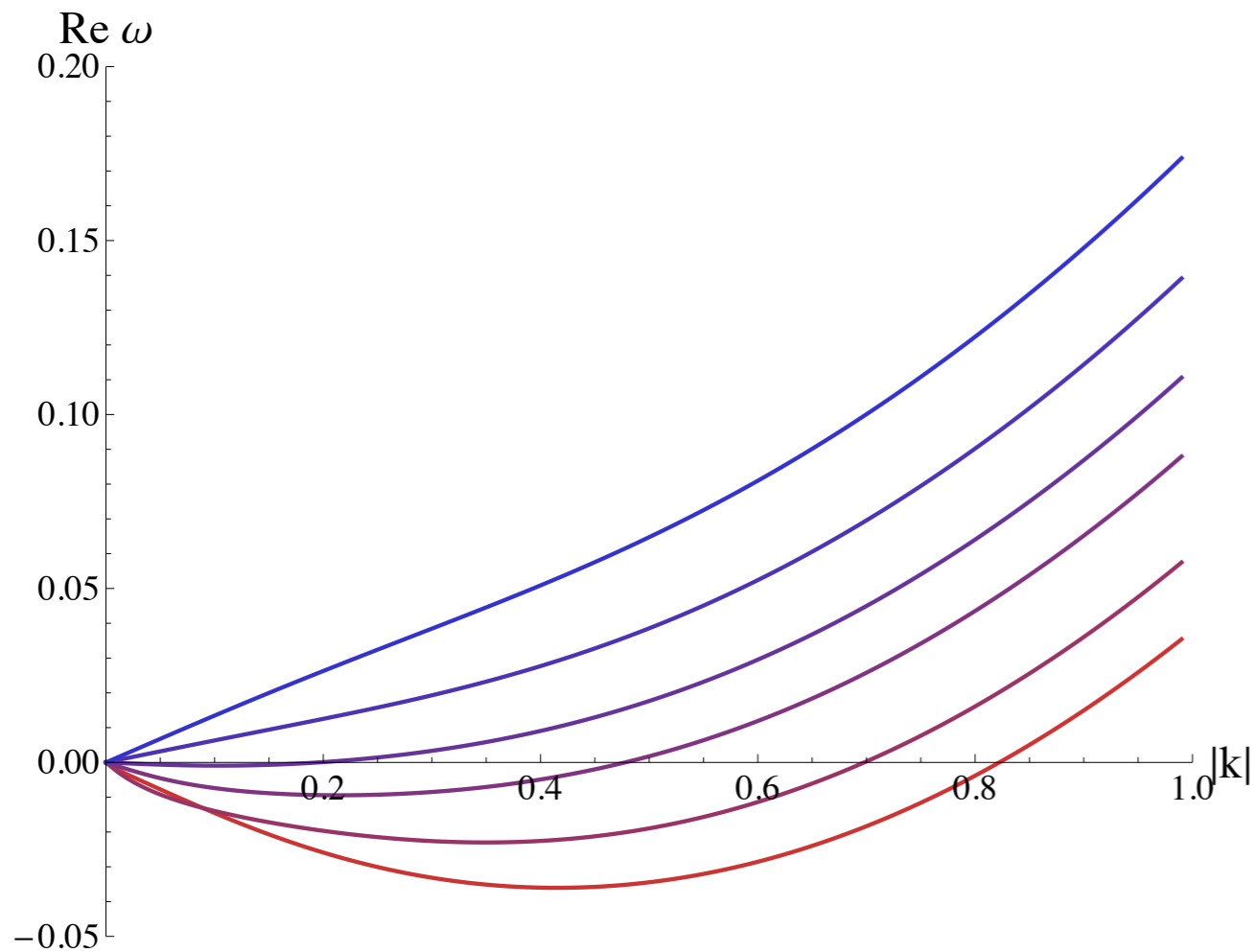
(courtesy: A. Schmitt)

