

# A Holographic Model of the Kondo Effect (Part I)

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Quantum Field Theory, String Theory, and Condensed Matter Physics  
Kolymbari, Greece  
September 1, 2014

# Credits

Based on 1310.3271

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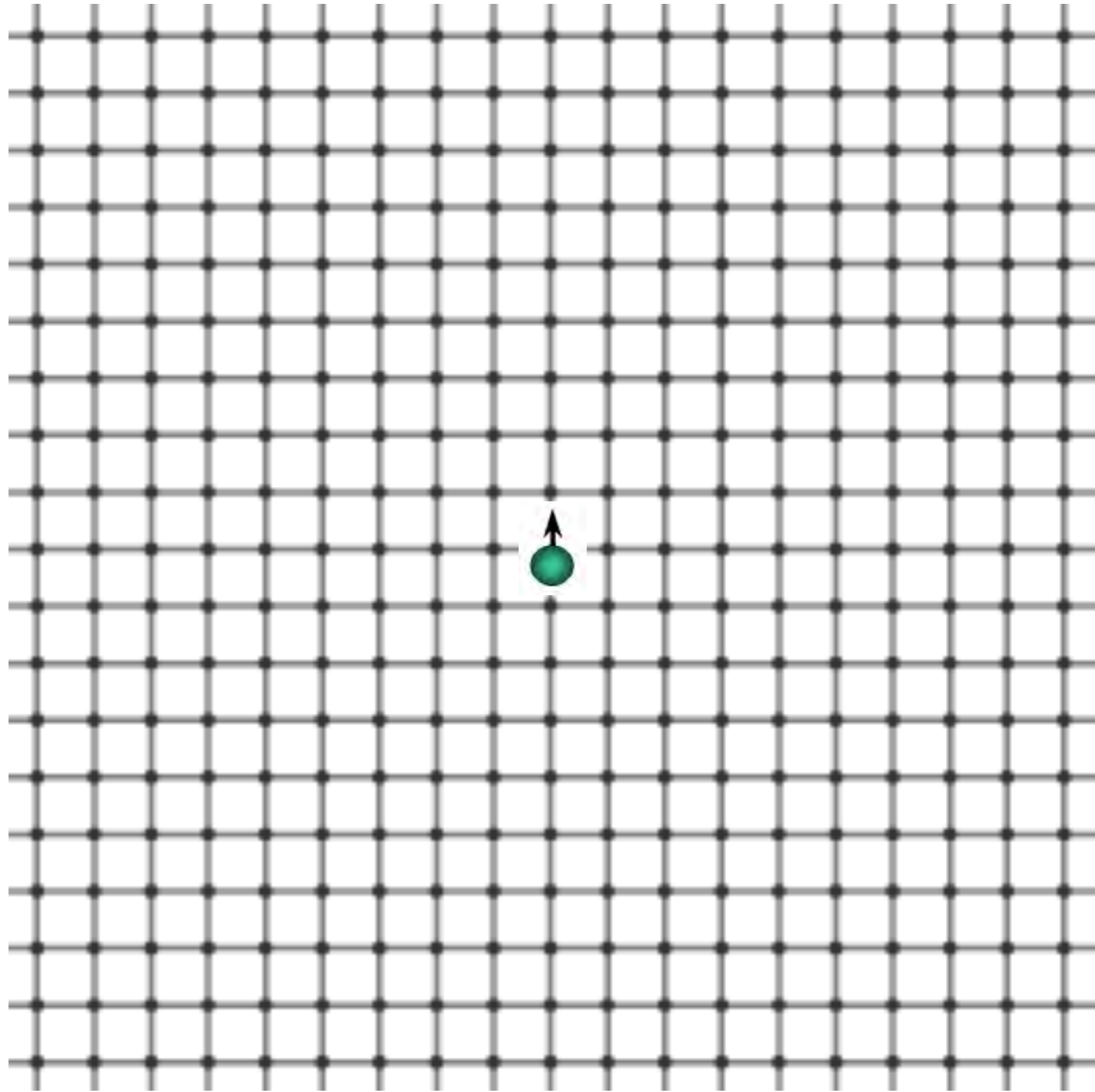
National Center for Theoretical Sciences, Taiwan

# Outline:

- The Kondo Effect
- The CFT Approach
- Top-Down Holographic Model
- Bottom-Up Holographic Model
- Summary and Outlook

# The Kondo Effect

The **screening**  
of a **magnetic moment**  
by **conduction electrons**  
at low **temperatures**



$$\vec{\mu} \propto g \vec{S}$$

# The Kondo Hamiltonian

$$H_K = \sum_{k,\sigma} \varepsilon(k) c_{k\sigma}^\dagger c_{k\sigma} + g_K \vec{S} \cdot \sum_{k\sigma k'\sigma'} c_{k\sigma}^\dagger \frac{1}{2} \vec{\tau}_{\sigma\sigma'} c_{k'\sigma'}$$

$$c_{k\sigma}^\dagger, c_{k\sigma}$$

Conduction electrons

$$\sigma = \uparrow, \downarrow$$

Spin  $SU(2)$

$$c_{k\sigma} \rightarrow e^{i\alpha} c_{k\sigma}$$

Charge  $U(1)$

$$\varepsilon(k) = \frac{k^2}{2m} - \varepsilon_F$$

Dispersion relation

# The Kondo Hamiltonian

$$H_K = \sum_{k,\sigma} \varepsilon(k) c_{k\sigma}^\dagger c_{k\sigma} + g_K \vec{S} \cdot \sum_{k\sigma k'\sigma'} c_{k\sigma}^\dagger \frac{1}{2} \vec{\tau}_{\sigma\sigma'} c_{k'\sigma'}$$

$\vec{S}$

Spin of magnetic moment

$\vec{\tau}$

Pauli matrices

$g_K$

Kondo coupling

# Running of the Coupling

$$\beta_{g_K} \propto -g_K^2 + \mathcal{O}(g_K^3)$$

Asymptotic freedom!

UV

$$g_K \rightarrow 0$$

The Kondo Temperature

$$T_K \sim \Lambda_{\text{QCD}}$$



## Running of the Coupling

$$\beta_{g_K} \propto -g_K^2 + \mathcal{O}(g_K^3)$$

At low energy, the coupling diverges!

IR

$$g_K \rightarrow \infty$$

The Kondo Problem

What is the ground state?

# Solutions of the Kondo Problem

Numerical RG (Wilson 1975)

Fermi liquid description (Nozières 1975)

Bethe Ansatz/Integrability  
(Andrei, Wiegmann, Tsvelick, Destri, ... 1980s)

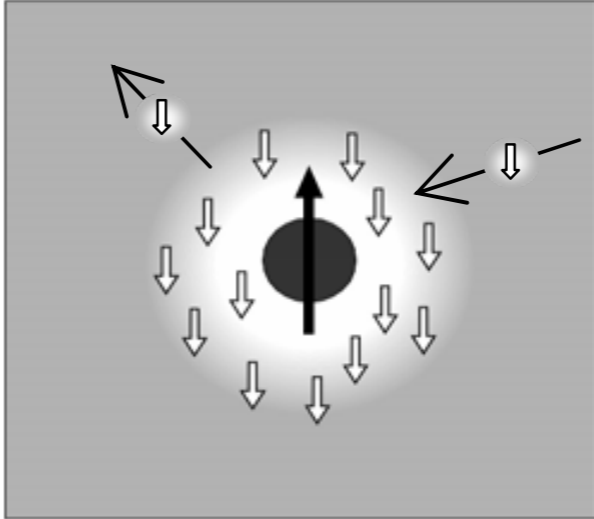
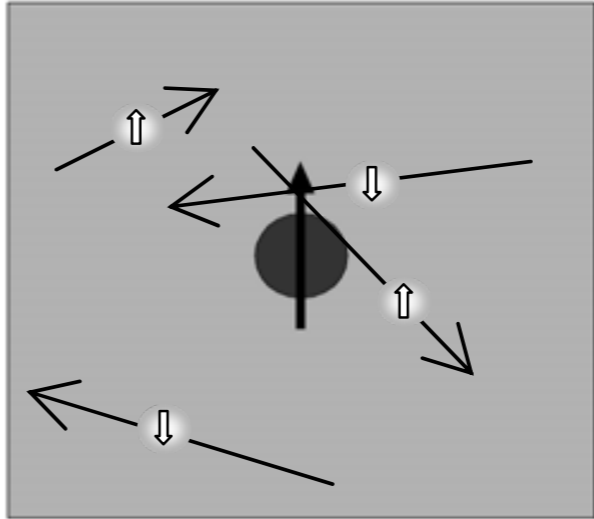
Large-N expansion  
(Anderson, Read, Newns, Doniach, Coleman, ... 1970-80s)

Quantum Monte Carlo  
(Hirsch, Fye, Gubernatis, Scalapino, ... 1980s)

Conformal Field Theory (CFT)  
(Affleck and Ludwig 1990s)

