Disentangling Topological Insulators by Tensor Networks

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Table of Contents

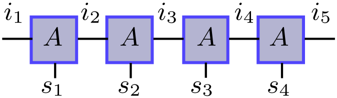
Tensor network approach to quantum manybody systems

Tensor network:

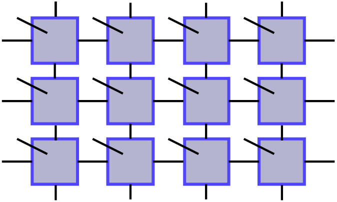
way to avoid exponential complexity of many-body problems

Tensor network wave functions of various kinds:

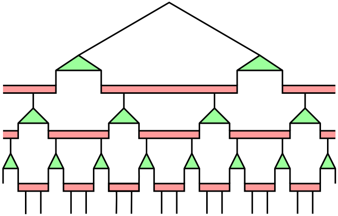
MPS (matrix product state) or DMRG



PEPS (projected entangled pair state)



MERA
(multiscale entanglement
renormalization ansatz)



Tensor network approach to quantum manybody systems

- Representing many-body wavefunctions by contracting many tensors
  DMRG, MPS, MERA, PEPS, etc.

$$|\Psi\rangle = \sum_{s_1, s_2, s_3, s_4 \cdots} C^{s_1, s_2, s_3, s_4 \cdots} |s_1, s_2, s_3, s_4 \ldots\rangle$$

Product state:

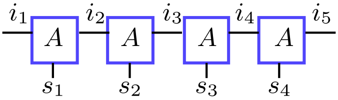
$$|\Psi\rangle = \sum_{\{s_a\}} A^{s_1} A^{s_2} A^{s_3} \cdots |s_1, s_2, s_3, s_4 \ldots\rangle = \prod_i \sum_{s_i} A^{s_i} |s_i\rangle$$

EE = 0

physical degrees of freedom

MPS (matrix product state) :

$$|\Psi\rangle = \sum_{\{s_a\}} \sum_{\{i_n = 1, \cdots, \chi\}} A^{s_1}_{i_1, i_2} A^{s_2}_{i_2, i_3} A^{s_3}_{i_3, i_4} A^{s_4}_{i_4, i_5} \cdots |s_1, s_2, s_3, s_4 \ldots\rangle$$

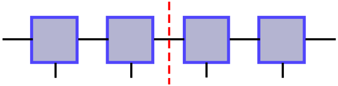


auxiliary index

structure of tensor-network and entanglement entropy

Matrix product state (DMRG):

$$|\Psi\rangle = \sum_{\{s_a\}} \sum_{\{i_n=1,\cdots,\chi\}} A^{s_1}_{i_1,i_2} A^{s_2}_{i_2,i_3} A^{s_3}_{i_3,i_4} A^{s_4}_{i_4,i_5} \cdots |s_1, s_2, s_3, s_4 \ldots\rangle$$



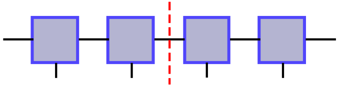
EE: $S_A \leq \log \chi$

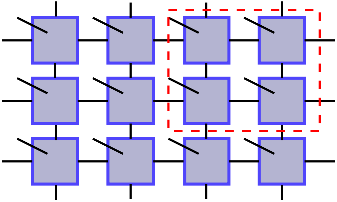
χ :dimension of the aux space ("bond dimension")

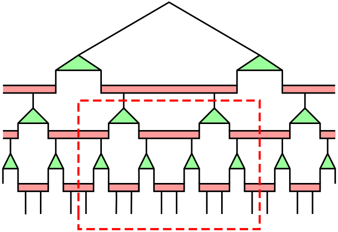
Area law scaling in 1D: quite generic in gapped quantum ground states.

Can we extract information from tensor network in an effeicent way ?

--> Cutting up tensor networks!



