

# Stringy Horizons and UV/IR mixing

Based on:

**1502.03633** with Amit Giveon and David Kutasov

**1506.07323** with Roy Ben-Israel, Amit Giveon and Lior Liram

And work to appear

It is widely believed that stringy corrections ( $g_s = 0$  and finite  $\alpha' = l_s^2$ ) are boring for large BH.

Indeed perturbative stringy corrections are.

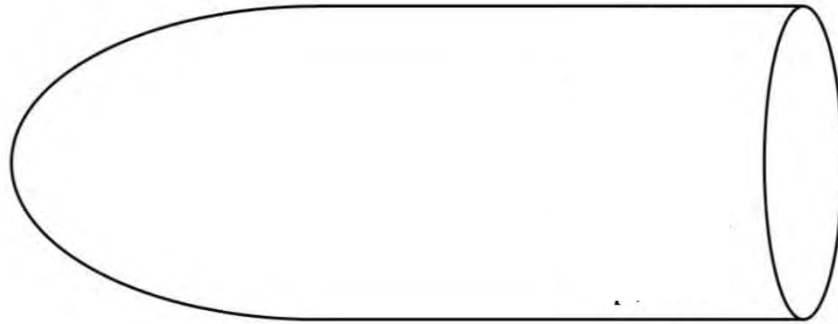
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Goal of the talk: show that non-perturbative stringy corrections are far from being trivial.

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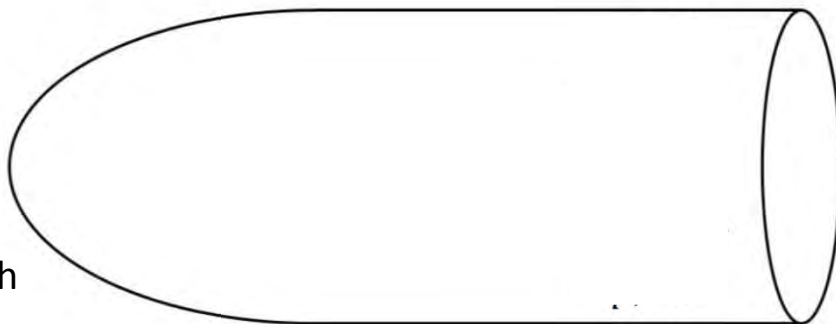


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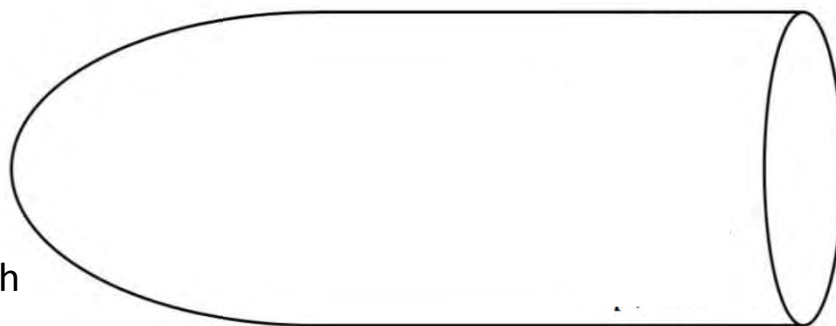


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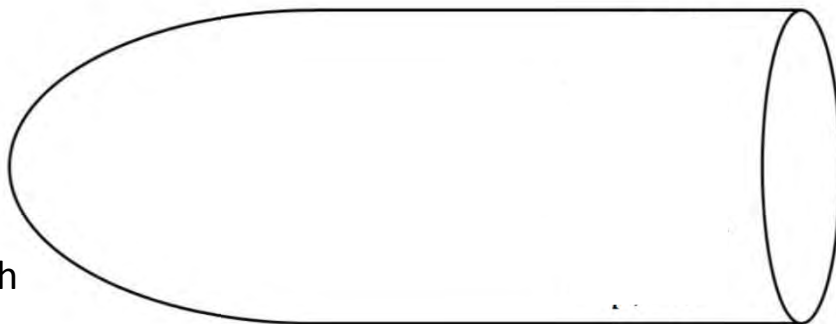
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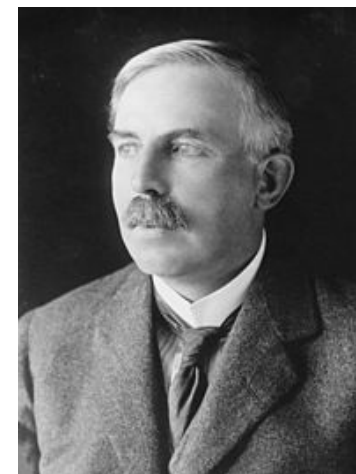
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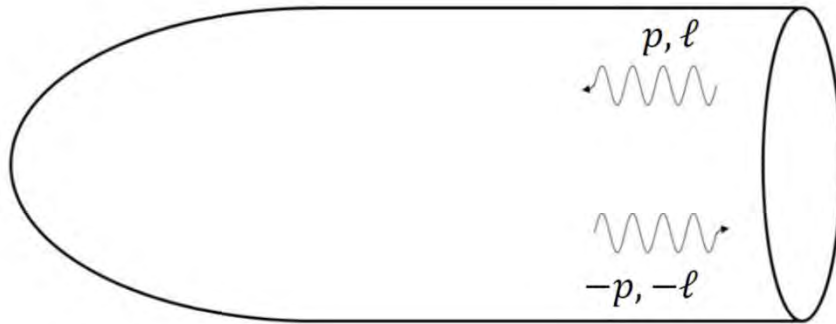
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We do so is by following this guy