UV

Fermi liquid + decoupled spin



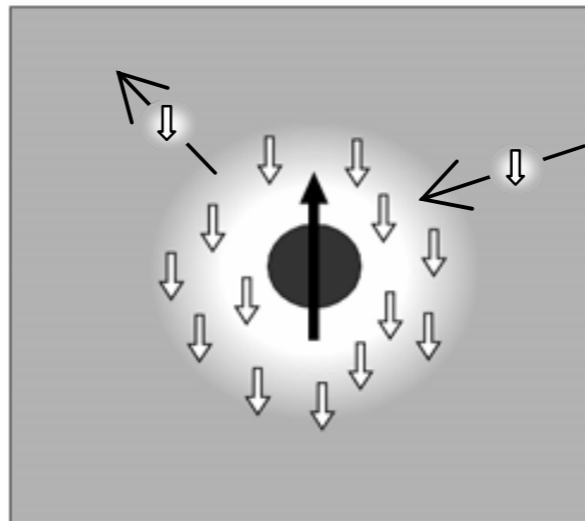
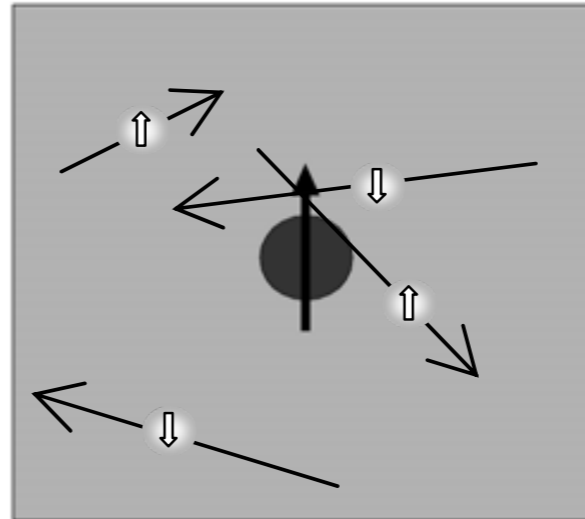
size  $\propto 1/T_K$

IR

“Kondo screening cloud”

UV

Fermi liquid + decoupled spin



**size**  $\propto 1/T_K$

IR

On average, the NET spin of the Kondo screening cloud equals that of a single electron

UV

Fermi liquid + decoupled spin

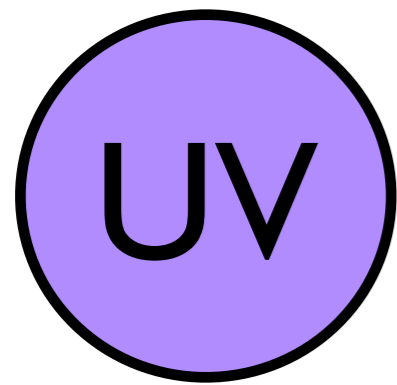
The Kondo screening cloud spin binds with the impurity spin

Anti-symmetric singlet of  $SU(2)$

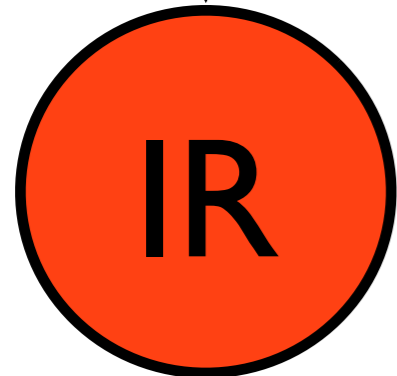
$$\frac{1}{\sqrt{2}} (|\uparrow_i \downarrow_e\rangle - |\downarrow_i \uparrow_e\rangle)$$

IR

“Kondo singlet”



Fermi liquid + decoupled spin



Fermi liquid

+ NO magnetic moment

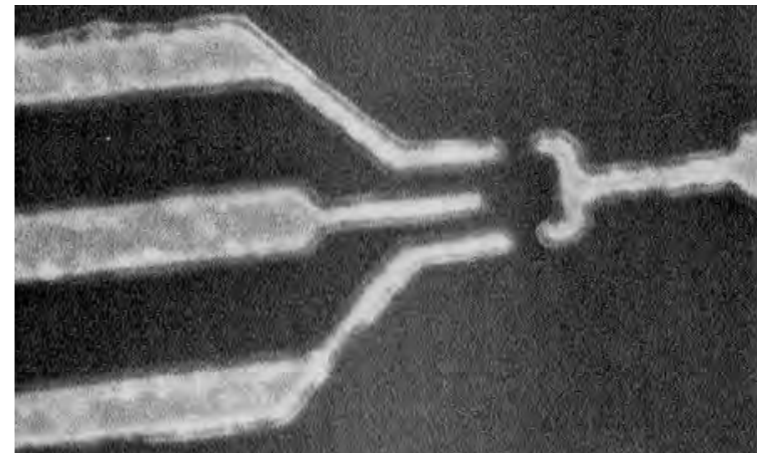
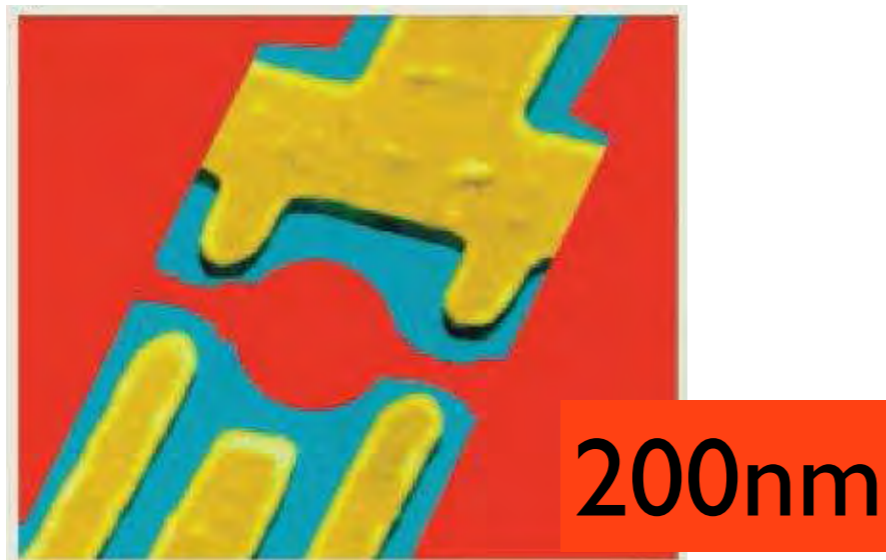
+ electrons EXCLUDED  
from impurity location

# Kondo Effect in Many Systems

## Alloys

Cu, Ag, Au, Mg, Zn, ... doped with Cr, Fe, Mo, Mn, Re, Os, ...

## Quantum dots



Goldhaber-Gordon, et al., *Nature* 391 (1998), 156-159.

Cronenwett, et al., *Science* 281 (1998), no. 5376, 540-544.

# Generalizations

Enhance the spin group

$$SU(2) \rightarrow SU(N)$$

Representation of impurity spin

$$s_{\text{imp}} = 1/2 \longrightarrow R_{\text{imp}}$$

Multiple “channels” or “flavors”

$$c \longrightarrow c^{\alpha} \quad \alpha = 1, \dots, k$$

$$U(1) \times SU(k)$$



# Generalizations

Kondo model specified by

$$N, k, R_{\text{imp}}$$

Apply the techniques mentioned above...

IR fixed point:

NOT always  
a fermi liquid

“Non-Fermi liquids”

# Open Problems

## Entanglement Entropy

Affleck, Laflorencie, Sørensen 0906.1809

## Quantum Quenches

Latta et al. 1102.3982

## Multiple Impurities

Kondo:

Form singlets with electrons

$$\vec{S}_i \cdot \vec{S}_j$$

Form singlets with each other

Competition between these can produce a

**QUANTUM PHASE TRANSITION**

# Open Problems

## Entanglement Entropy

Affleck, Laflorencie, Sørensen 0906.1809

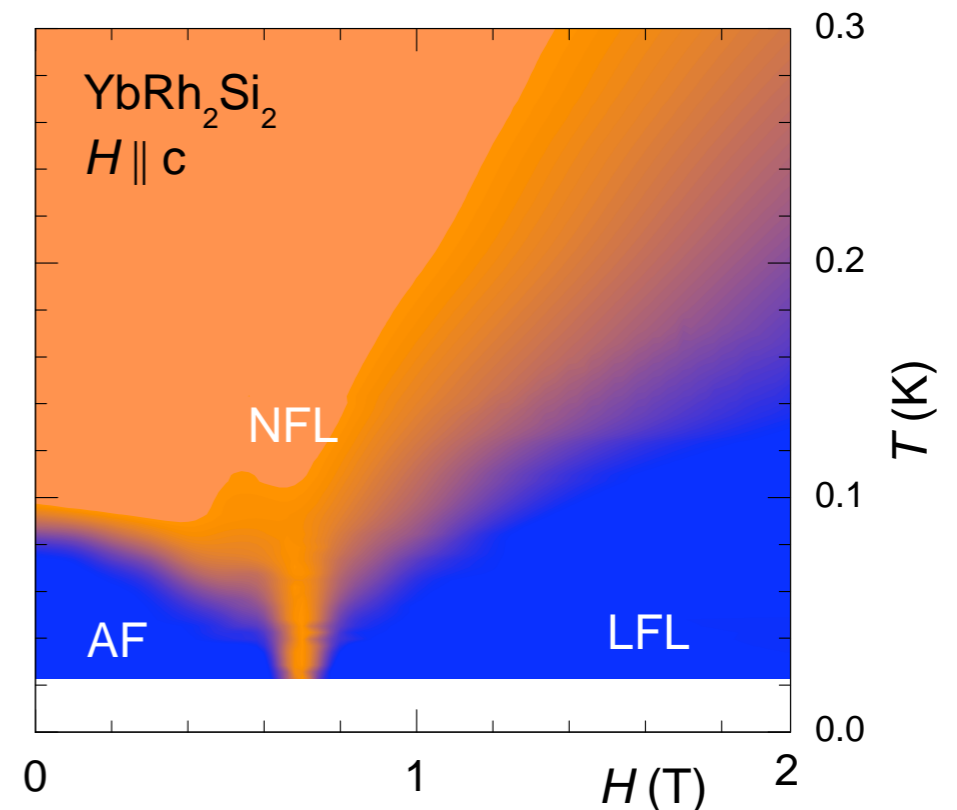
## Quantum Quenches

Latta et al. 1102.3982

## Multiple Impurities

Heavy fermion compounds

Kondo lattice



# Open Problems

## Entanglement Entropy

Affleck, Laflorencie, Sørensen 0906.1809

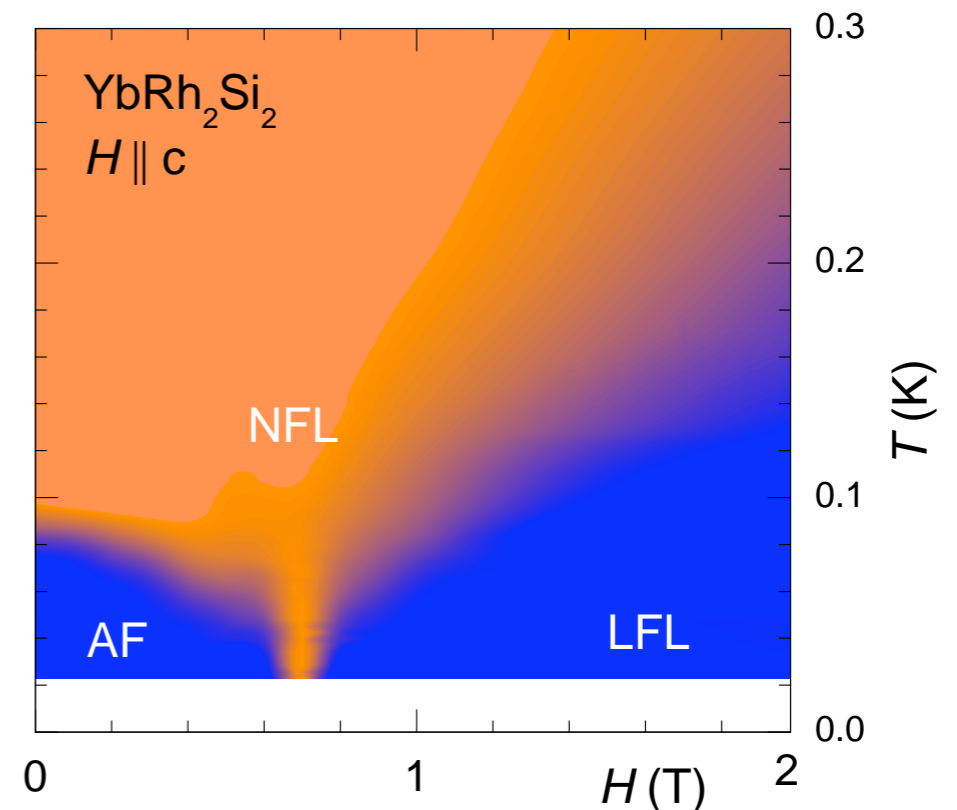
Quantum Quench

Let's try AdS/CFT!

Multiple impurities

Heavy fermion compounds

Kondo lattice



# GOAL

Find a holographic description  
of the  
Kondo Effect

# GOAL

Find a holographic description

**Single Impurity ONLY**

Kondo Effect

# Solutions of the Kondo Problem

Numerical RG (Wilson 1975)

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Bethe Ansatz/Integrability  
(Andrei, Wiegmann, Tsvetick, Destri, ... 1980s)

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(Affleck and Ludwig 1990s)

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# CFT Approach to the Kondo Effect

Affleck and Ludwig 1990s

Reduction to one spatial dimension

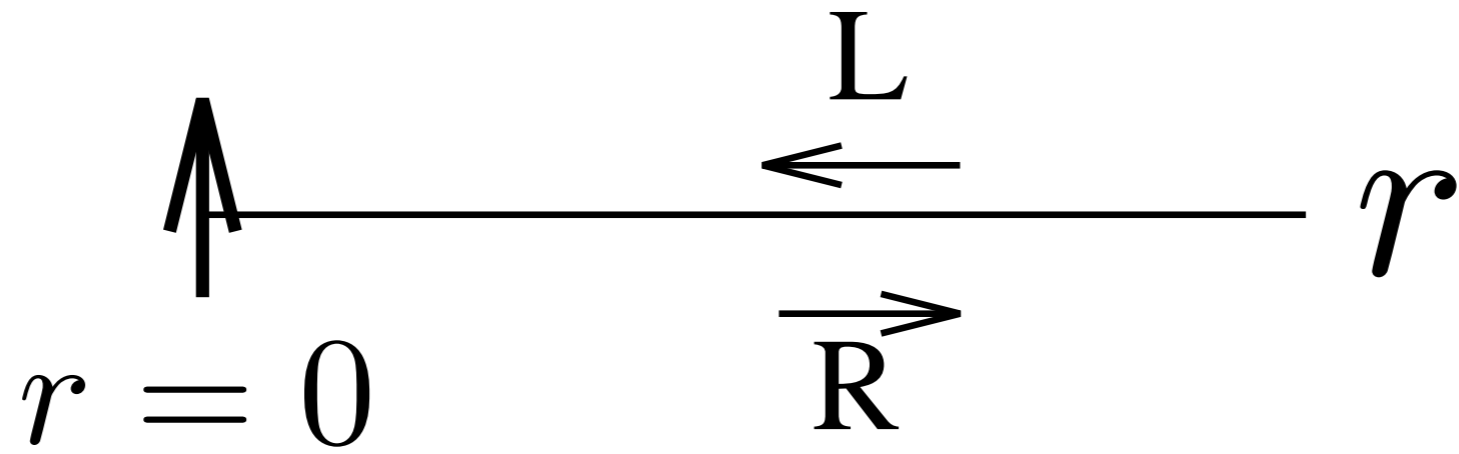
Kondo interaction preserves spherical symmetry

$$g_K \delta^3(\vec{x}) \vec{S} \cdot c^\dagger(\vec{x}) \frac{\vec{\tau}}{2} c(\vec{x})$$

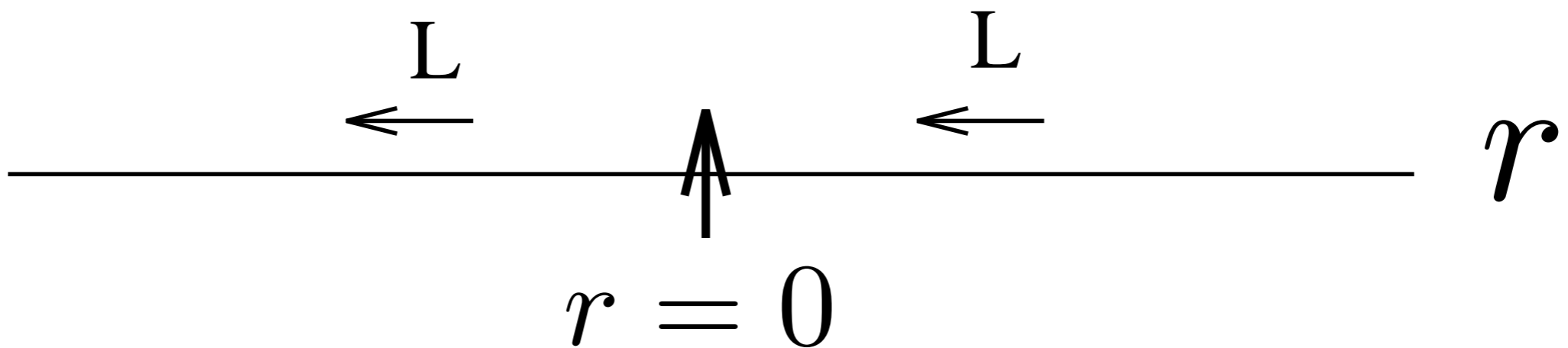
restrict to s-wave

restrict to momenta near  $k_F$

$$c(\vec{x}) \approx \frac{1}{r} \left[ e^{-ik_F r} \psi_L(r) - e^{+ik_F r} \psi_R(r) \right]$$



$$\psi_L(-r) \equiv \psi_R(+r)$$



# CFT Approach to the Kondo Effect

$$H_K = \frac{v_F}{2\pi} \int_{-\infty}^{+\infty} dr \left[ \psi_L^\dagger i\partial_r \psi_L + \delta(r) \tilde{g}_K \vec{S} \cdot \psi_L^\dagger \vec{\tau} \psi_L \right]$$

$$\tilde{g}_K \equiv \frac{k_F^2}{2\pi^2 v_F} \times g_K$$

RELATIVISTIC chiral fermions

$v_F$  = “speed of light”

**chiral CFT!**

Spin  $SU(N)$

$k \geq 1$

$$J = \psi_L^\dagger \psi_L$$

$U(1)$

$$\vec{J} = \psi_L^\dagger \vec{\tau} \psi_L$$

$SU(N)$

$$J^A = \psi_L^\dagger t^A \psi_L$$

$SU(k)$

$$z \equiv \tau + ir$$

$$J^A(z) = \sum_{n \in \mathbb{Z}} z^{-n-1} J_n^A$$

$$[J_n^A, J_m^B] = if^{ABC} J_{n+m}^C + N \frac{n}{2} \delta^{AB} \delta_{n,-m}$$

$SU(k)_N$  Kac-Moody Algebra

**N counts net number of chiral fermions**

# CFT Approach to the Kondo Effect

$$H_K = \frac{v_F}{2\pi} \int_{-\infty}^{+\infty} dr \left[ \psi_L^\dagger i \partial_r \psi_L + \delta(r) \tilde{g}_K \vec{S} \cdot \psi_L^\dagger \vec{\tau} \psi_L \right]$$

**Full symmetry:**

$(1 + 1)d$  chiral conformal symmetry

$$SU(N)_k \times SU(k)_N \times U(1)_{kN}$$

# CFT Approach to the Kondo Effect

$$H_K = \frac{v_F}{2\pi} \int_{-\infty}^{+\infty} dr \left[ \psi_L^\dagger i \partial_r \psi_L + \delta(r) \tilde{g}_K \vec{S} \cdot \psi_L^\dagger \vec{\tau} \psi_L \right]$$