[Alford, Mallavarpu, Schmitt, Stetina] to appear



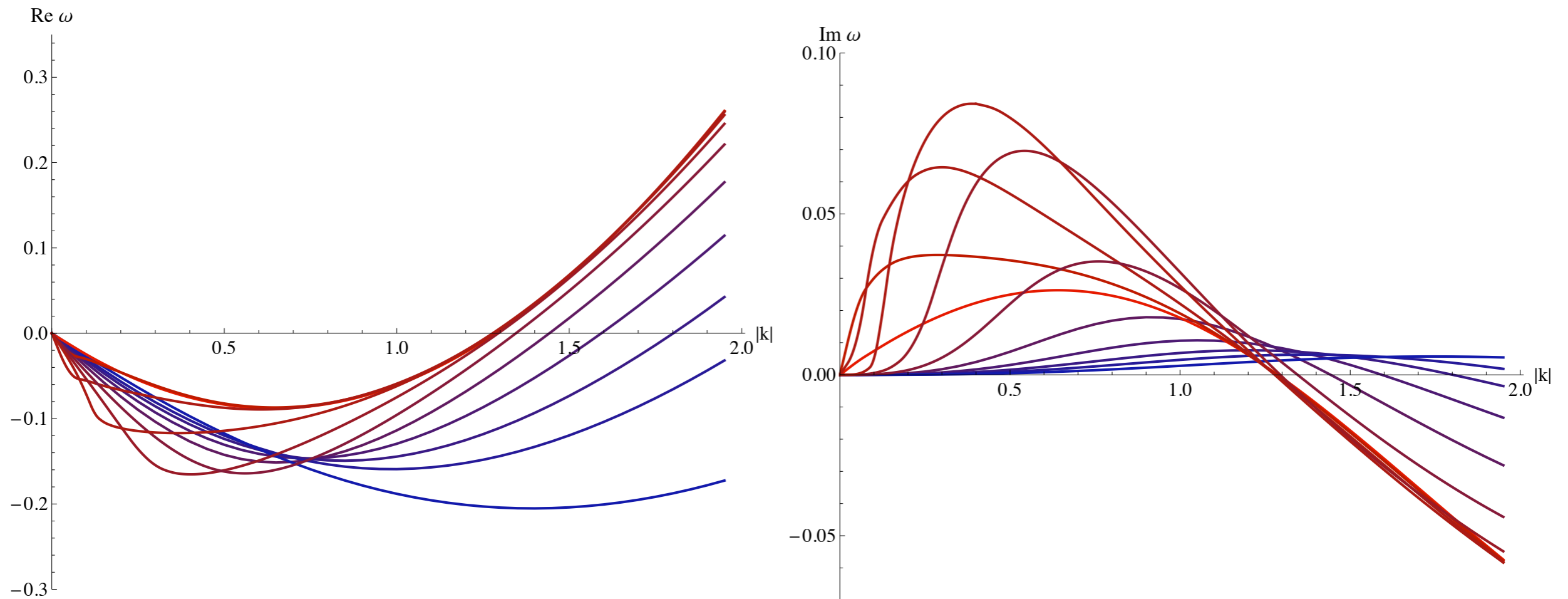
Landau criterion

Type I Goldstone



Landau criterion

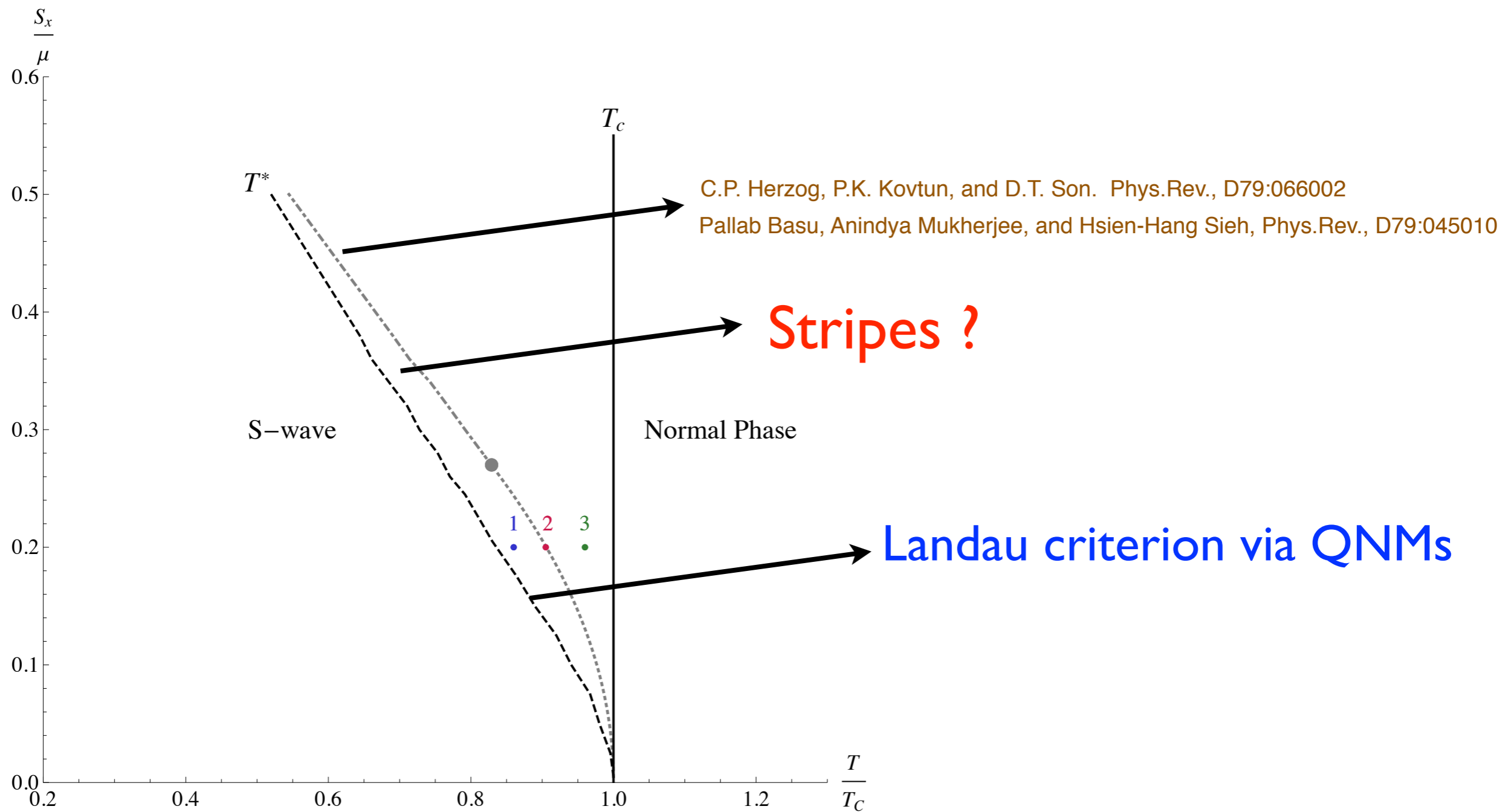
Type II Goldstone



