

Five dimensional gauge theories and Higgs branch at infinite coupling

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Five dimensional gauge theory

- Existence of SCFTs in 4 dimensions and higher
- Brane systems
- D4 - D8
- Webs of five branes
- M theory of CY singularity
- F theory with 8 supercharges

Gauge theory in 5d

SQCD

- Today focus on gauge theories with a single gauge group
- Relatively few theories which have a UV fixed point in 5d
- Number of flavors should be small enough
- Conditions similar to the condition for asymptotic freedom in 4d, but not the same

CS coupling

- The CS level k can be turned on, but fewer number of flavors are needed to have a 5d UV SCFT

SU(n) gauge theory

- Further restrict to SU(n) with N flavors and CS level k
- The level k is $1/2 + \text{integer}$ for N odd
- Integer for N even

Coulomb branch parameters

- Inverse gauge coupling $1/g^2$
- Masses for flavors
- Both have dimension 1
- Contribute to the central charge formula for BPS states

$SU(n)$ with $2n+4$ flavors and 0 CS level

- This theory has a 6d UV fixed point
- Trade 1 flavor by $1/2$ CS level increase by integrating out a massive flavor
- Get a condition for the existence of a 5d UV fixed point
- $N + 2k \leq 2n + 4$
- Excluding the case $N=2n+4, k=0$ which has 6d UV SCFT

Classical flavor symmetry

- $U(N) \times U(1)$
- Two $U(1)$ global symmetries
- Baryon number and the topological $U(1)$ symmetry which counts instanton number
- For $n=2$ this changes to $SO(2N) \times U(1)$

Gauge instantons in 5d

- This is a particle with a contribution to the mass that scales like $|I|/g^2$, with instanton charge I
- At infinite coupling, at the origin of the coulomb branch, becomes massless
- Increases the global symmetry
- Many new flat directions along the Higgs branch

Higgs branch at infinite coupling

- Given a theory with n , N , k set all parameters to 0 and ask what is the Higgs branch at infinite coupling

Higgs branch hyper Kähler cone

- Recall all parameters are set to zero - have a cone
- $SU(2)$ R symmetry
- Chiral ring
- Operators characterized by representation under $SU(2)$ R and representations under the global symmetry

HyperKähler cones

Global symmetry

- Theorem
- All operators with spin 1 under $SU(2)_R$ transform in the adjoint representation of the global symmetry

Instanton operators

- Assume there exists chiral operators in the chiral ring which transform under the $U(1)$ instanton, $SU(2)_R$, and global symmetries
- Non perturbative operators which do not show in the classical theory and become crucially relevant to the UV SCFT
- Analogs of 't Hooft monopole operators in 3d

Instantons and global symmetry

- If the instanton carries spin 1 under $SU(2)_R$ the global symmetry is enhanced
- All such instantons are highly restricted as they must complete the existing classical global symmetry C to form the adjoint representation of a bigger global symmetry F which contains the classical symmetry as a subgroup
- In all cases we know, C is a Levi subgroup of F - the ranks are equal

Global symmetry at infinite coupling

- $N = 2n + 3, k = 1/2: SO(4n + 8)$
- $N = 2n + 2, k = 1: SO(4n + 4) \times SU(2)$
- $N = 2n + 2, k = 0: SU(2n + 4)$
- $N = 2n + 1, k = 3/2: SO(4n + 2) \times U(1)$
- $N = 2n + 1, k = 1/2: SO(2n + 2) \times SU(2)$
- $N = 2n, k = 2: SO(4n) \times U(1)$

More global symmetry

- $N = 2n, k = 1: SU(2n + 1) \times U(1)$
- $N = 2n, k = 0: SU(2n) \times SU(2) \times SU(2)$
- $N = 2n - 1, k = 5/2: SO(4n - 2) \times U(1)$
- $N = 2n - 1, k = 3/2: SU(2n) \times U(1)$
- $N = 2n - 1, k = 1/2: SU(2n - 1) \times SU(2) \times U(1)$

Instantons and Higgs branch

- If the instanton has a spin under $SU(2)_R$ which is higher than 1, the Higgs branch becomes larger, by a significant amount
- Roughly double the dimension

Instantons transform under $SU(2)$ R

- To compute this quantity, we need to evaluate the energy of a vacuum state in the topological sector of 1 instanton
- For a gauge group G this is $h/2$ where h is the dual Coxeter number of G
- Closest analog to the 4d beta function for a SYM theory

Instanton zero modes

- Each flavor of quarks contributes a fermionic zero mode
- Easy to see in the D4-D8 system by looking at D0-D8
- Or D1-D7 for a five brane web
- As a result, the instanton transforms under the global symmetry as
- spinor representation if it is a rotation group
- Antisymmetric representation if it is a unitary group

Higgs branch at infinite coupling

- Given this data, how to proceed?
- Two ways of constructing hyperKahler cones
- HyperKahler quotient (Higgs branch) F & D terms
- 3D N=4 Coulomb branch

Construction using 3d techniques

- Assume
- The answer is given by a Coulomb branch of 3d $N=4$ theory
- Use the global symmetry as an input
- The instantons transform in spin $n/2$ of $SU(2)_R$ and spinor or maximally antisymmetric representation of the global symmetry

3d $N=4$ Coulomb branch

Global symmetry

- An imbalance of a node is the number of flavors minus twice the number of colors
- A node is balanced if the imbalance is 0
- The subset of balanced nodes forms the Dynkin diagram of the global symmetry
- If the imbalance of a node is $n-2$ then there is an operator in the chiral ring with spin $n/2$ under $SU(2)_R$ and representation under the global symmetry given by the node is attached to