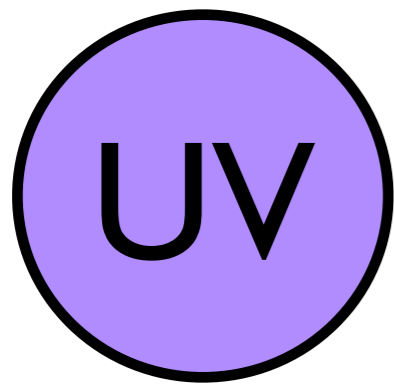
$$J = \psi_L^\dagger \psi_L \quad U(1)$$

$$\vec{J} = \psi_L^\dagger \vec{\tau} \psi_L \quad SU(N)$$

$$J^A = \psi_L^\dagger t^A \psi_L \quad SU(k)$$

Kondo coupling:  $\vec{S} \cdot \vec{J}$

marginal

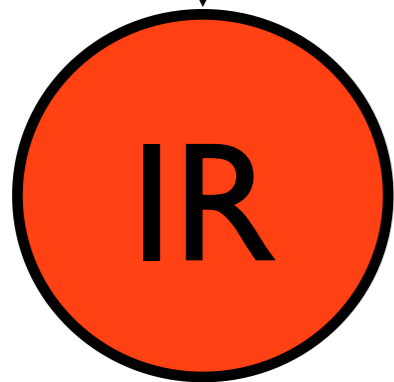


$$SU(N)_k \times SU(k)_N \times U(1)_{Nk}$$

Eigenstates are representations of the Kac-Moody algebra

Determine how representations re-arrange between UV and IR

$$R_{\text{highest weight}}^{UV} \otimes R_{\text{imp}} = R_{\text{highest weight}}^{IR}$$



$$SU(N)_k \times SU(k)_N \times U(1)_{Nk}$$



# CFT Approach to the Kondo Effect

Take-Away Messages

Central role of the  
Kac-Moody Algebra

Kondo coupling:  $\vec{S} \cdot \vec{J}$

# Outline:

- The Kondo Effect
- The CFT Approach
- Top-Down Holographic Model
- Bottom-Up Holographic Model
- Summary and Outlook

# GOAL

Find a holographic description  
of the  
Kondo Effect

# STRATEGY

Follow the CFT approach

Reproduce the Symmetries

Reproduce the Kondo coupling  $\vec{S} \cdot \vec{J}$

What classical action do we write  
on the gravity side of the correspondence?

# How do we describe holographically...

- ① The chiral fermions?
- ② The impurity?
- ③ The Kondo coupling?

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

Open strings

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

**3-3** and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

CFT with holographic dual



# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

Decouple



# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

(1+1)-dimensional  
chiral fermions

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

the impurity

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

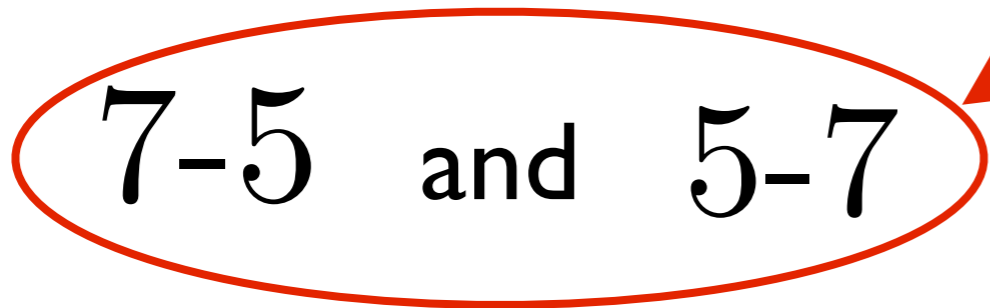
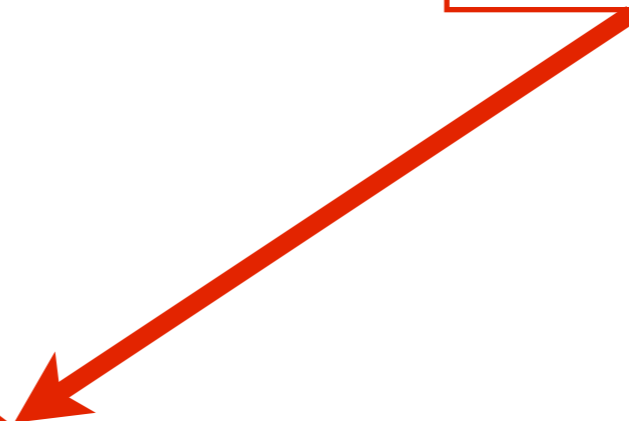
3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

Kondo interaction



# The D3-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						

3-3 strings

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

$$\int_{S^5} F_5 \propto N_c$$

$$F_5 = dC_4$$

# The D3-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						

3-3 strings

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=

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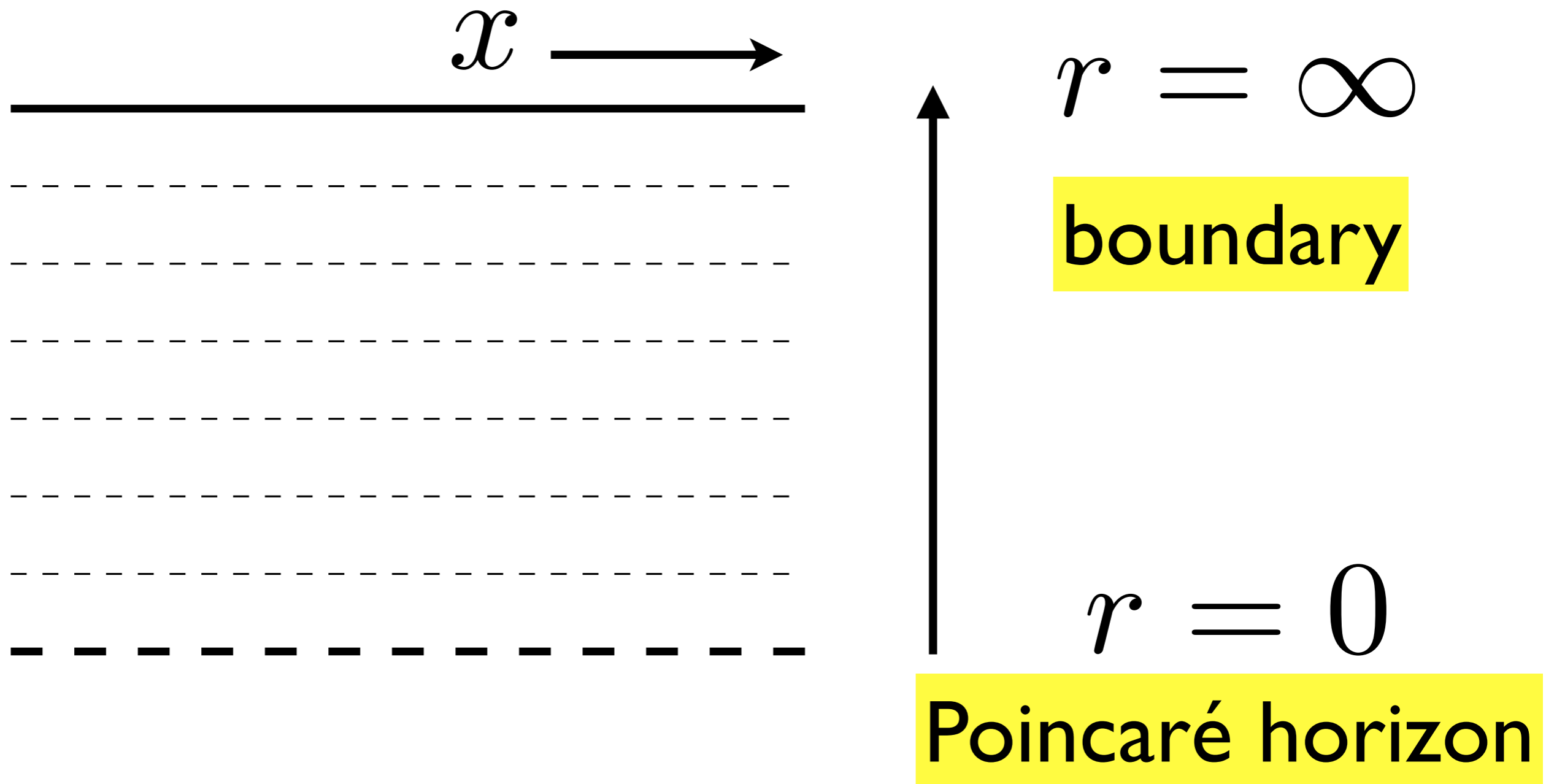
$AdS_5 \times S^5$

Our Kondo model will have TWO coupling constants

't Hooft and Kondo

# Anti-de Sitter Space

$$ds^2 = \frac{dr^2}{r^2} + r^2 (-dt^2 + dx^2 + dy^2 + dz^2)$$



# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

Decouple





# Probe Limit

$$N_7/N_c \rightarrow 0 \text{ and } N_5/N_c \rightarrow 0$$

$U(N_7) \times U(N_5)$  becomes a global symmetry

Total symmetry:

$$\underbrace{SU(N_c)}_{\text{gauged}} \times \underbrace{U(N_7) \times U(N_5)}_{\text{global}}$$

(plus R-symmetry)

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

(1+1)-dimensional  
chiral fermions

# The D7-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X

Skenderis, Taylor hep-th/0204054

Harvey and Royston 0709.1482, 0804.2854

Buchbinder, Gomis, Passerini 0710.5170

(1+1)-dimensional chiral fermions  $\psi_L$

$$SU(N_c) \times U(N_7) \times U(N_5)$$

$$N_c \quad \overline{N}_7 \quad \text{singlet}$$

# The D7-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X

Skenderis, Taylor hep-th/0204054

Harvey and Royston 0709.1482, 0804.2854

Buchbinder, Gomis, Passerini 0710.5170

(1+1)-dimensional chiral fermions  $\psi_L$

Kac-Moody algebra

$$SU(N_c)_{N_7} \times SU(N_7)_{N_c} \times U(1)_{N_c N_7}$$

# The D7-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X

## Differences from Kondo

(1+1)-dimensional chiral fermions  $\psi_L$

Do not come from reduction from (3+1) dimensions

Genuinely relativistic

# The D7-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X

(1+1)-dimensional chiral fermions  $\psi_L$

Differences from Kondo

$SU(N_c)$  is gauged!

$$\vec{J} = \psi_L^\dagger \vec{\tau} \psi_L$$

# The D7-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X

$SU(N_c)$  is gauged!