**Does not support finite superflow!
Superconductor but not Superfluid**

Landau criterion

Phase diagram based on Landau criterion



Ungauged Model

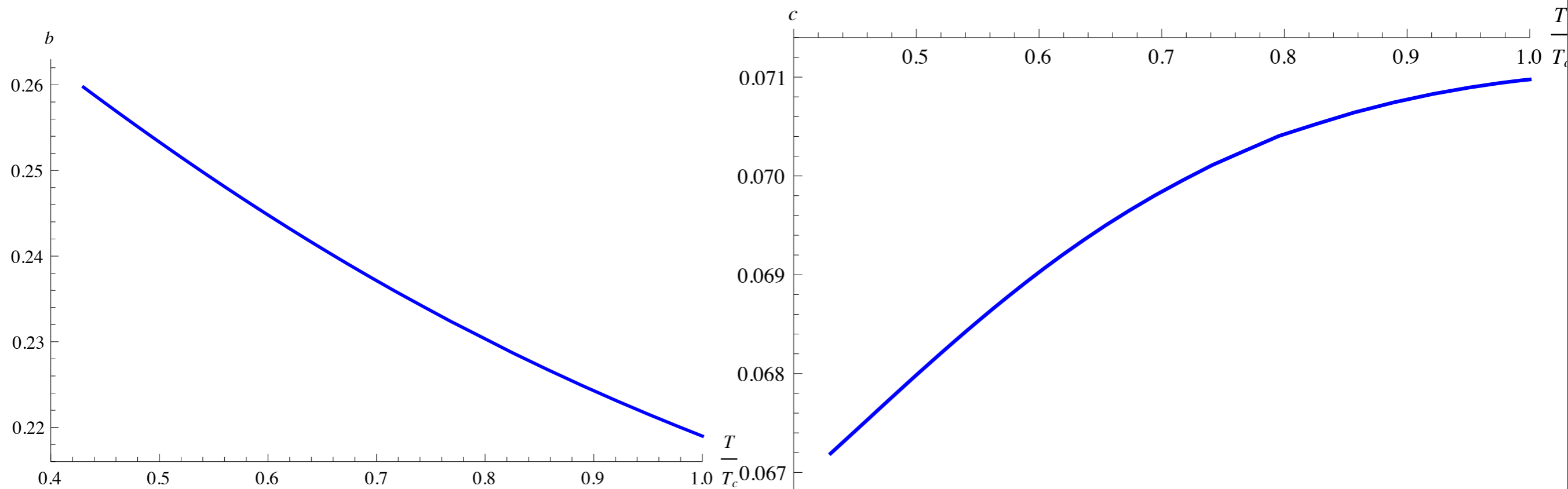
What if $SU(2)$ only global in AdS bulk?