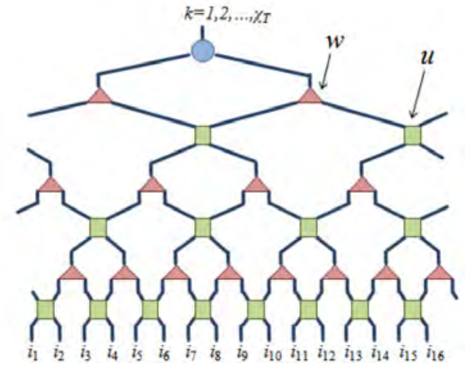


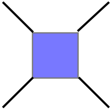


Cutting up = defining "reduced density matrix"

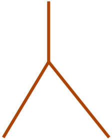
multiscale entanglement renormalization ansatz (MERA)

[Vidal (07-08)]





"disentangler"



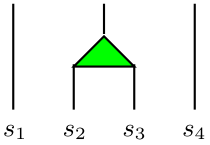
isometry
(coarse-graining)

block spin decimation and disentangler

- Block spin decimation

$$\rho_{tot} = |\Psi\rangle\langle\Psi|$$

$$\rho_{23} = \mathrm{Tr}_{14}\rho_{tot} = \sum_i p_i |\phi_i\rangle\langle\phi_i|$$



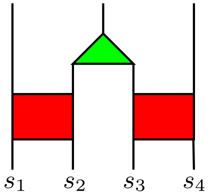
small p_i --> throw away

- Disentangler

$$|\Psi\rangle = \frac{(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)_{12}}{\sqrt{2}} \frac{(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)_{34}}{\sqrt{2}} \qquad p_i = \frac{1}{4}(\forall i)$$

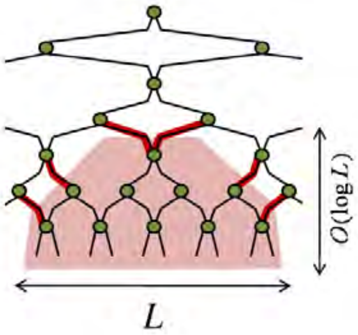
$$U \frac{(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)_{12}}{\sqrt{2}} = |\uparrow\uparrow\rangle_{12}$$

$$\mathrm{Tr}_{14}\left[U_{12} \otimes U_{34} \rho_{tot} (U_{12} \otimes U_{34})^\dagger\right]$$



MERA and holographich entanglement entropy

[Swingle (09)]

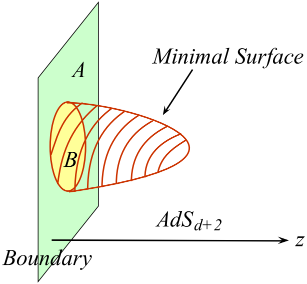


EE for (1+1)D case: $S_A \sim \log(l/a)$

Geometry <--> Entanglement

- Holographic formula for EE

$$\boxed{S_A = \frac{\text{Area of minimal surface } \gamma_A}{4G_N}}$$



- Entanglement <--> geometry

$$ds^2 = g_{uu}du^2 + \frac{e^{2u}}{\epsilon^2}d\vec{x}^2 + g_{tt}dt^2$$

$$S_A = \frac{1}{4G_N}\cdot\frac{V_{d-1}}{\epsilon^{d-1}}\int_{u_{\rm IR}(=-\infty)}^{u_{\rm UV}(=0)} du\sqrt{g_{uu}}e^{(d-1)u}$$

Can we make "AdS/MERA" more precise ?

continuous MERA (cMERA)

[Haegeman-Osborne-Verschelde-Verstraete (11)]

- Quantum circuit representation of the target states:

$$|\Psi(u_{\rm IR})\rangle \equiv |\Omega\rangle \qquad |\Psi(u_{\rm UV})\rangle \equiv |\Psi\rangle$$