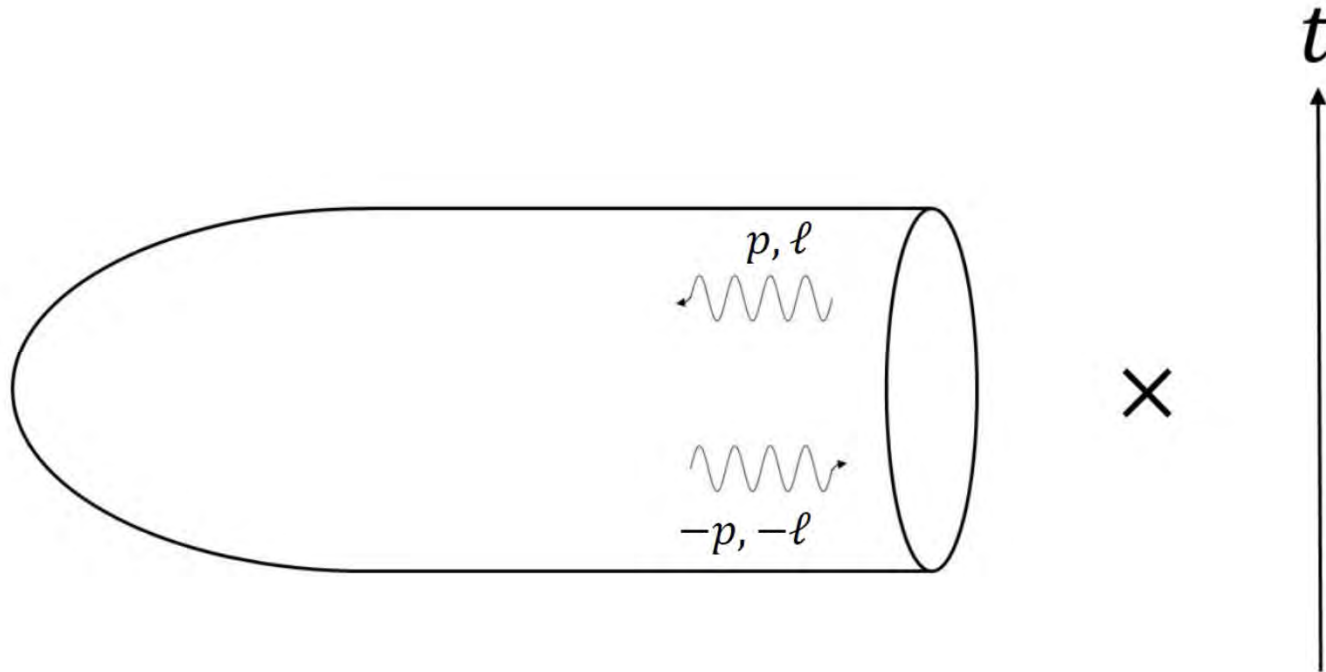


Following Rutherford we do this experiment

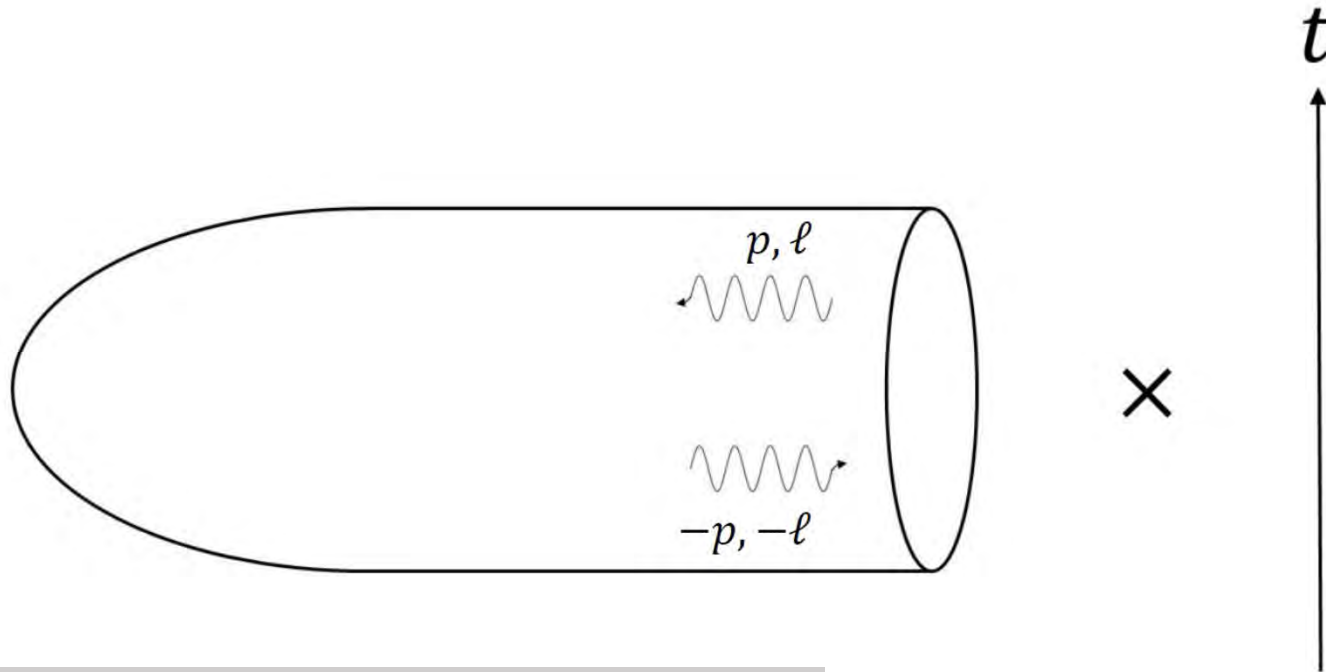




Following Rutherford we do this experiment. To set it up in string theory we add an external time direction