Generators of the chiral ring

- The global symmetry of this Coulomb branch $SO(4n + 8)$
- An adjoint valued operator at spin 1 of $SU(2)_R$
- A spinor valued operator at spin $n/2$ of $SU(2)_R$
- Mesons, baryons, gaugino bilinear, instantons
- satisfying relations that are dictated by the quiver and the symmetries

E7 sequences

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(2) \end{array} - \begin{array}{c} \blacksquare \\ SO(12) \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{cccccccc} & & & \bullet 2 & & & & \\ & & & | & & & & \\ \circ & - & \circ & - & \circ & - & \circ & - & \circ & - & \circ & - & \circ \\ 1 & & 2 & & 3 & & 4 & & 3 & & 2 & & 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_1 \end{array} - \begin{array}{c} \blacksquare \\ 2n+2 \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{cccccccc} & & & \circ n & & & & \\ & & & | & & & & \\ \circ & - & \dots & - & \circ & - & \circ & - & \circ & - & \bullet & - & \circ \\ 1 & & & & 2n-1 & & 2n & & n+1 & & 2 & & 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_0 \end{array} - \begin{array}{c} \blacksquare \\ 2n+2 \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{cccccccc} & & & \bullet 2 & & & & \\ & & & | & & & & \\ \circ & - & \dots & - & \circ & - & \circ & - & \circ & - & \dots & - & \circ \\ 1 & & & & n+1 & & n+2 & & n+1 & & & & 1 \end{array} \right)$$

E5 sequences

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(2) \end{array} - \begin{array}{c} \blacksquare \\ SO(8) \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{c} 1 \quad 2 \quad 1 \\ \circ - \circ - \circ \\ \circ - \circ - \circ \\ 1 \quad 2 \quad 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_2 \end{array} - \begin{array}{c} \blacksquare \\ 2n \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{c} \quad \quad \quad \\ - - \dots - - - - \\ 1 \quad 2 \quad \quad 2n-3 \quad 2n-2 \quad n \quad 1 \\ \quad \quad \quad \quad \quad \bullet 1 \\ \quad \quad \quad \quad \quad - - \bullet 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_1 \end{array} - \begin{array}{c} \blacksquare \\ 2n \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{c} \quad \quad \quad \quad \\ - - \dots - - - - - \\ 1 \quad 2 \quad \quad n-1 \quad n \quad n \quad n-1 \quad \quad 1 \\ \quad \quad \quad \quad \bullet 1 \quad \bullet 1 \quad \quad \quad \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_0 \end{array} - \begin{array}{c} \blacksquare \\ 2n \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{c} \quad \quad \\ - \bullet - \\ - - \dots - - - \\ 1 \quad \quad n \quad \quad 1 \end{array} \right)$$

E4 sequences

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_{\frac{5}{2}} \\ \text{---} \quad \blacksquare \\ 2n-1 \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{ccccccc} & & & & & \bullet 1 & \\ & & & & & | & \\ & & & & & \circ n-1 & \\ & & & & & | & \\ \circ 1 & \text{---} & \circ 2 & \text{---} & \dots & \text{---} & \circ 2n-4 & \text{---} & \circ 2n-3 & \text{---} & \circ n-1 & \text{---} & \bullet 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_{\frac{3}{2}} \\ \text{---} \quad \blacksquare \\ 2n-1 \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{ccccccccccc} & & & & & & \bullet 1 & \text{---} & & & \bullet 1 & \\ & & & & & & | & & & & | & \\ \circ 1 & \text{---} & \circ 2 & \text{---} & \dots & \text{---} & \circ n-2 & \text{---} & \circ n-1 & \text{---} & \circ n-1 & \text{---} & \circ n-1 & \text{---} & \circ n-2 & \text{---} & \circ 2 & \text{---} & \circ 1 \end{array} \right)$$

$$\mathcal{H}_\infty \left(\begin{array}{c} \circ \\ SU(n)_{\frac{3}{2}} \\ \text{---} \quad \blacksquare \\ 2n-1 \end{array} \right) = \mathcal{C}_{3d} \left(\begin{array}{ccccccccccc} & & & & & & & & \circ 1 & & & & & & & & & & & & \\ & & & & & & & & / & & \backslash & & & & & & & & & & \\ & & & & & & & & \bullet 1 & & \bullet 1 & & & & & & & & & & \\ & & & & & & & & | & & | & & & & & & & & & & \\ \circ 1 & \text{---} & \circ 2 & \text{---} & \dots & \text{---} & \circ n-2 & \text{---} & \circ n-1 & \text{---} & \circ n-1 & \text{---} & \circ n-2 & \text{---} & \circ 2 & \text{---} & \circ 1 \end{array} \right)$$

Minimally unbalanced quivers

- Quivers with very few unbalanced nodes
- Global symmetry
- 1 unbalanced node gives either a single non Abelian factor or 2 non Abelian factors
- If there are p unbalanced nodes, expect $p-1$ $U(1)$ factors
- Gradually complicated moduli spaces

Summary

- The Higgs branch at infinite coupling of a 5d theory is conveniently encoded by a coulomb branch of a 3d $N=4$ theory
- The chiral ring can be derived from this description
- The generators of the chiral ring are simple and describe the solution to this problem
- Get a window to non perturbative effects made by instantons and the precise way they correct classical relations in the chiral ring

Thank you !