$$\infty$$

$$|\Psi(u)\rangle = U(u, u_{\rm IR})|\Omega\rangle$$

$$u_{\rm UV} = 0$$

$$|\Psi\rangle = U(0, u)|\Psi(u)\rangle$$

$$u_{\rm IR} = -$$

- MERA evolution operator

$$U(u_1, u_2) = P\exp\left[-i\int_{u_2}^{u_1}(K(u) + L)du\right]$$

disentangler

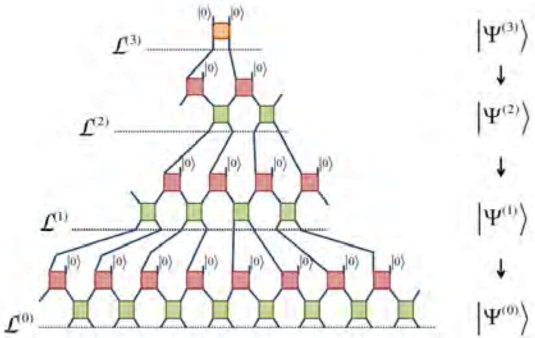
coarse-graining

- Optimizing |Omega>, U --> true ground state

MERA and quantum circuit

- Tensor network method can be formulated as a quantum circuit (successive applications of unitary transformations)

- For MERA: add dummy states |0>



cMERA for free boson

- free boson in d+1 dim:

$$H = \frac{1}{2} \int d^d k \left[\pi(k)\pi(-k) + \epsilon_k^2 \cdot \phi(k)\phi(-k) \right]$$

$$\phi(k) = \frac{a_k + a^\dagger_{-k}}{\sqrt{2\epsilon_k}} \qquad \pi(k) = \sqrt{2\epsilon_k} \left(\frac{a_k - a^\dagger_{-k}}{2i} \right)$$

- IR state:

$$\left(\sqrt{M}\phi(x) + \frac{i}{\sqrt{M}}\pi(x) \right) |\Omega\rangle = 0$$

completely uncorrelated

$$(\alpha_k a_k + \beta_k a^\dagger_{-k})|\Omega\rangle = 0$$

$$\alpha_k = \frac{1}{2} \left(\sqrt{\frac{M}{\epsilon_k}} + \sqrt{\frac{\epsilon_k}{M}} \right) \qquad \beta_k = \frac{1}{2} \left(\sqrt{\frac{M}{\epsilon_k}} - \sqrt{\frac{\epsilon_k}{M}} \right)$$

- coarse-graining

$$e^{-iuL}\phi(k)e^{iuL} = e^{-\frac{d}{2}u}\phi(e^{-u}k)$$

$$e^{-iuL}\pi(k)e^{iuL} = e^{-\frac{d}{2}u}\pi(e^{-u}k)$$

- disentangler

$$K(u) = \frac{1}{2}\int d^d k\,[g(k,u)(\phi(k)\pi(-k) + \pi(k)\phi(-k))]$$

$$g(k,u) = \chi(u)\cdot\Gamma\left(|k|/\Lambda\right)$$

cutoff function

- variational principle:

$$E = \langle\Psi|H|\Psi\rangle = \langle\Omega|H(u_{\mathrm{IR}})|\Omega\rangle$$

$$\chi(u) = \frac{1}{2}\cdot\frac{e^{2u}}{e^{2u} + m^2/\Lambda^2}, \quad M = \sqrt{\Lambda^2 + m^2}$$

- Scale-dependent Bogoliubov transformation:

$$(\alpha_k, \beta_k) \cdot \begin{pmatrix} a_k \\ a_{-k}^\dagger \end{pmatrix} |\Omega\rangle = 0$$ IR

$$U(u)(\alpha_k, \beta_k) \cdot \begin{pmatrix} a_k \\ a_{-k}^\dagger \end{pmatrix} |\Omega\rangle = 0$$

$$(\alpha_k, \beta_k) \cdot U(u) \begin{pmatrix} a_k \\ a_{-k}^\dagger \end{pmatrix} U^\dagger(u) \cdot U(u)|\Omega\rangle = 0$$

$$(\alpha_k, \beta_k) \cdot M(u) \begin{pmatrix} a_k \\ a_{-k}^\dagger \end{pmatrix} |\Psi(u)\rangle_L = 0$$

$$(A_k, B_k) \cdot \begin{pmatrix} a_k \\ a_{-k}^\dagger \end{pmatrix} |\Psi(u)\rangle_L = 0$$ UV

Bures distance (quantum distance)

- Bures distance: $D_{\mathrm{B}}(\rho_1, \rho_2) := 2\left(1 - \mathrm{Tr}\sqrt{\rho_1^{1/2}\rho_2\rho_1^{1/2}}\right)$