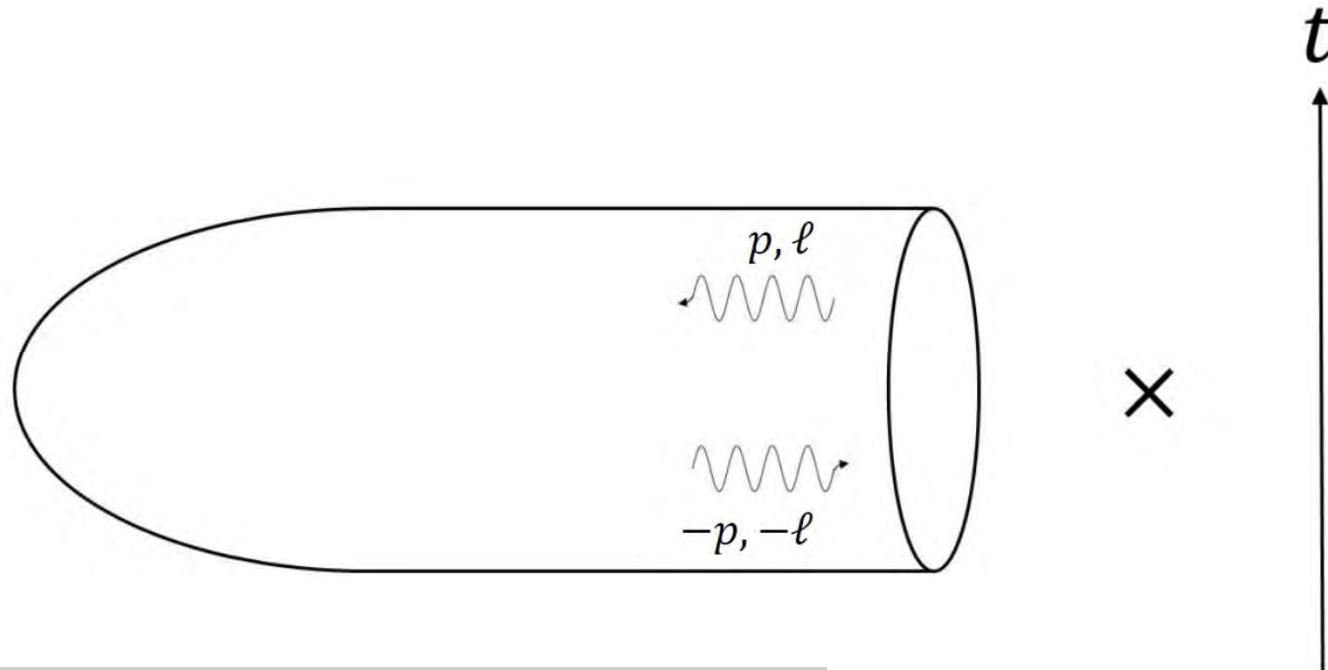


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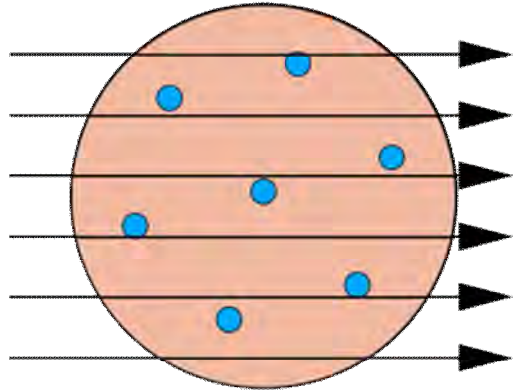
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Our job is data analysis:

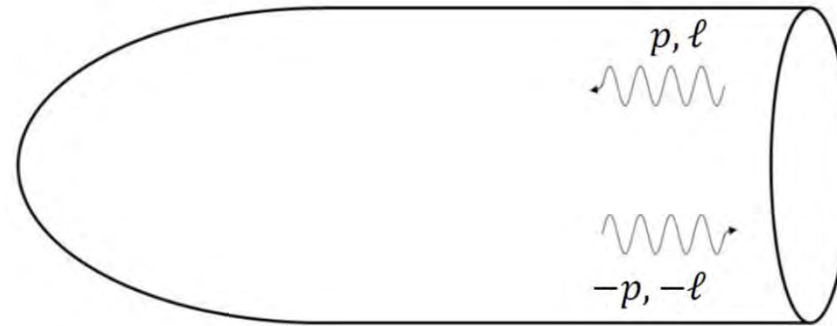
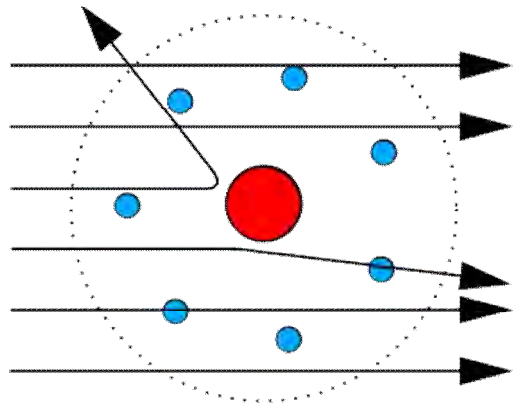
translate the CFT results into target-space info.

Just like in Rutherford's case, there are two options

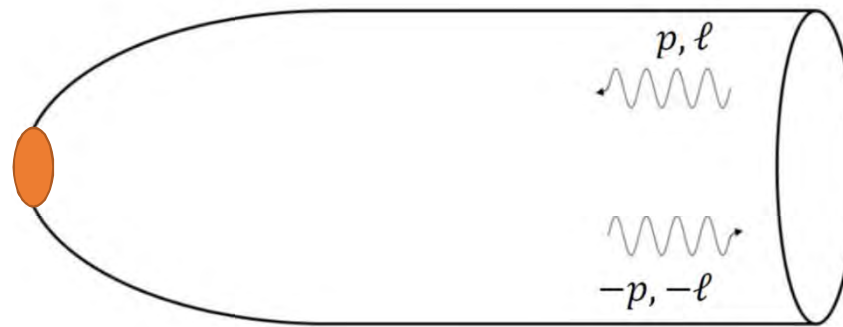
THOMSON



RUTHERFORD

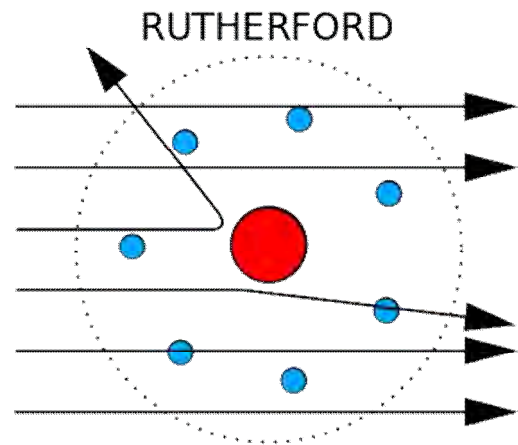
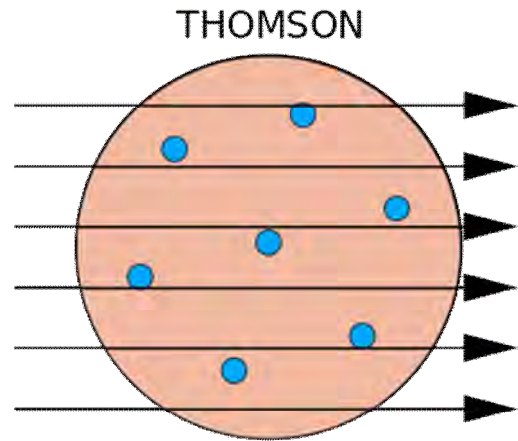


The tip is not special.

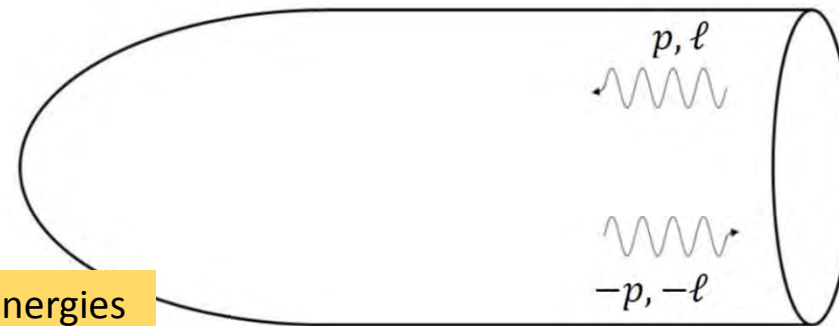


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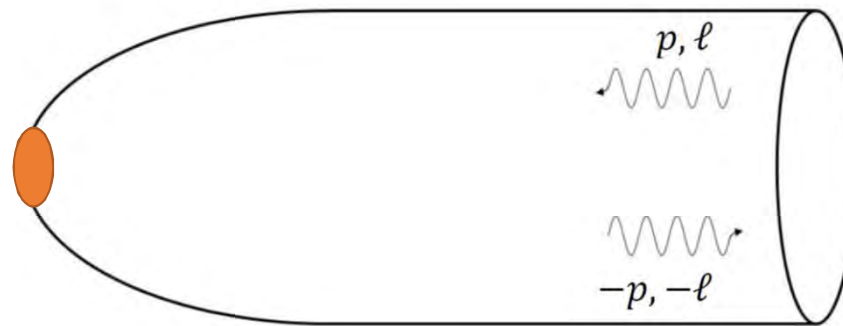
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Going to high energies  
was and is the key.