Gauge Anomaly!



Harvey and Royston 0709.1482, 0804.2854

Buchbinder, Gomis, Passerini 0710.5170

# Probe Limit

$$N_7/N_c \rightarrow 0$$

In the probe limit, the gauge anomaly is suppressed...

$$SU(N_c)_{N_7} \rightarrow SU(N_c)$$

... but the global anomalies are not.

$$SU(N_7)_{N_c} \times U(1)_{N_c N_7} \rightarrow SU(N_7)_{N_c} \times U(1)_{N_c N_7}$$



$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\psi_L$

$=$

Probe D7-branes

$AdS_3 \times S^5$

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

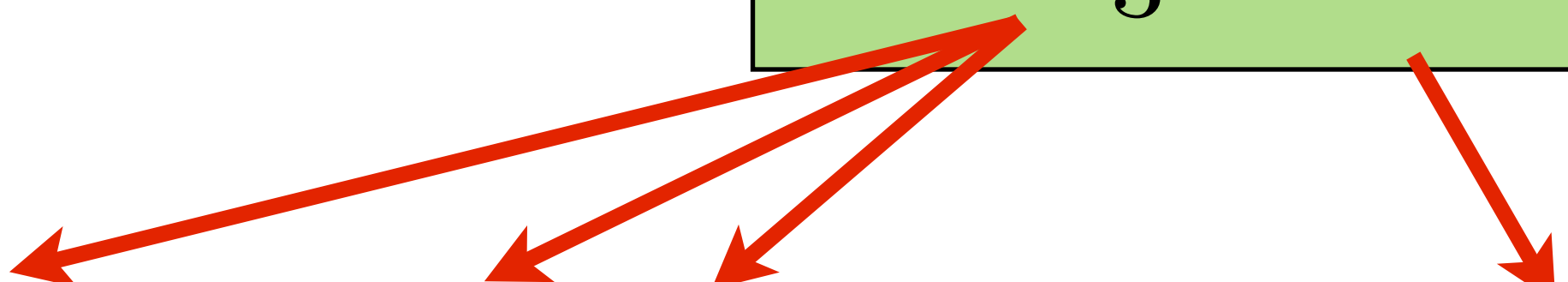
$AdS_5 \times S^5$

Probe  $\psi_L$

$=$

Probe D7-branes

$AdS_3 \times S^5$


$$ds^2 = \frac{dr^2}{r^2} + r^2 (-dt^2 + dx^2 + dy^2 + dz^2) + ds_{S^5}^2$$

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\psi_L$

$=$

Probe D7-branes

$AdS_3 \times S^5$

$U(N_7)$  Current  $J$

$=$

$U(N_7)$  Gauge field  $A$

Current  $J$  = Gauge field  $A$

Kac-Moody Algebra = Chern-Simons Gauge Field

rank and level of algebra = rank and level of gauge field

Gukov, Martinec, Moore, Strominger  
hep-th/0403225

Kraus and Larsen  
hep-th/0607138

Probe D7-branes along  $AdS_3 \times S^5$

$$S_{D7} = +\frac{1}{2}T_{D7}(2\pi\alpha')^2 \int P[C_4] \wedge \text{tr} F \wedge F + \dots$$

$$= -\frac{1}{2}T_{D7}(2\pi\alpha')^2 \int P[F_5] \wedge \text{tr} \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right) + \dots$$

$$= -\frac{N_c}{4\pi} \int_{AdS_3} \text{tr} \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right) + \dots$$

$U(N_7)_{N_c}$  Chern-Simons gauge field

# Answer #1

The chiral fermions:

Chern-Simons Gauge Field in  $AdS_3$

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

the impurity

# The D5-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_5$ D5	X				X	X	X	X	X	

Skenderis, Taylor hep-th/0204054

Camino, Paredes, Ramallo hep-th/0104082

Gomis and Passerini hep-th/0604007

(0+1)-dimensional fermions  $\chi$

$$SU(N_c) \times U(N_7) \times U(N_5)$$

$N_c$

singlet

$\overline{N}_5$



# The D5-branes

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_5$ D5	X				X	X	X	X	X	

$SU(N_c)$  is “spin”

$$\vec{S} = \chi^\dagger \vec{\tau} \chi$$

“slave fermions”

“Abrikosov pseudo-fermions”

Abrikosov, **Physics** 2, p.5 (1965)

$$N_5 = 1$$

Integrate out  $\chi$

$$\text{Det} (\not{D}) = \text{Tr}_R P \exp \left[ i \int dt A_t \right]$$

$$R = \left. \begin{array}{c} \square \\ \square \\ \vdots \\ \square \\ \square \\ \square \end{array} \right\} Q = \chi^\dagger \chi$$

$$U(N_5) = U(1) \text{ charge}$$

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\chi$

$=$

Probe D5-branes

$AdS_2 \times S^4$

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\chi$

$=$

Probe D5-branes

$AdS_2 \times S^4$


$$ds^2 = \frac{dr^2}{r^2} + r^2 (-dt^2 + dx^2 + dy^2 + dz^2) + ds_{S^5}^2$$

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\chi$

$=$

Probe D5-branes

$AdS_2 \times S^4$

$U(N_5)$  Current  $J$

$=$

$U(N_5)$  Gauge field  $a$

$Q$

$=$

Electric flux

Probe D5-brane along  $AdS_2 \times S^4$

Camino, Paredes, Ramallo hep-th/0104082

Dissolve  $Q$  strings into the D5-brane

$AdS_2$  electric field  $f_{rt}$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} = Q = \chi^\dagger \chi$$

# Answer #2

The impurity:

Yang-Mills Gauge Field in  $AdS_2$

# Top-Down Model

	0	1	2	3	4	5	6	7	8	9
$N_c$ D3	X	X	X	X						
$N_7$ D7	X	X			X	X	X	X	X	X
$N_5$ D5	X				X	X	X	X	X	

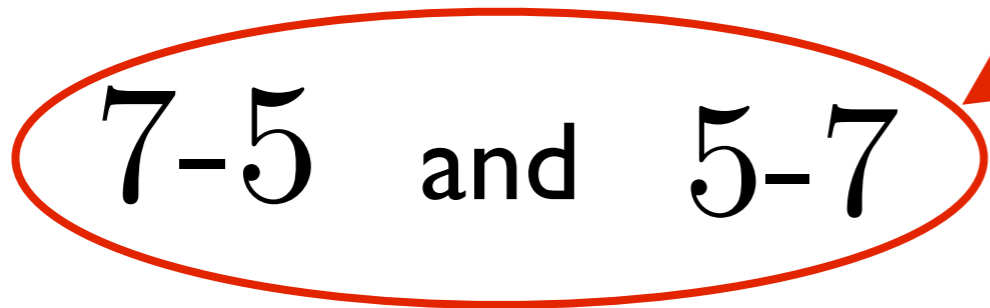
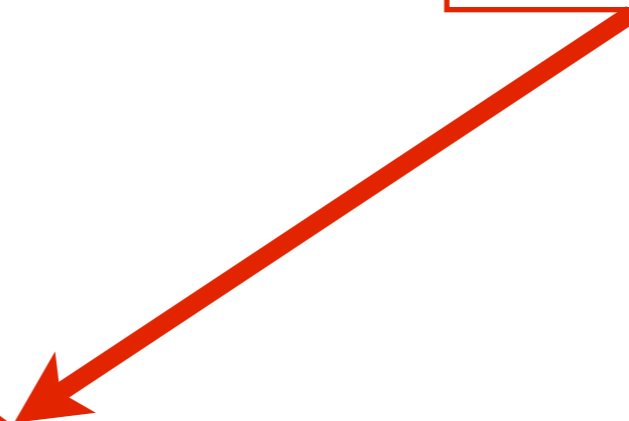
3-3 and 5-5 and 7-7

3-7 and 7-3

3-5 and 5-3

7-5 and 5-7

Kondo interaction





# The Kondo Interaction

	0	1	2	3	4	5	6	7	8	9
$N_5$ D5	X				X	X	X	X	X	
$N_7$ D7	X	X			X	X	X	X	X	X

Complex scalar!

$$\begin{array}{ccc}
 SU(N_c) \times U(N_7) \times U(N_5) & & \\
 \text{singlet} & \overline{N}_7 & N_5
 \end{array}$$

$$\mathcal{O} \equiv \psi_L^\dagger \chi$$

# The Kondo Interaction

$SU(N_c)$  is “spin”

$$\vec{J} = \psi_L^\dagger \vec{\tau} \psi_L$$

$$\vec{S} = \chi^\dagger \vec{\tau} \chi$$

$$\vec{S} \cdot \vec{J} = \chi^\dagger \vec{\tau} \chi \cdot \psi_L^\dagger \vec{\tau} \psi_L$$

$$\vec{\tau}_{ij} \cdot \vec{\tau}_{kl} = \delta_{il} \delta_{jk} - \frac{1}{N_c} \delta_{ij} \delta_{kl}$$

$$\vec{S} \cdot \vec{J} = |\psi_L^\dagger \chi|^2 + \mathcal{O}(1/N_c)$$

“double trace”