- For pure states: $\rho_1 = |\psi_1\rangle\langle\psi_1| \qquad \rho_2 = |\psi_2\rangle\langle\psi_2|$

$$D_{\mathrm{B}}(\psi_1, \psi_2) = 2\left(1 - |\langle\psi_1|\psi_2\rangle|\right)$$

- For infinitesimally close state:

$$D_{\mathrm{B}}[\psi(\xi), \psi(\xi + d\xi)] = g_{ij}(\xi) d\xi_i d\xi_j$$

$$\boxed{g_{ij}(\xi) = \mathrm{Re}\,\langle\partial_i\psi(\xi)|\partial_j\psi(\xi)\rangle - \langle\partial_i\psi(\xi)|\psi\rangle\langle\psi|\partial_j\psi(\xi)\rangle}$$

- Berry gauge field $\boxed{A_i(\xi) = -i\langle\psi(\xi)|\partial_i\psi(\xi)\rangle}$

$$\langle\partial_i\psi(\xi)|[1 - \rho(\xi)]|\partial_j\psi(\xi)\rangle = g_{ij} + iF_{ij}$$

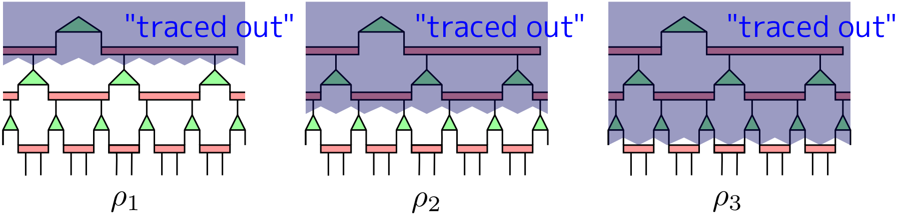
Introducing metric in MERA

- Proposal for a metric in radial direction: arXiv:1208.3469

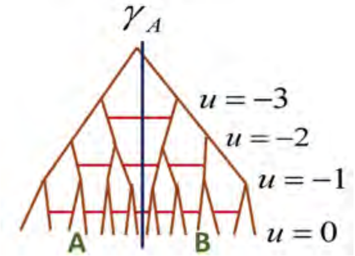
$$\boxed{g_{uu}(u)du^2 = \mathcal{N}^{-1}\left(1 - |{}_L\langle\Psi(u)|\Psi(u+du)\rangle_L|^2\right)}$$

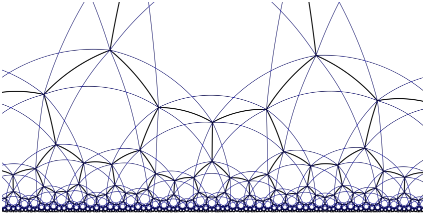
where $|\Psi(u+du)\rangle_L = e^{iLu}|\Psi(u)\rangle$ wfn in "interaction picture"

$\mathcal{N} = \mathrm{Vol.}\int_{|k|\leq\Lambda e^u} d^dk$ normalization



Motivation for the metric





$$S_A \propto L^{d-1} \sum_{u=-\infty}^{0} n(u) \cdot 2^{(d-1)u}$$

$$S_A = \frac{1}{4G_N} \cdot \frac{V_{d-1}}{\epsilon^{d-1}} \int_{u_{\mathrm{IR}}(=-\infty)}^{u_{\mathrm{UV}}(=0)} du \sqrt{g_{uu}}\, e^{(d-1)u}$$

strength of disentangler

Case study for the metric

- Relativistic free scalar:

$$ds^2 = g_{uu} du^2 + \frac{e^{2u}}{\epsilon^2} d\vec{x}^2 + g_{tt} dt^2$$

$$g_{uu}(u) = \chi(u)^2 = \frac{e^{4u}}{4(e^{2u} + m^2/\Lambda^2)^2}$$

massless limit:

$$g_{uu}(u) = \text{const.}$$

AdS metric

massive case:

$$e^{2u} = \frac{1}{\Lambda^2 z^2} - \frac{m^2}{\Lambda^2}$$

AdS soliton

$$ds^2 = \frac{dz^2}{4z^2} + \left(\frac{1}{\Lambda^2 z^2} - \frac{m^2}{\Lambda^2}\right) dx^2 + g_{tt} dt^2$$

- Flat space:

$$H = \int d^d x \, \phi(x) e^{A(-\partial^2)^{w/2}} \phi(x)$$

$$\epsilon_k \propto e^{A_1 k^w}$$

$$g_{uu}(u) = g(u)^2 \propto e^{2wu}$$

c.f. Li-Takayanagi (10)

Issues

- Large-N ? higher spin ? [cf. Swingle (12)]
- Diffeo invariance ?