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The reflection coefficient takes the form

$$R(p, l) = R_{per}(p, l)R_{non-per}(p)$$



Obtained by solving the KG  
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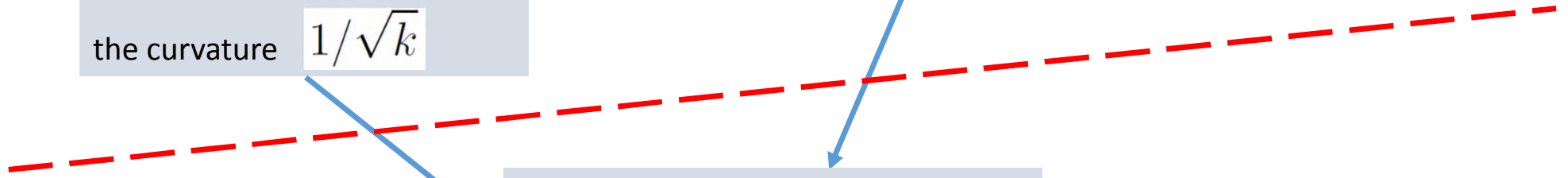
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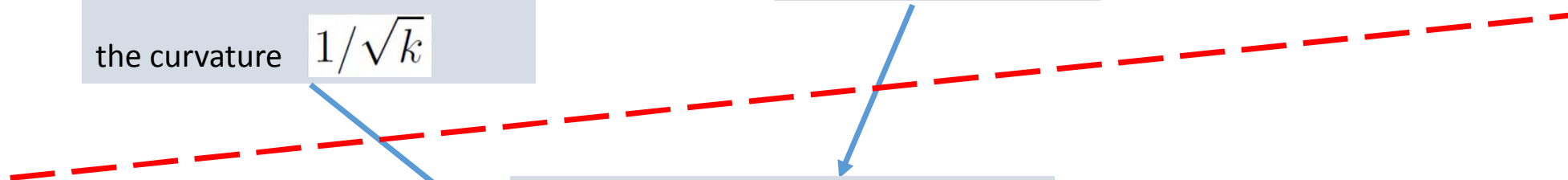
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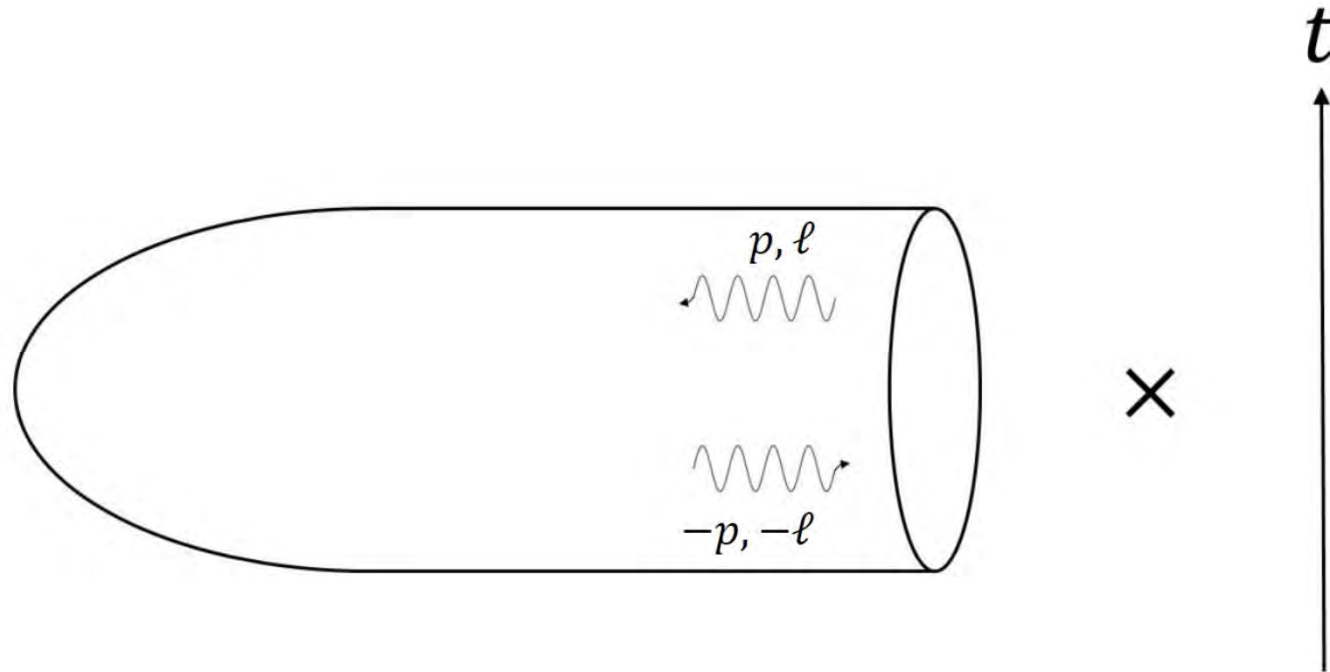
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A bit more details: For simplicity we take  $l=0$  and the on-shell condition  $E=p$

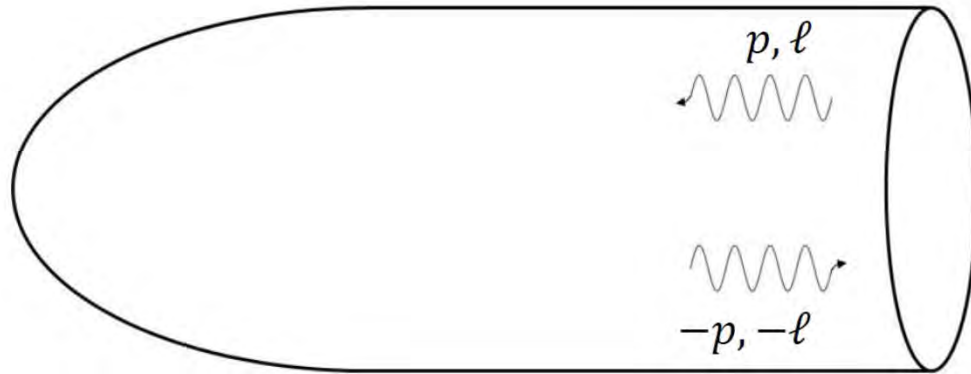


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Then we have the reflection coefficient as a function of  $E$ .

Fourier transform

gives  $f(t)$ .