$\mathcal{N} = 4$  SYM

$N_c \rightarrow \infty$

$\lambda \rightarrow \infty$

$=$

Type IIB Supergravity

$AdS_5 \times S^5$

Probe  $\psi_L$

$=$

Probe D7-branes

$AdS_3 \times S^5$

Probe  $\chi$

$=$

Probe D5-branes

$AdS_2 \times S^4$

$\mathcal{O} \equiv \psi_L^\dagger \chi$

$=$

Bi-fundamental scalar

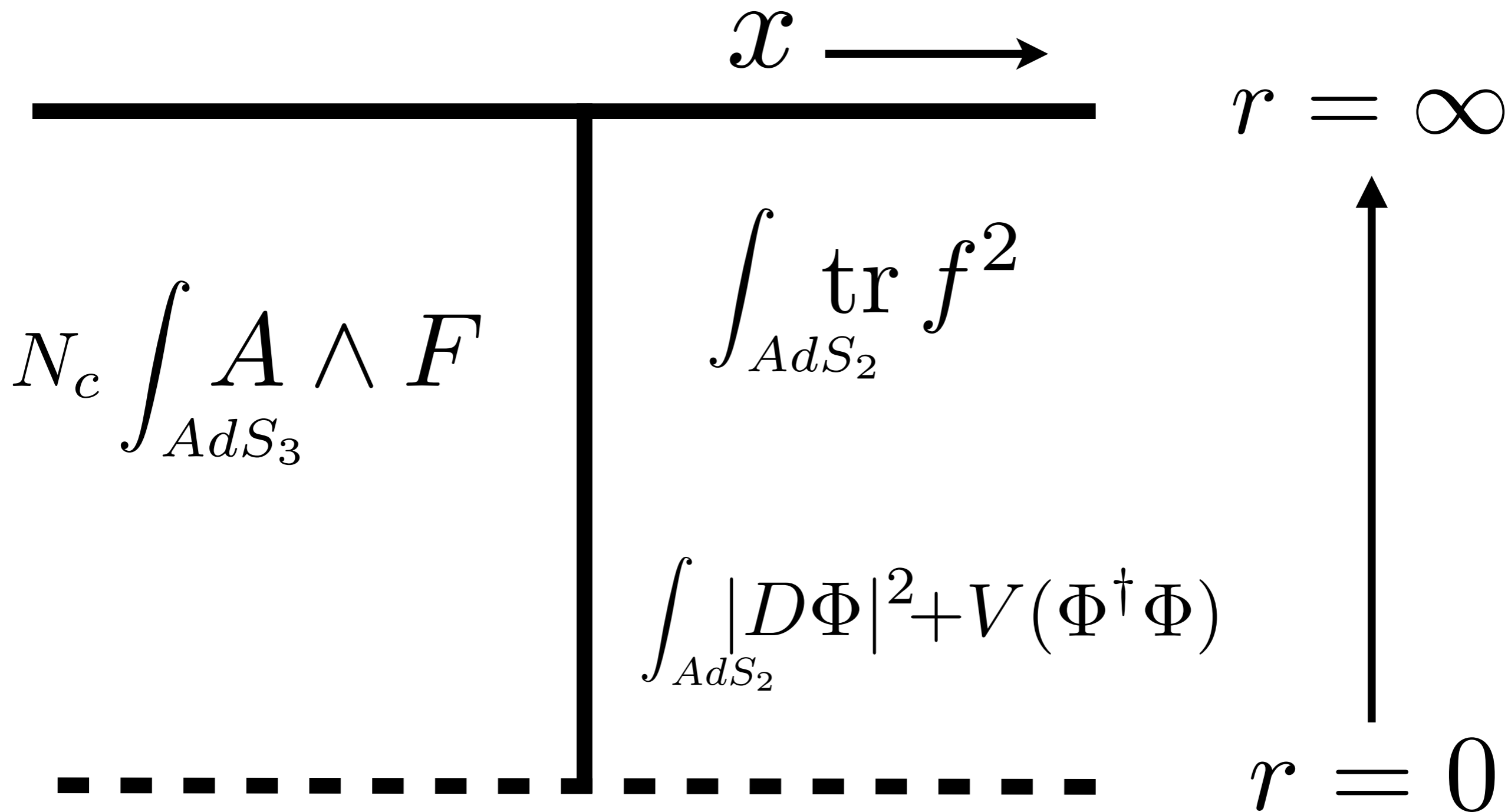
$AdS_2 \times S^4$

# Answer #3

The Kondo interaction:

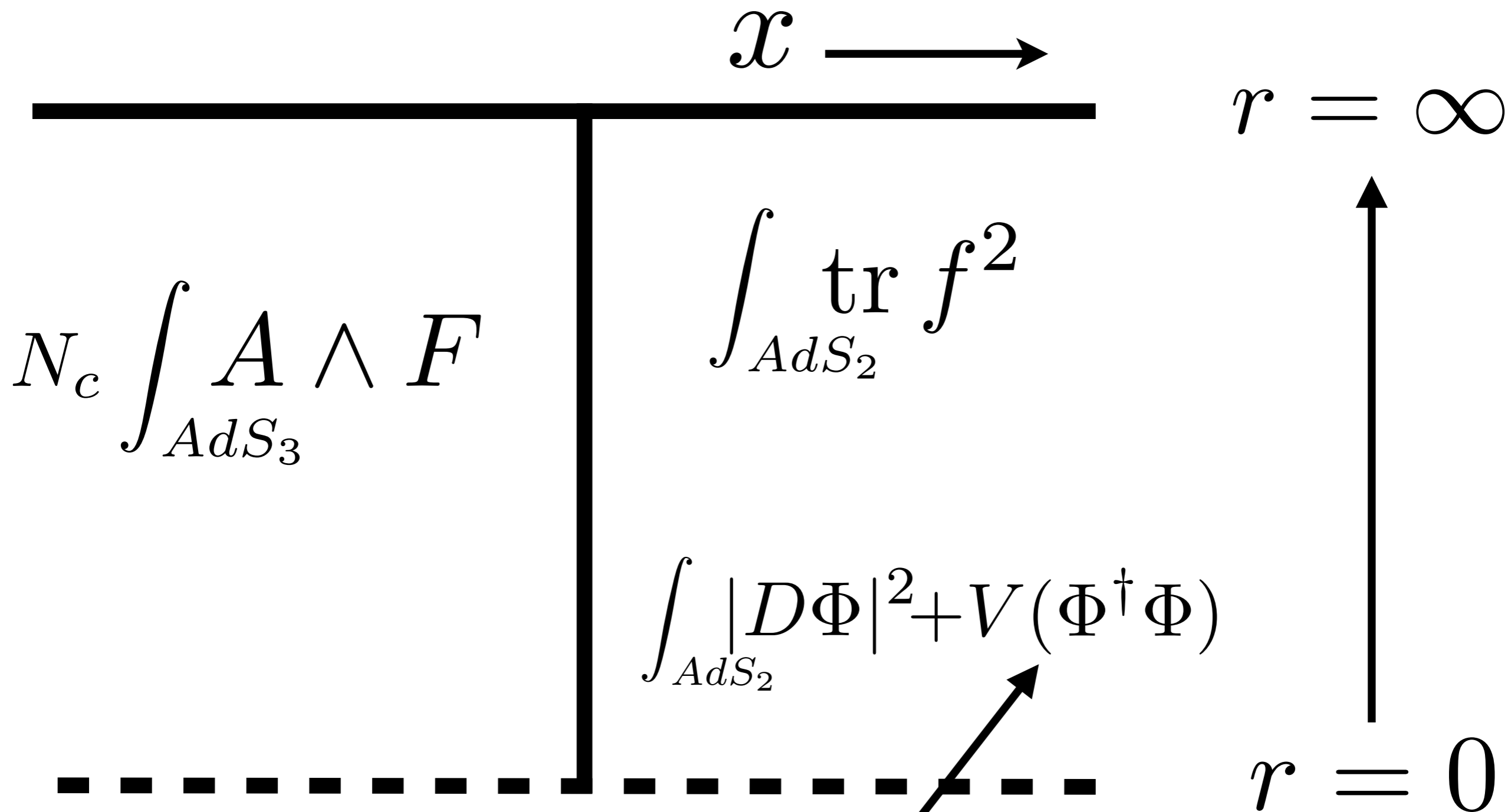
Bi-fundamental scalar in  $AdS_2$

# Top-Down Model



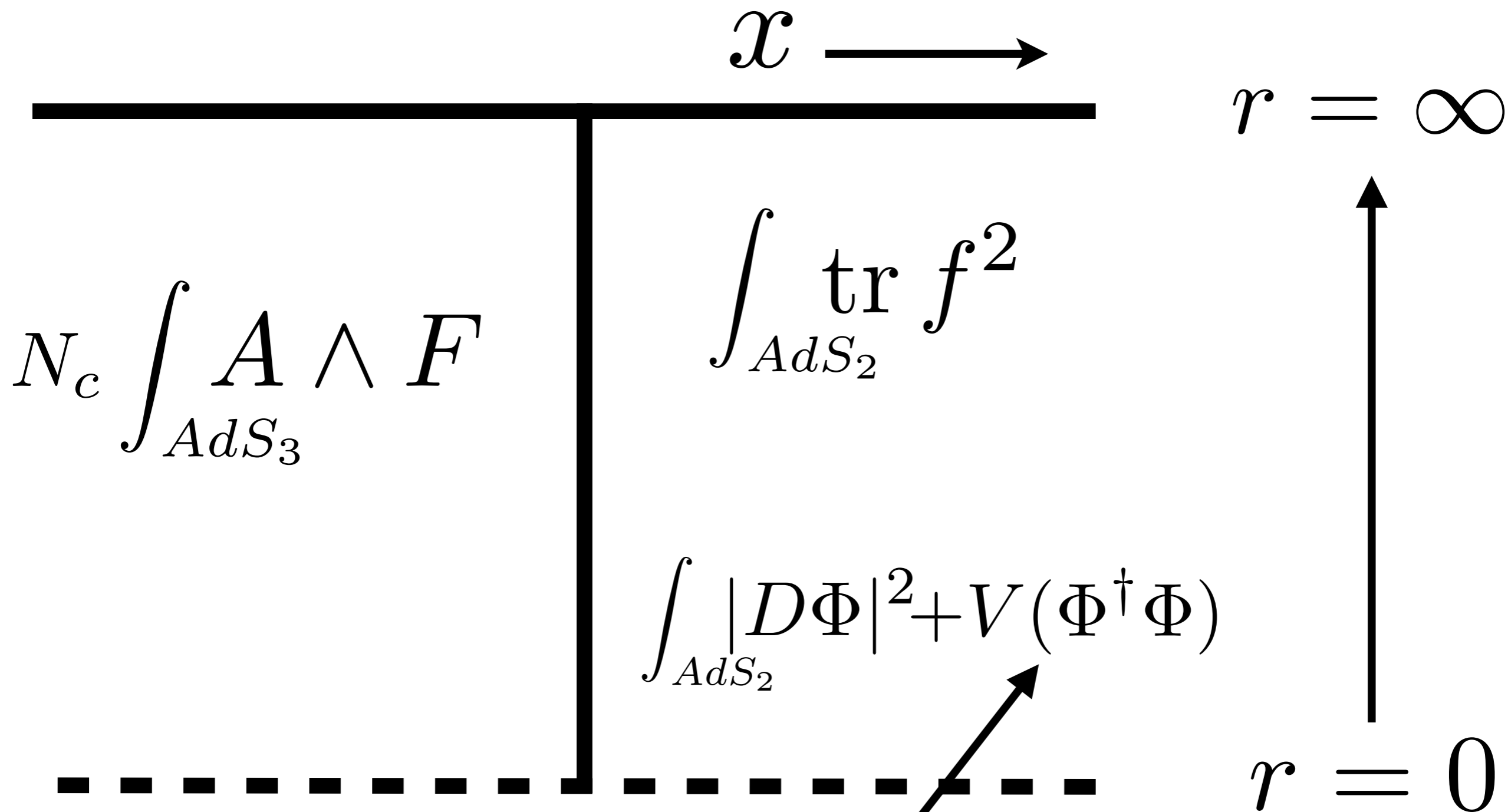
$$D\Phi = \partial\Phi + iA\Phi - ia\Phi$$