- Time-component of metric g_tt ?
- Effects of interactions ?
- Einstein equation?

[cf. Faulkner-Guica-Hartman-Myers-Van Raamsdonk 13, Nozaki-Numasawa-Prudenziati-Takayanagi 13, Bhattacharya-Takayanagi 13, etc]

Advantages of AdS/MERA:

- No need for large-N
- Can define geometry for generic many-body states

MERA for Topological Phases of Matter

- Topological phases: gapped phases of matter with a lot of entanglement

  e.g. QH states (described by Chern-Simons theories)

- No classical order parameter, highly entangled quantum states of matter

- There is no adiabatic path to topologically trivial states (e.g. atomic insulator v.s. QHE)

- Topological phases: we cannot completely remove entanglement

  What happens if we try to construct disentanglers and MERA ?

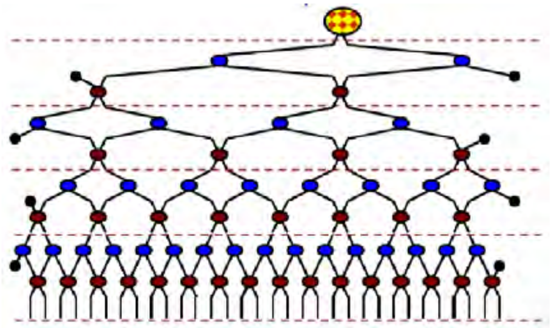
[Cf. McGreevy-Swingle 14]

MERA for topological phase

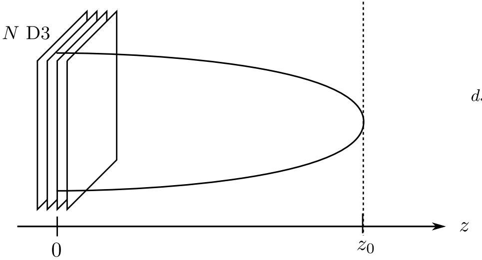
- MERA for Kitaev toric code and Levin-Wen String nets:

[Koenig, Reichardt, Vidal (08)]
[Aguado-Vidal (09)]

Topological infomation is strored in "top tensor"



Holographic dual of gapped phases

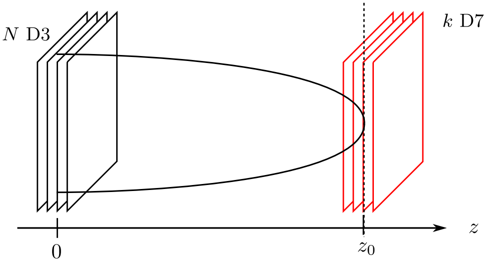


Holographic dual of pure YM in (2+1)D

[Witten (98)]

$$ds^2 = \frac{R^2}{z^2}\left(-dt^2 + f(z)dy^2 + dx_1^2 + dx_2^2 + f^{-1}(z)dz^2\right)$$

$$f(z) = 1 - (z/z_0)^4$$



Holographic dual of CS

[Fujita-Li-Ryu-Takayanagi (10)]

$$S_{\mathrm{D3}} = \frac{k}{4\pi}\int \mathrm{tr}\left(A \wedge dA + \frac{2}{3}A \wedge A \wedge A\right)$$

$$S_{top} \sim \frac{k^2}{2}\log N$$

Branes = Top tensor ?