- $U(1)$ gauge field + scalar doublet
- Global symmetry sufficient to chose vacuum
- $SU(2)$ = “Outer Automorphism”
- Decomposition: HHH-superfluid + scalar
- Still type II Goldstone?
- Massive ?

Ungauged Model

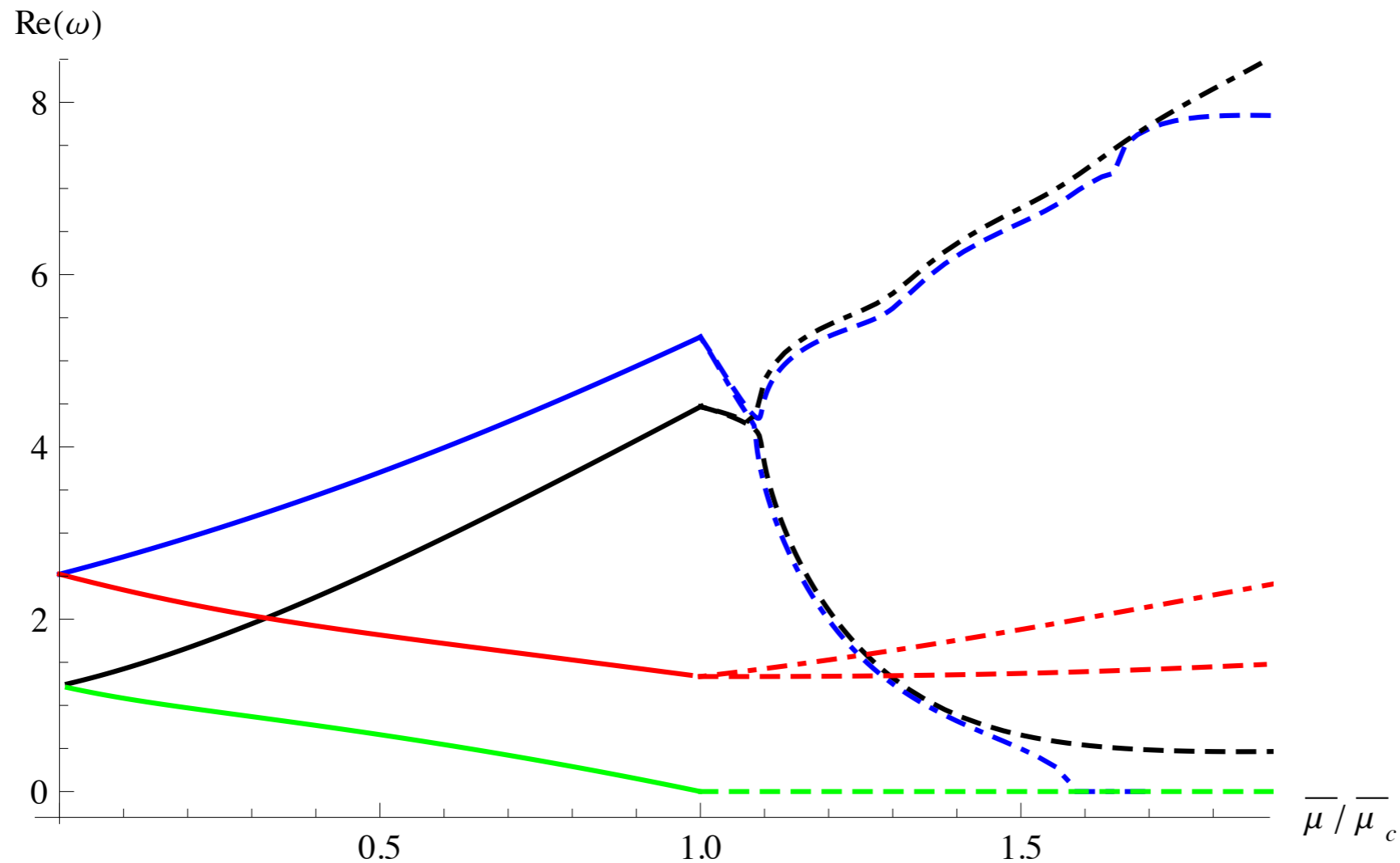
Massless mode in 2nd scalar!