# Top-Down Model



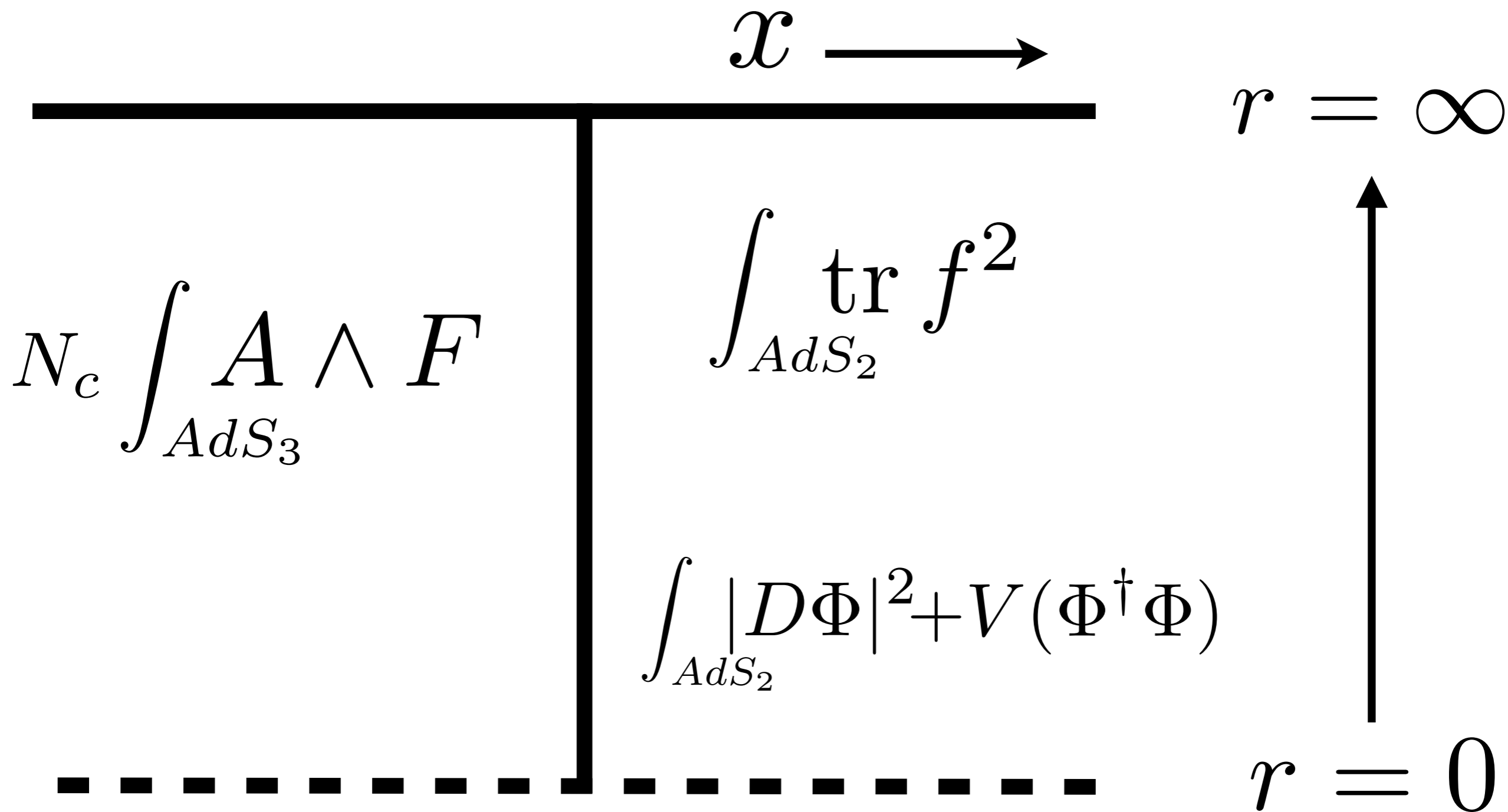
What is  $V(\Phi^\dagger \Phi)$ ?

# Top-Down Model



We don't know.

# Top-Down Model



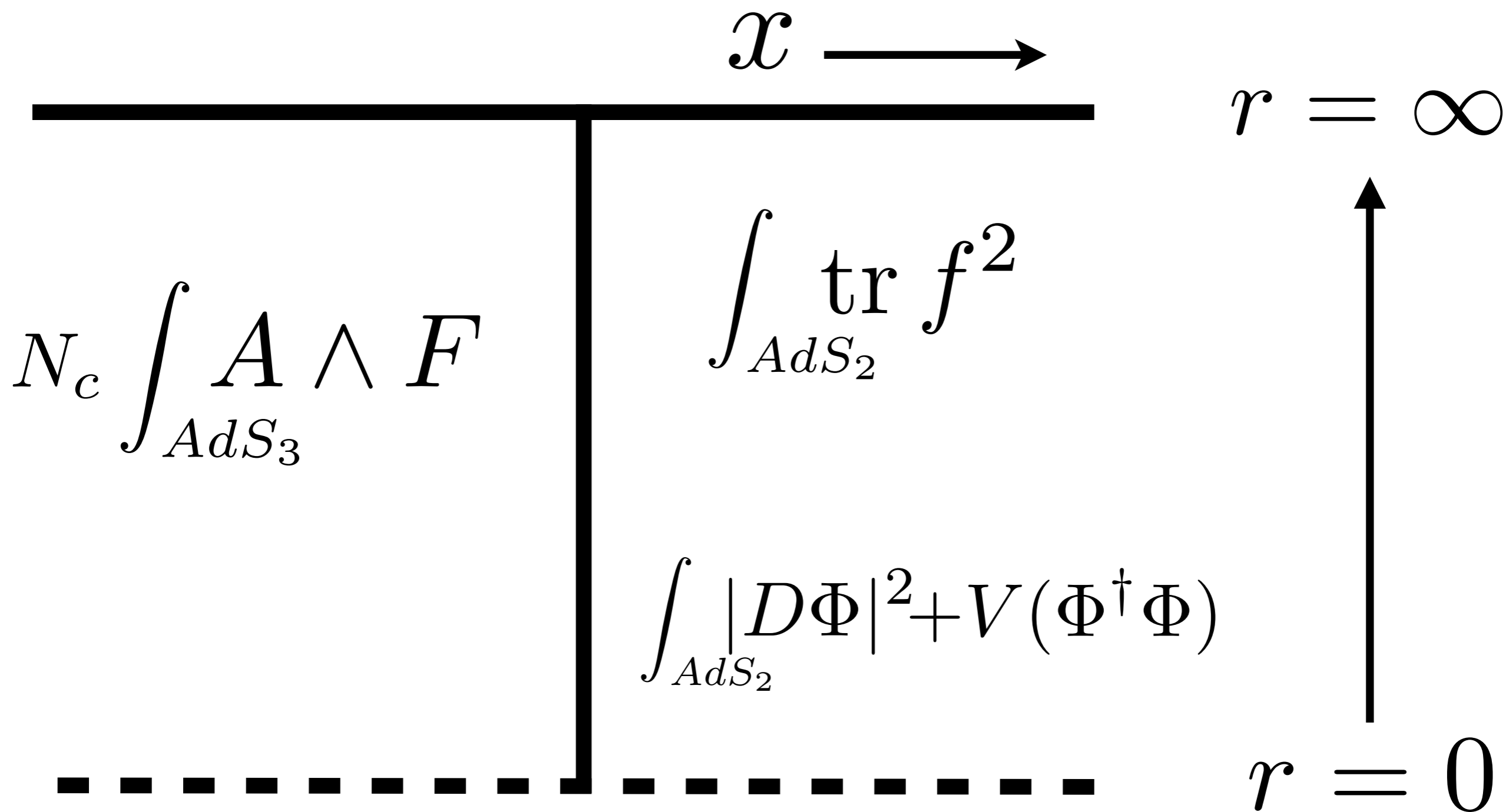
Switch to bottom-up model!