Chern Insulators

- QHE without uniform magnetic field

- Continuum Dirac model with k-dependent mass:

$$H = \int d^2\mathbf{k}\,\psi^\dagger(\mathbf{k})\left[\boldsymbol{R}(\mathbf{k})\cdot\boldsymbol{\sigma}\right]\psi(\mathbf{k})$$

$$\boldsymbol{R}(\mathbf{k}) = (k_x,\ k_y,\ m - \boldsymbol{k}\cdot\boldsymbol{k})$$

- Two phases having different topological invariant (Hall conductance)

$$\sigma_{xy} = \frac{1}{8\pi}\int d^2k\,\boldsymbol{R}\cdot(\partial_{k_x}\boldsymbol{R}\times\partial_{k_y}\boldsymbol{R})$$

$\sigma_{xy} = +1 \quad m > 0$ "topological"

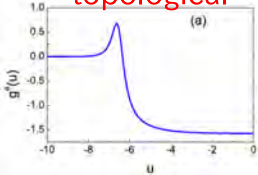
$\sigma_{xy} = 0 \quad m < 0$ "trivial"

Disentangler and Metric for Chern Insulators

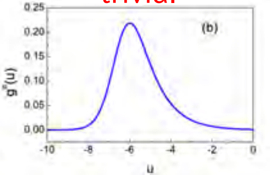
- IR state: $$|\Omega^{\text{trivial}}\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- Disentangler: $$\hat{K}(u) = i \int d^2\mathbf{k} \left[g_{\mathbf{k}}(u)\psi_1^\dagger(\mathbf{k})\psi_2(\mathbf{k}) + g_{\mathbf{k}}^*(u)\psi_1(\mathbf{k})\psi_2^\dagger(\mathbf{k}) \right]$$

topological



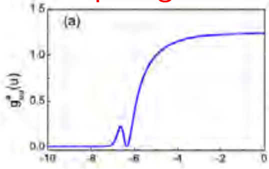
trivial



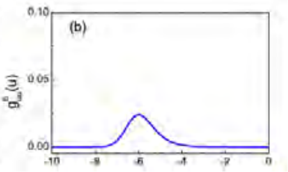
$$g_k(u) = g(u)\Gamma\left(\frac{k}{\Lambda e^u}\right)\frac{k}{\Lambda e^u}e^{-i\theta}$$

- Metric:

topological



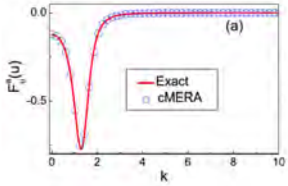
trivial



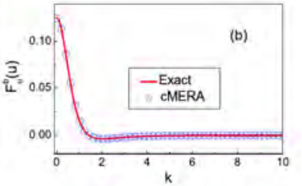
Berry flux at UV

- Small k:

topological

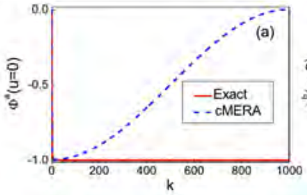


trivial

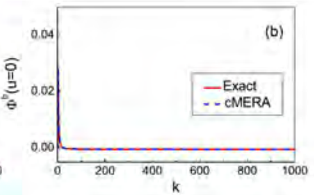


- Large k:

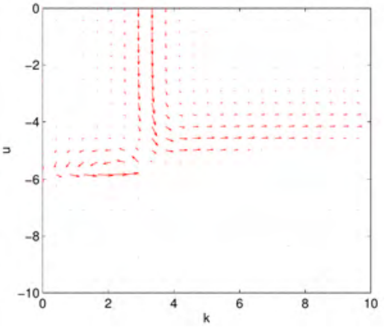
topological

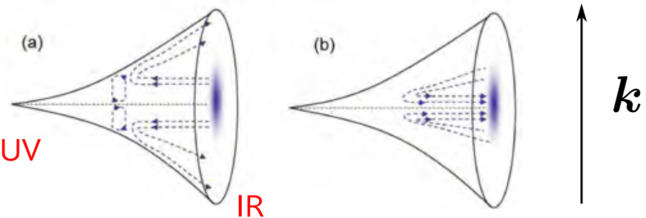


trivial



topological case

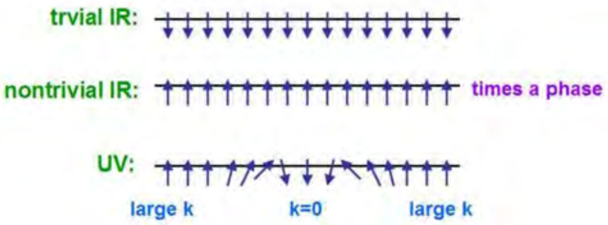


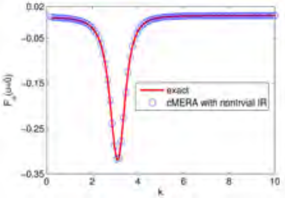


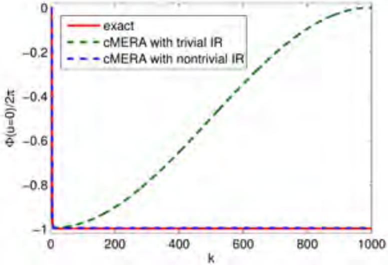
Choice of IR state

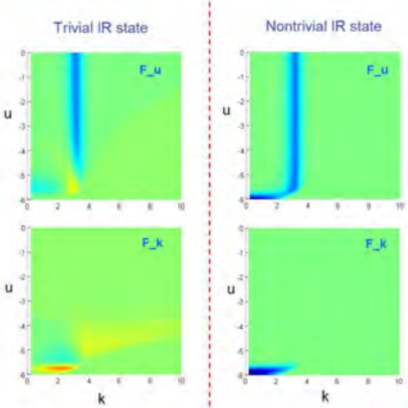
$$|\Omega^{\text{trivial}}\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

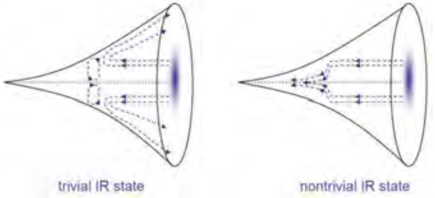
$$|\Omega^{\text{nontrivial}}\rangle = |\Psi(\mathbf{k}, u = u_{\text{IR}})\rangle = \begin{pmatrix} -e^{-i\theta} \\ 0 \end{pmatrix}$$




exact
cMERA with nontrivial IR
k


exact
cMERA with trivial IR
cMERA with nontrivial IR
Φ(u=0)/2π
k

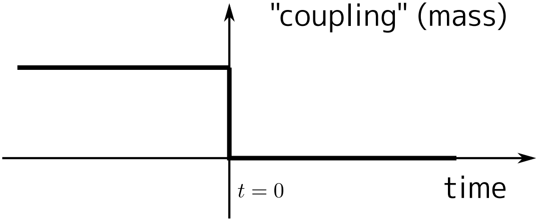

Trivial IR state
Nontrivial IR state
F_u
F_k
u
k


trivial IR state
nontrivial IR state

Quantum quench and finite T

- MERA at finite T (mixed state)

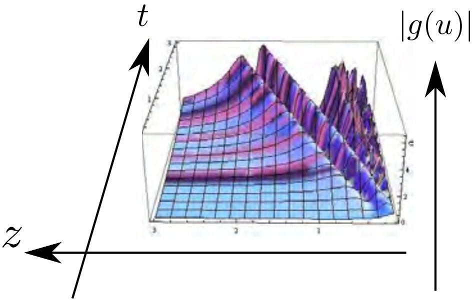
- Quantum quench (pure state)

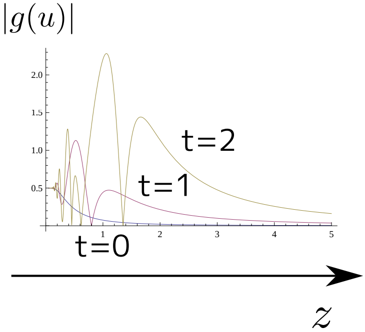


$$|\Psi(t=0)\rangle = e^{-\frac{\beta}{4}H}|B\rangle$$ "Boundary state" [Calabrese-Cardy 06]

$$|\Psi(t=0)\rangle \propto \exp\left(-\frac{1}{2}\int dk\, e^{-\frac{\beta\varepsilon_k}{2}} a_k^\dagger a_{-k}^\dagger\right)|0\rangle$$