$\times$

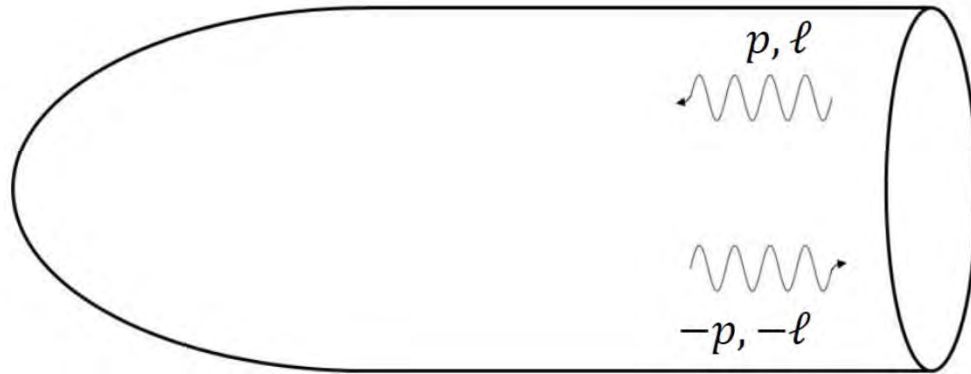
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A vertical arrow pointing upwards, indicating the direction of time  $t$ .

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For large  $t$ , up to tiny corrections,

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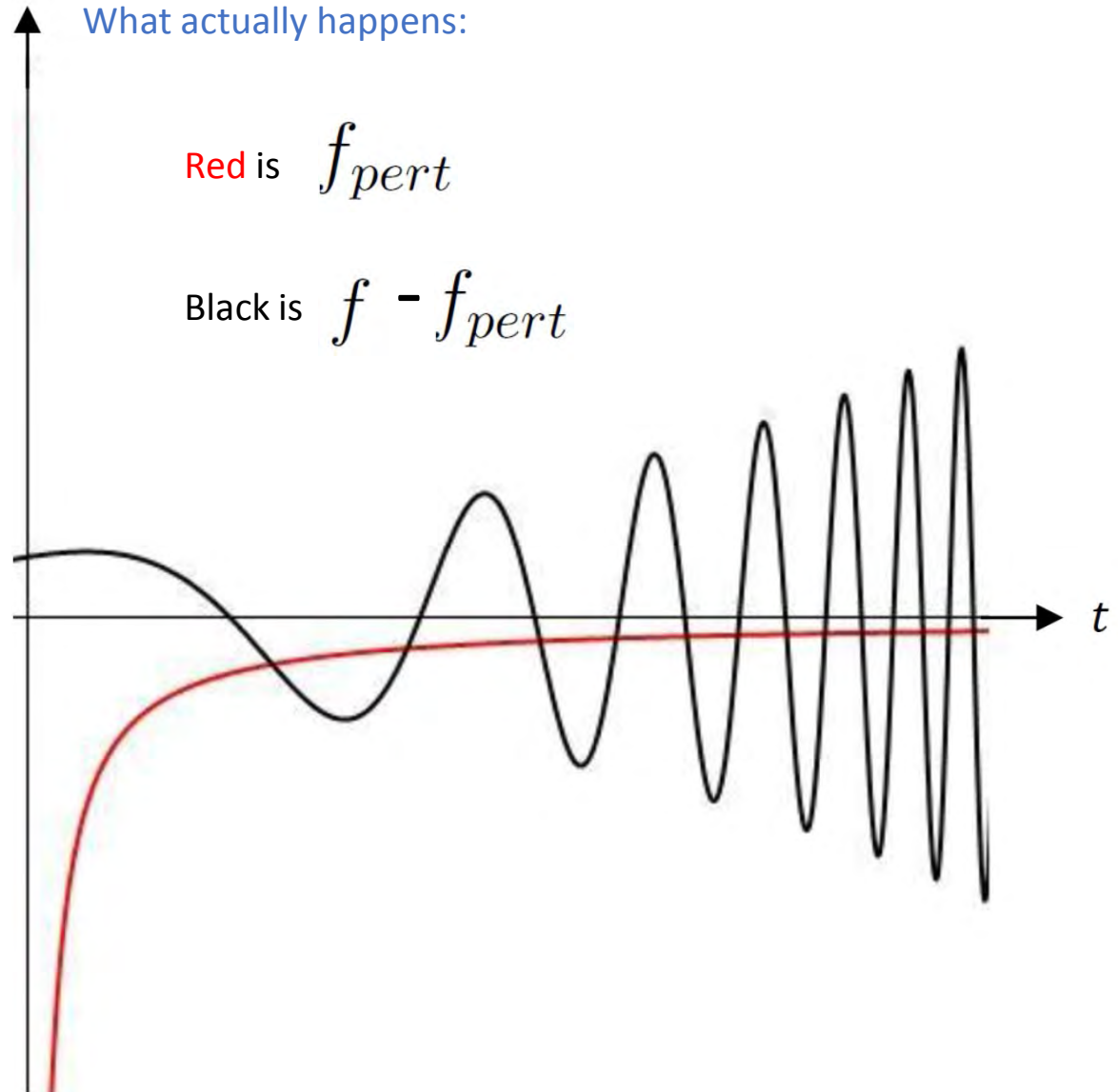
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What actually happens:

Red is  $f_{pert}$

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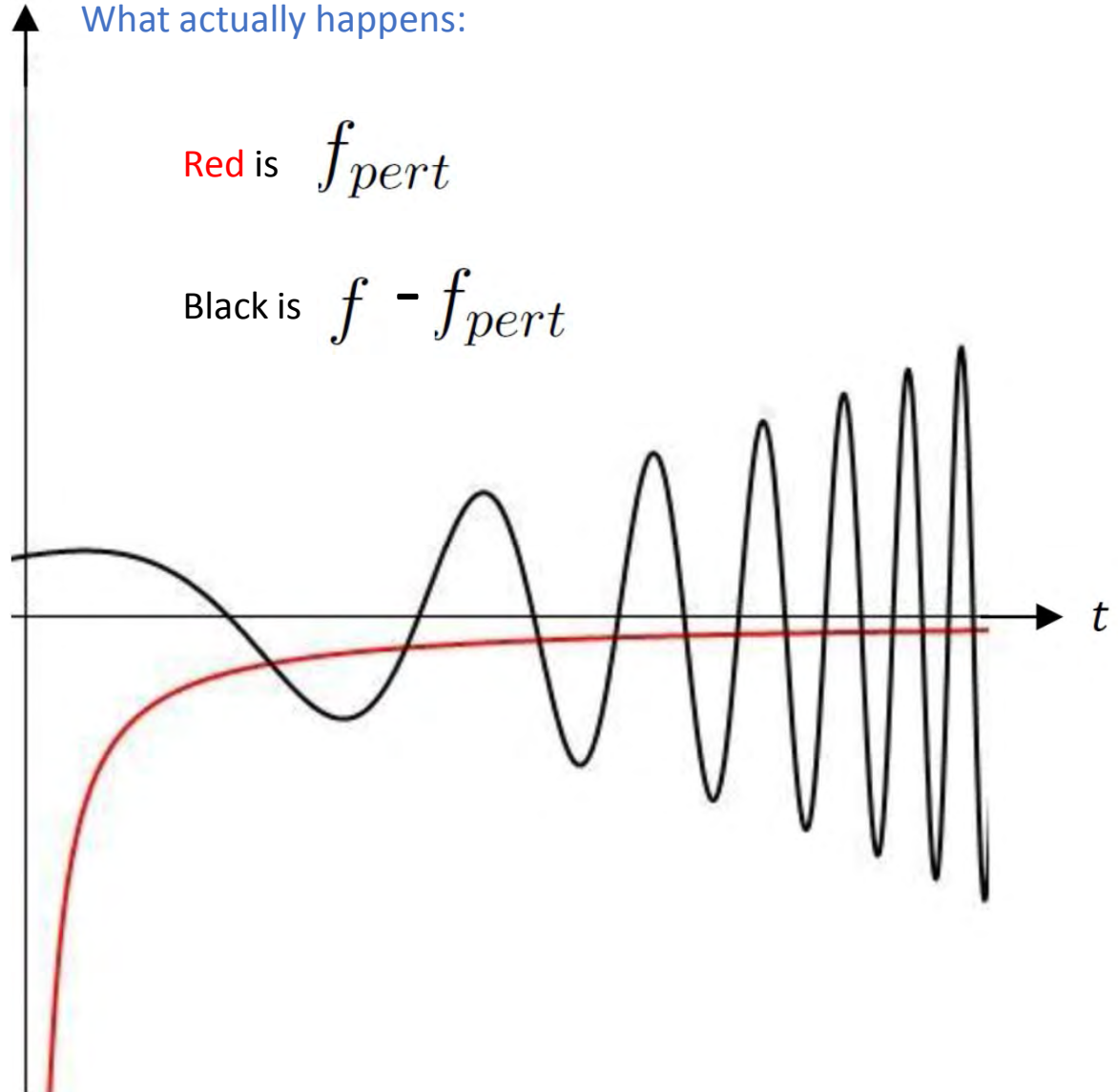
Rest of the talk:

1. How come?
2. Why this is relevant to the BH information puzzle?

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## How come?

At high energies (compared to the curvature scale  $1/\sqrt{k}$ ) we have  $R_{pert} = \text{const.} \iff \delta_{pert} = -\pi(1/2 + |\ell|)$

This is **not** a special property of the model.

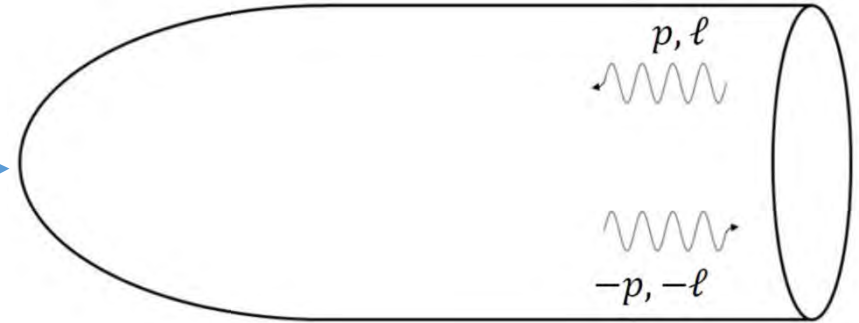


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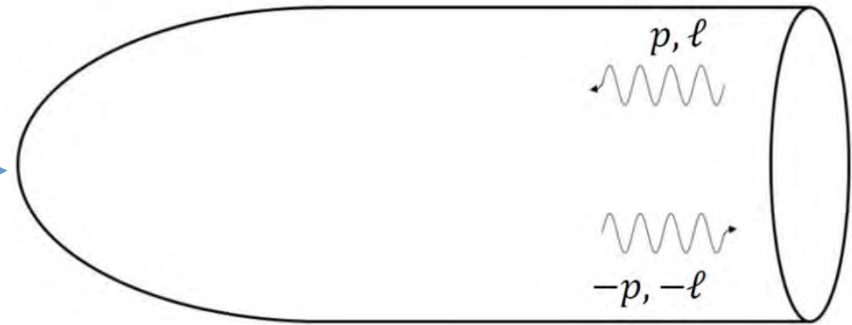


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Indeed it is at low energies.

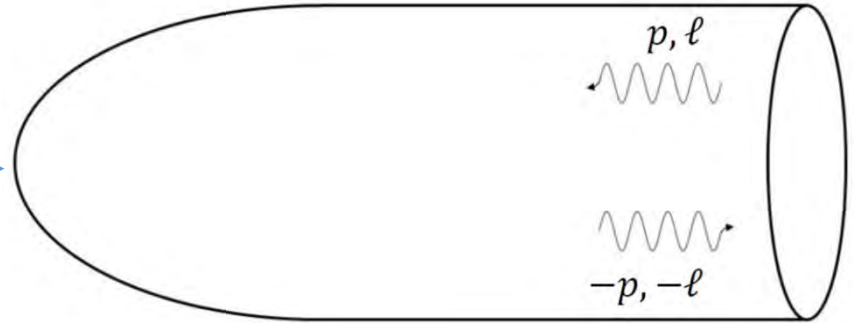
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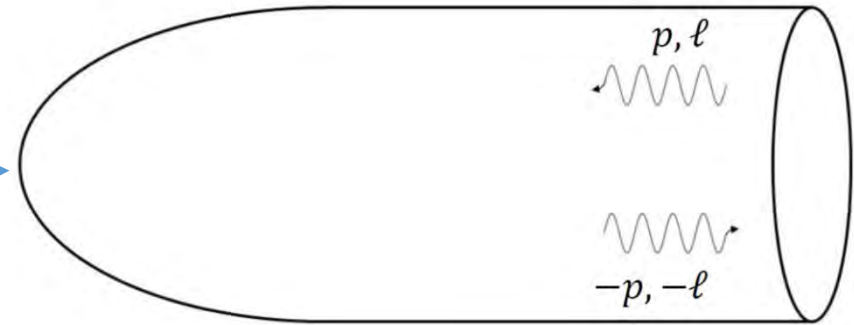
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The FZZ duality was believed to work in the following way:

Cigar is a good description at large  $k$  and

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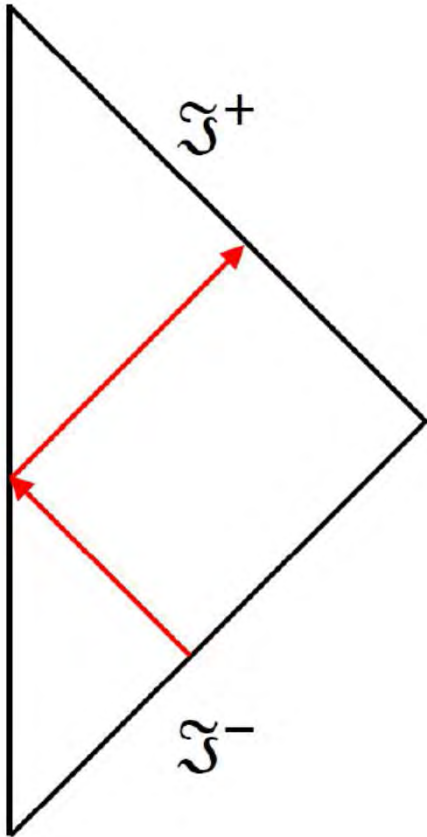
We see that in fact at large k we have:

Cigar is a good description at low energies.