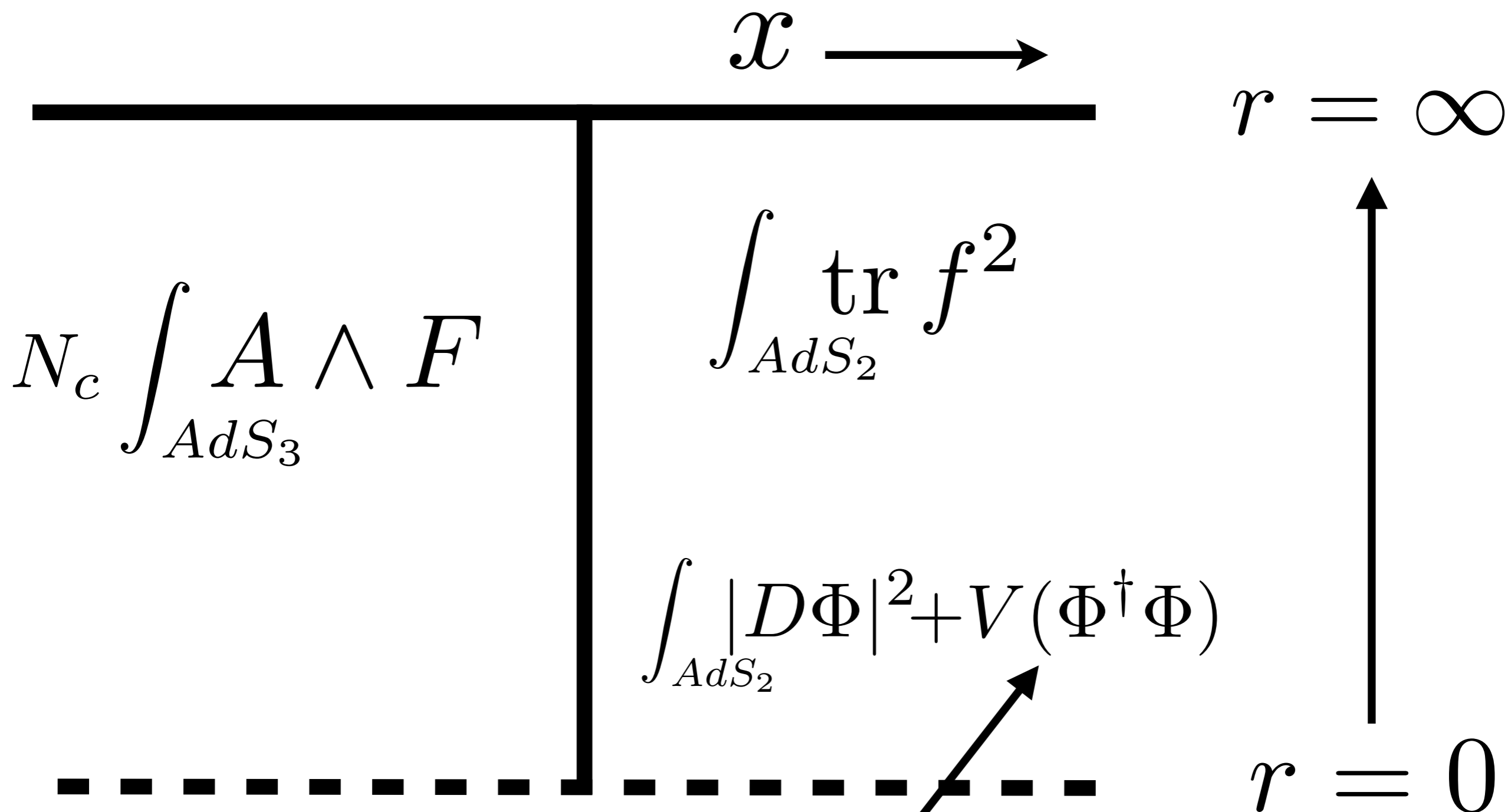
# Outline:

- The Kondo Effect
- The CFT Approach
- Top-Down Holographic Model
- Bottom-Up Holographic Model
- Summary and Outlook

# Bottom-Up Model

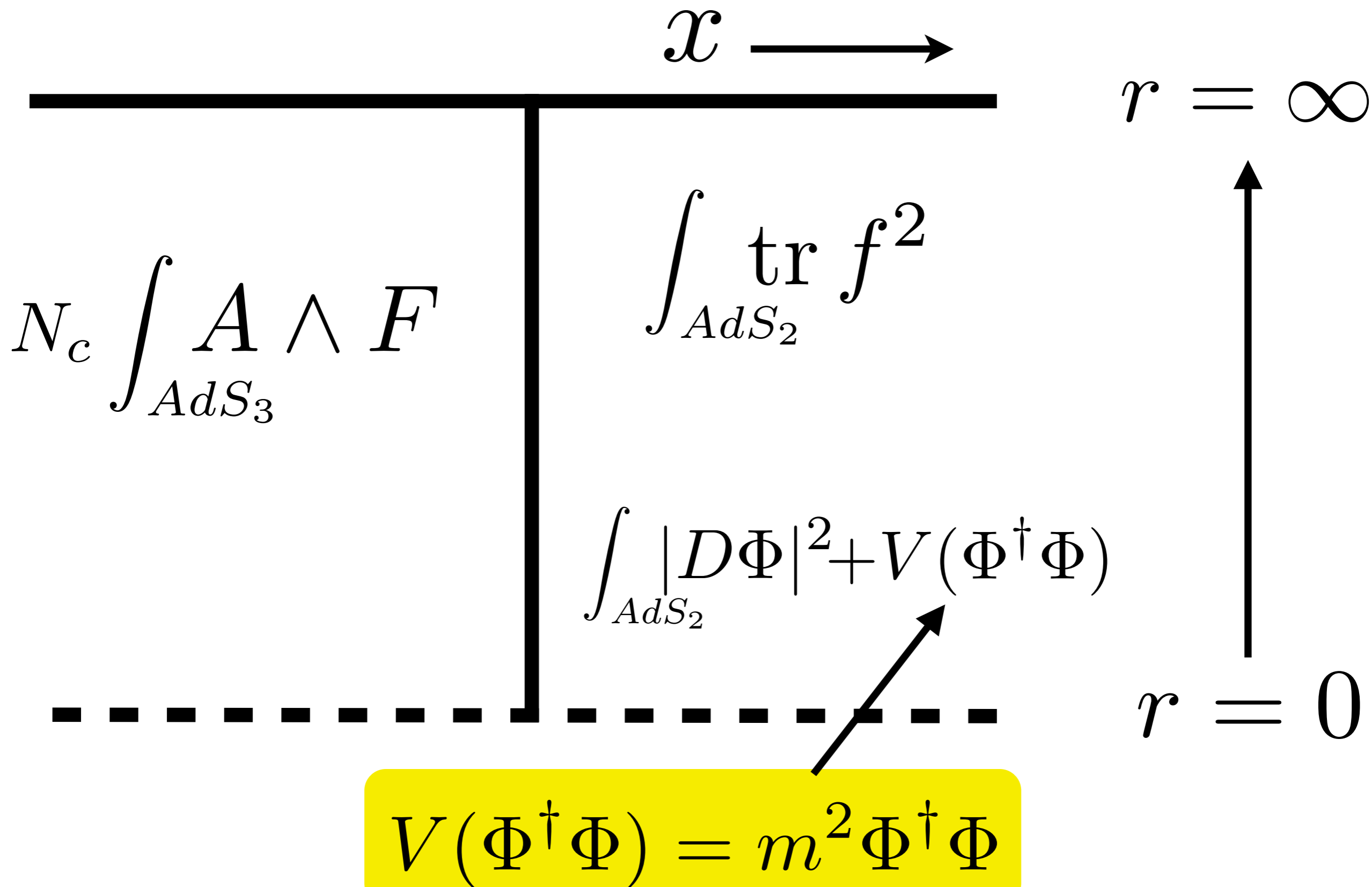


# Bottom-Up Model



We pick  $V(\Phi^\dagger \Phi)$

# Bottom-Up Model



# Boundary Conditions

We choose  $m^2 =$  Breitenlohner-Freedman bound

$$\Phi(r) = \tilde{c} r^{-1/2} + c r^{-1/2} \log r + \dots$$

Our double-trace (Kondo) coupling:

$$c = \tilde{g}_K \tilde{c}$$

Witten hep-th/0112258

Berkooz, Sever, Shomer hep-th/0112264

$$\sqrt{-g} f^{rt} \Big|_{\partial AdS_2} = Q$$

## Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

A holographic superconductor in  $AdS_2$

# Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

Superconductivity???

# Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

The large-N Kondo effect!



# Solutions of the Kondo Problem

Numerical RG (Wilson 1975)

Fermi liquid description (Nozières 1975)

Bethe Ansatz/Integrability  
(Andrei, Wiegmann, Tsvetick, Destri, ... 1980s)

Large-N expansion  
(Anderson, Read, Newns, Doniach, Coleman, ... 1970-80s)

Quantum Monte Carlo  
(Hirsch, Fye, Gubernatis, Scalapino, ... 1980s)

Conformal Field Theory (CFT)  
(Affleck and Ludwig 1990s)

# Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

$$T_c \propto T_K$$

# Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} |_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} |_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

Represents the formation of the Kondo singlet

## Phase Transition

$$T > T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) = 0$$

$$\langle \psi_L^\dagger \chi \rangle = 0$$

$$T < T_c$$

$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} \neq 0 \quad \phi(r) \neq 0$$

$$\langle \psi_L^\dagger \chi \rangle \neq 0$$

The phase transition is an ARTIFACT of the large-N limit!

The actual Kondo effect is a crossover

# Outline:

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# Summary

What is the holographic dual of the Kondo effect?

Holographic superconductor in  $AdS_2$

coupled as a defect

to a Chern-Simons gauge field in  $AdS_3$

# Outlook

- Entropy? Heat Capacity? Resistivity?
- Multi-channel?
- Other impurity representations?
- Entanglement entropy?
- Quantum Quenches?
- Multiple Impurities?

**Thank You.**