Metric after quantum quench for 2d free boson





- Light-cone like structure

$$g_{uu} = g(u)^2 \simeq \frac{1}{4}\left[1 + \frac{a_1 k^2 \beta^2 + a_2 k^2 t^2}{\sinh^2(k\beta/2)}\right]$$

- t-linear growth of S_A

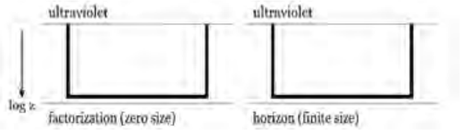
$$\Delta S_A \sim \int_{-\log\beta/\epsilon}^{0} du\,(\sqrt{g_{uu}} - 1/2) + \int_{-\infty}^{-\log\beta/\epsilon} du\sqrt{g_{uu}} \sim \frac{t}{\beta}$$

[Calabrese-Cardy (05)
Hartman-Maldacena (13)]

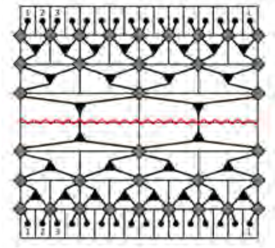
What can we say about finite T ?

- Thermofield double description:

[Swingle (12)]



[Matsueda, Molina-Vilaplana, etc]



- Concrete setup in cMERA (cf. Hartman-Maldacena, next slide):

$$
\begin{aligned}
|\Psi(0,t)\rangle_{th} &= \mathcal{N} \cdot e^{-it(H_1+H_2)} \cdot \prod_k \sum_{n_k=0}^{\infty} e^{-\beta\epsilon_k n_k/2} |n_k\rangle_1 |n_k\rangle_2 \\
&= \mathcal{N} \cdot \exp\left(\int dk e^{-\frac{\beta\epsilon_k}{2}} e^{-2i\epsilon_k t} a_k^\dagger \tilde{a}_k^\dagger \right) |0\rangle|\tilde{0}\rangle.
\end{aligned}
$$

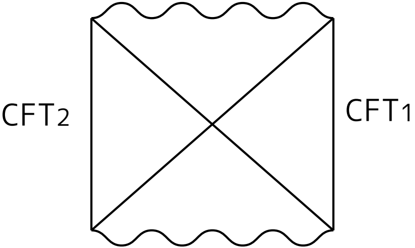
- Can use the same disentangler for quantum quench

$$
\tilde{a}_k \to a_{-k}, \quad |0\rangle|\tilde{0}\rangle \to |0\rangle.
$$

$$
|\Psi(0, t=0)\rangle_{th} = P e^{-i\int_{u_{IR}}^{0} \hat{K}(s)ds} \otimes P e^{-i\int_{u_{IR}}^{0} \hat{\tilde{K}}(\tilde{s})d\tilde{s}} |\Omega(\beta)\rangle.
$$

same metric as quench

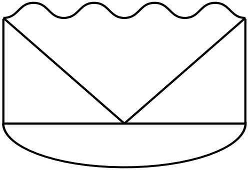
Hartman-Maldacena



$$ds^2 = -\frac{1-z^2/z_H^2}{z^2}d\tau^2 + \frac{dz^2}{z^2(1-z^2/z_H^2)} + \frac{dx^2}{z^2}$$

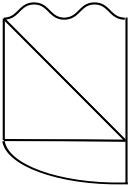
[Israel (76), Maldacena]

$$|\Psi\rangle = \sum_n e^{-\beta E_n/2}|n\rangle_1|n\rangle_2$$



thermofield double

$$|\Psi\rangle = \sum_n e^{iE_n t}e^{-\beta E_n/2}|n\rangle_1|n\rangle_2$$



quantum quench

$$|\Psi\rangle = \sum_n e^{iE_n t}e^{-\beta E_n/2}|n\rangle$$

Summary

- Proposal for definition of of bulk metric and gauge field
  for MERA representation of quantum states

- Behaviors expected from AdS/CFT

- Classical phases of matter <-- group theory (symmetry breaking)
  Quantum phases of matter <-- geometry (entanglement)
  Topological phases <--> D-branes (non-trivial IR)