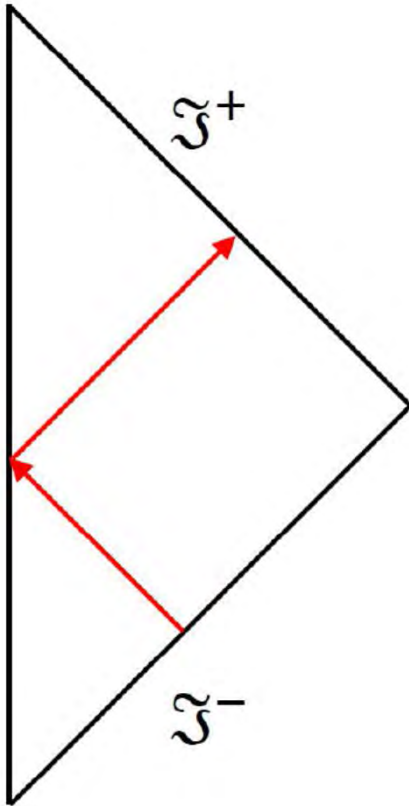
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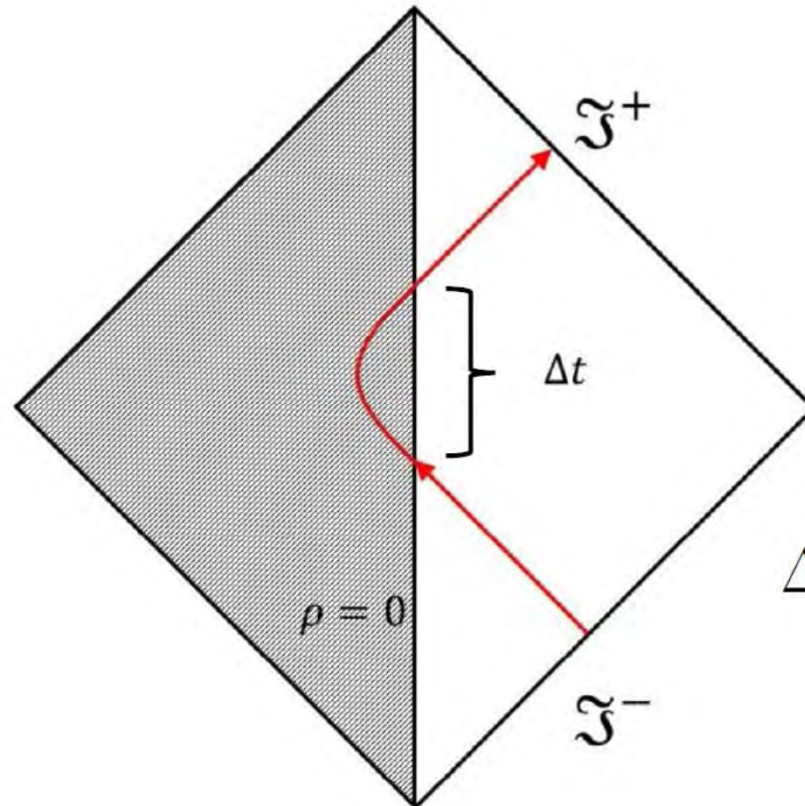
At low energies



At low energies



At high energies



$$\Delta t \sim \frac{1}{\sqrt{k}} \log(E)$$

UV/IR mixing

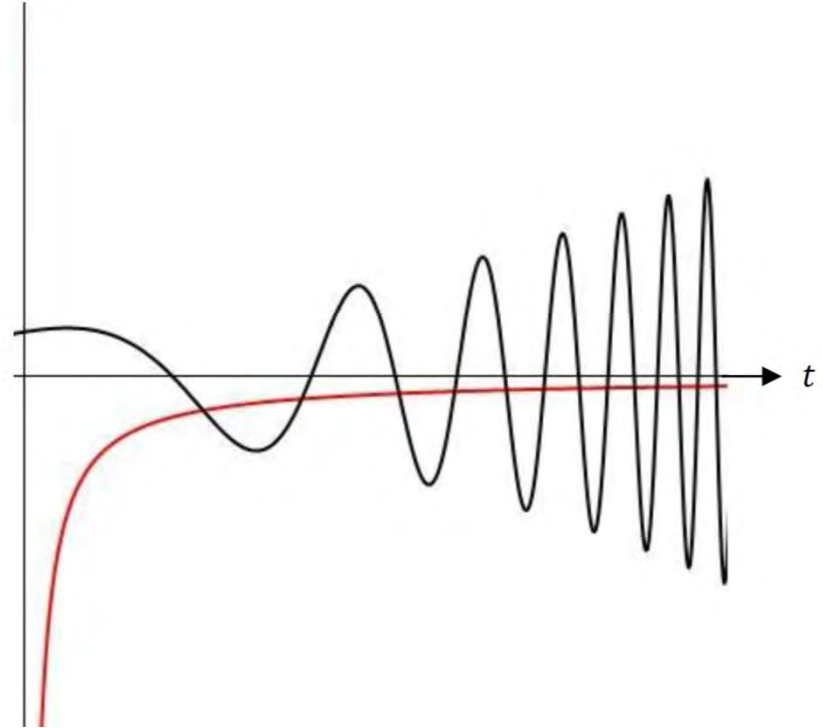
# Coarse graining

Suppose that we have an uncertainty in the time separation  $\delta t$   
then because the frequency blows up faster than the amplitude  
if we coarse grain

$$f_{\delta t}^{\text{coarse-grain}}(t) \equiv \frac{1}{\delta t} \int_t^{t+\delta t} f(t') dt'.$$