Dispersion relation: $\omega = (b + ic)k^2$



Ungauged Model

Massive mode:



Does not obey theory!! $\omega \neq q\mu$

Effective $U(2)$ symmetric action only
for lower Energies!

Summary

- Type II Goldstone QNMs !
- Compare to weak coupling
- Universality of Pseudo-diffusion ?
- “Un-gauged” model:
no $SU(2)$ gauge fields, violates some Theorems
- Backreacted models
1st, 2nd, 4th sound modes etc.
- Superflow: striped phases ?

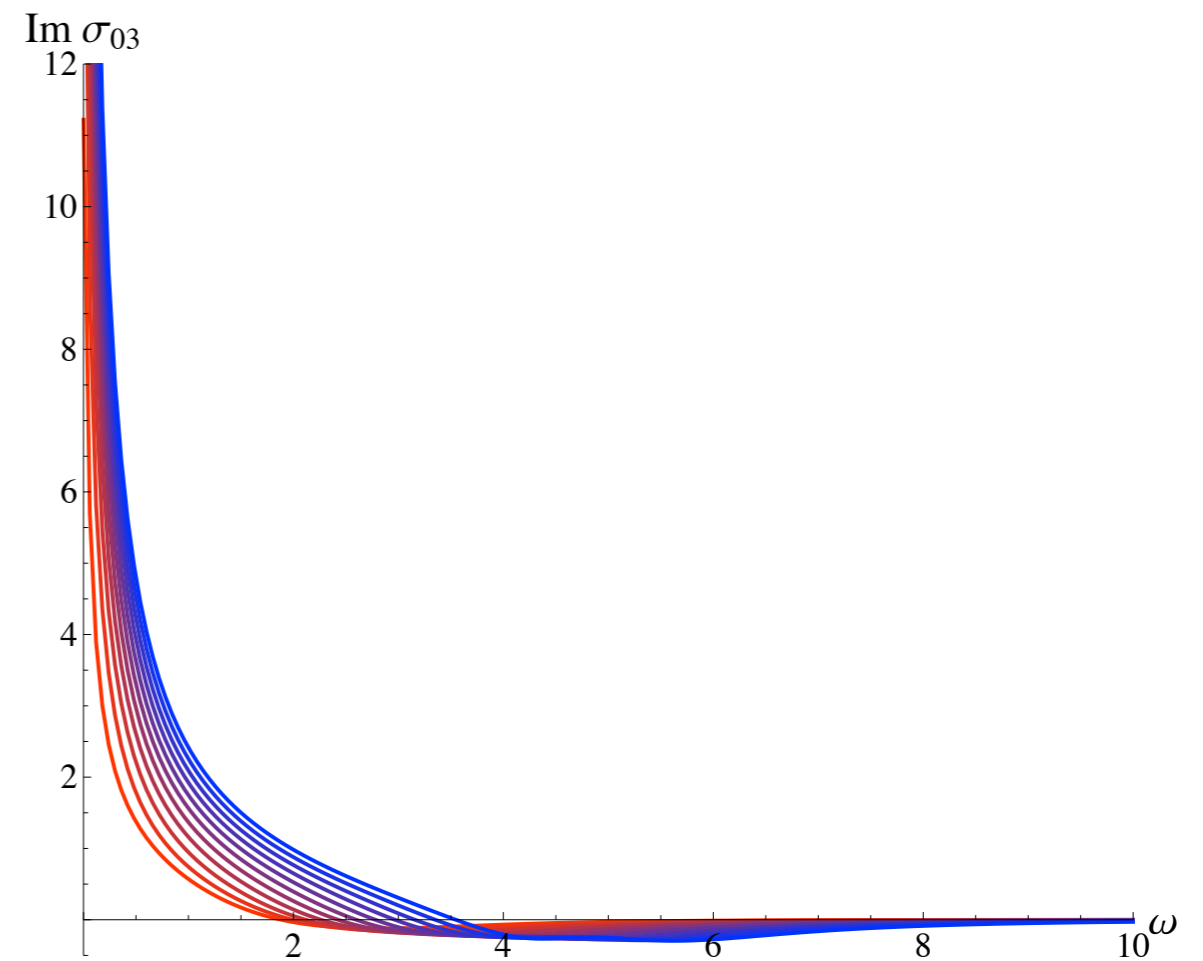
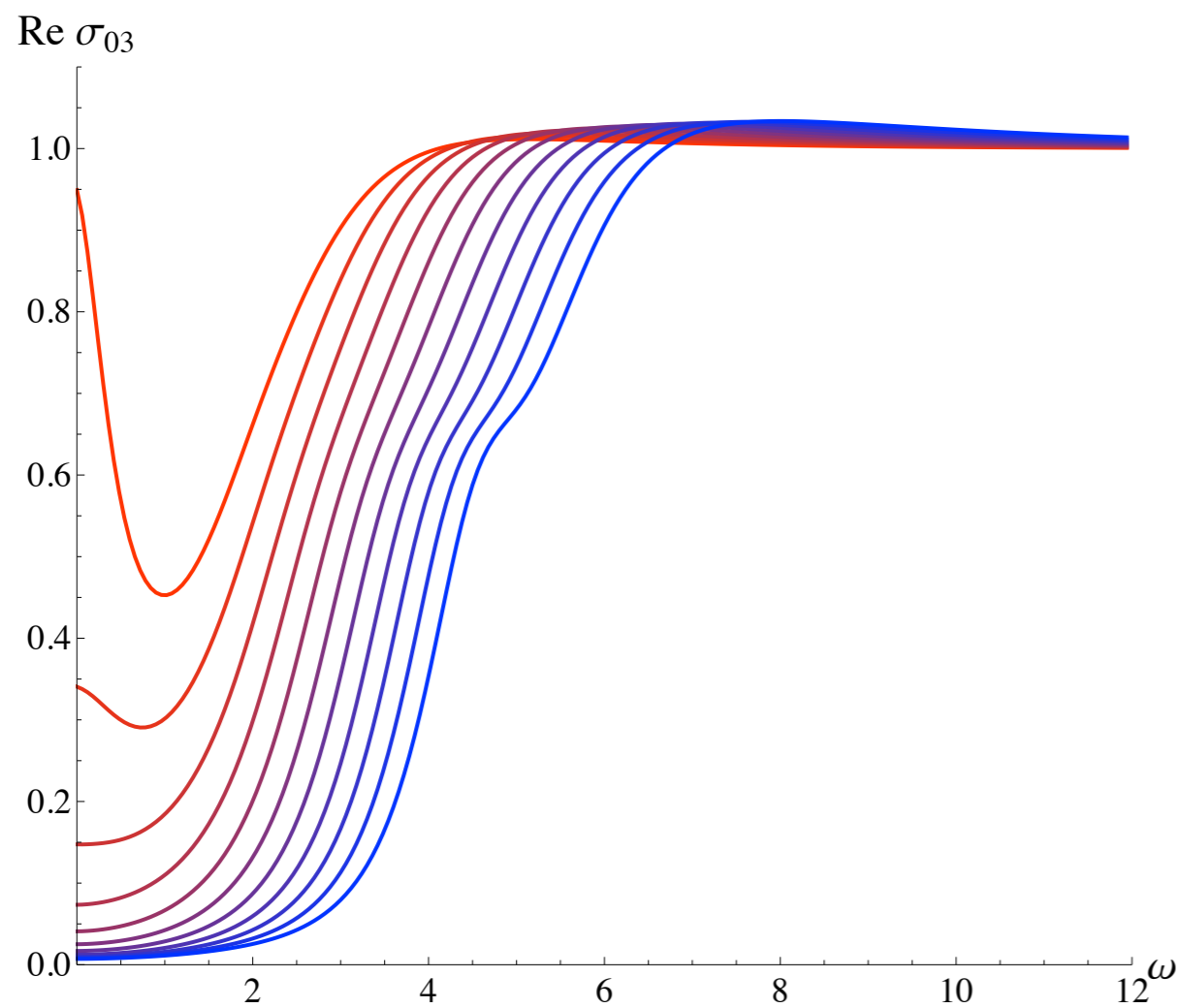
Summary

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THANK YOU!

Drude Peak

Due to gapped “Pseudo” diffusion mode !



Holography

Conductivities related to type II, off-diagonal
no Superconductor !