we get that for  $\delta t \gg e^{-\frac{1}{2}\sqrt{\frac{k}{2}}t}$

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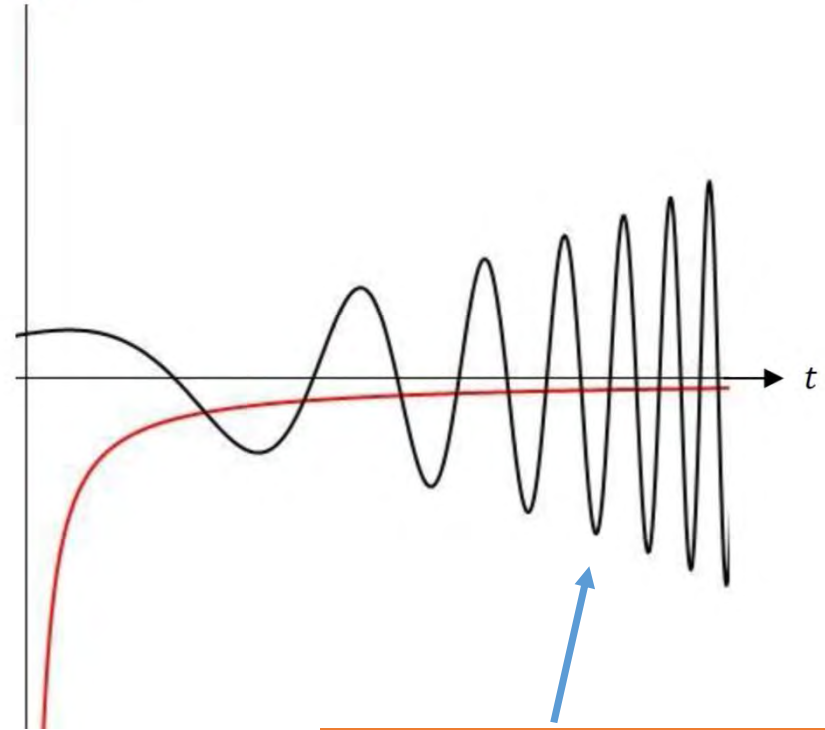
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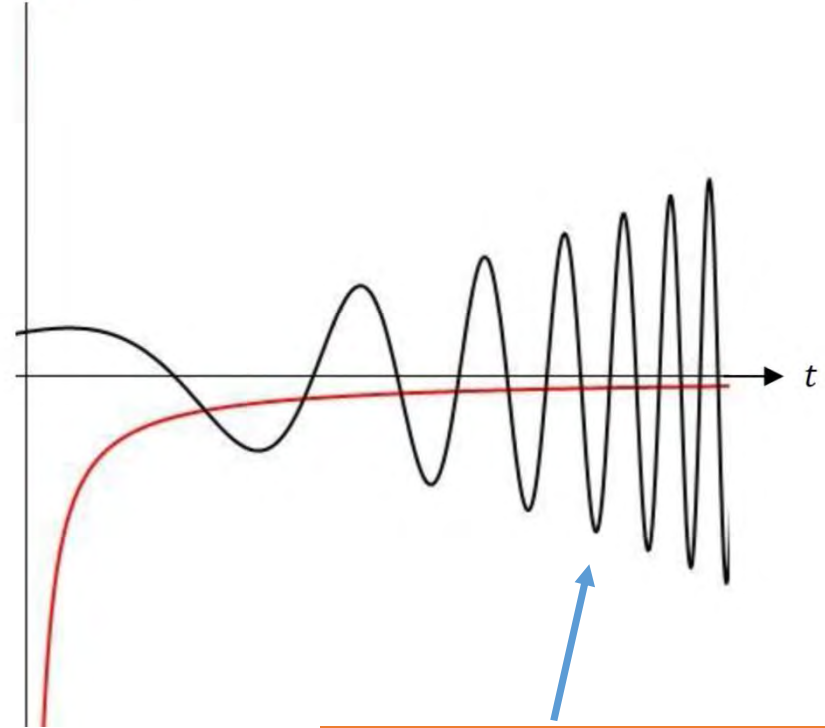
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This suggest an analogy with BH hair.

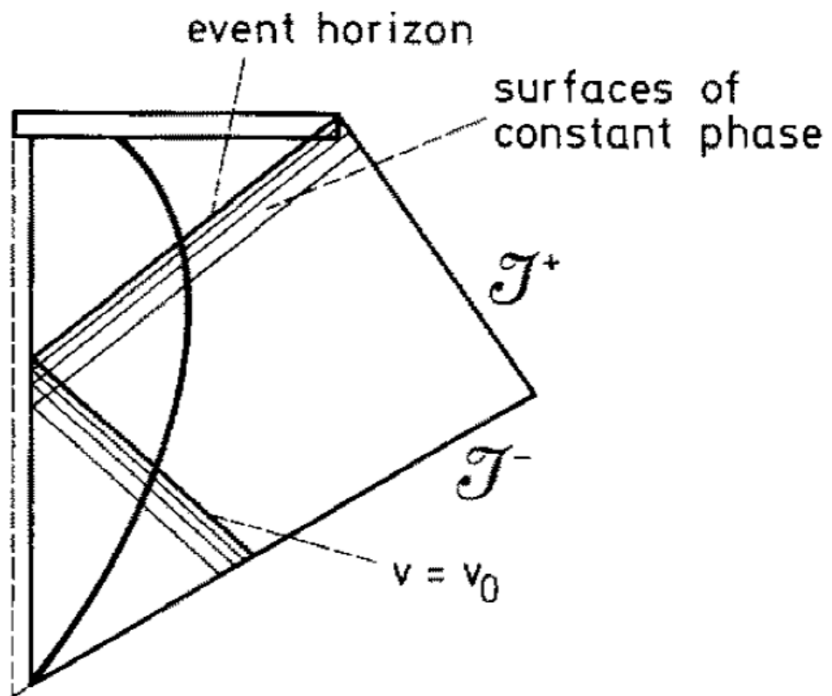
Can be made rather precise.



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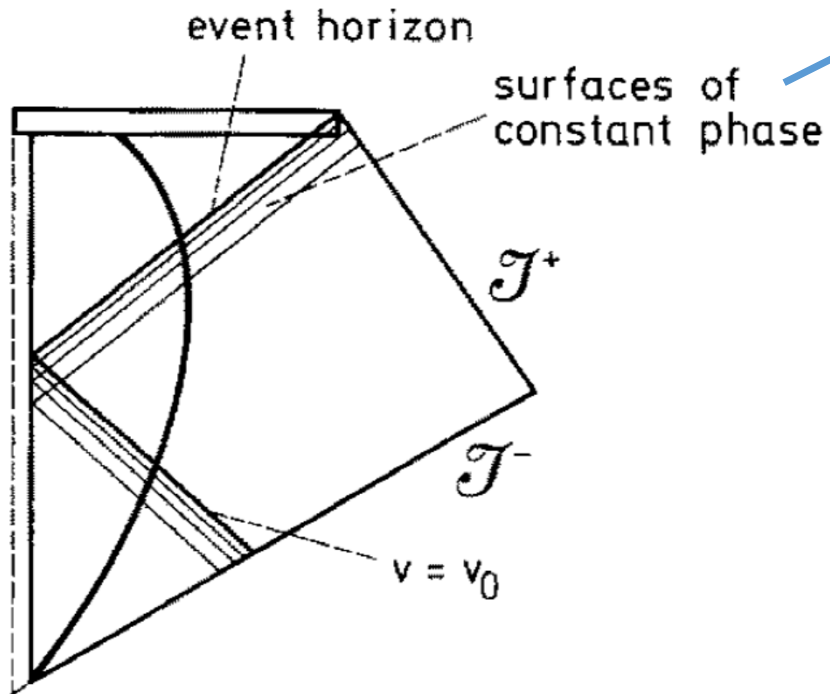
## Black Holes Information Puzzle

Hawking calculation in GR (figure from Hawking's paper)



Black Holes Information Puzzle

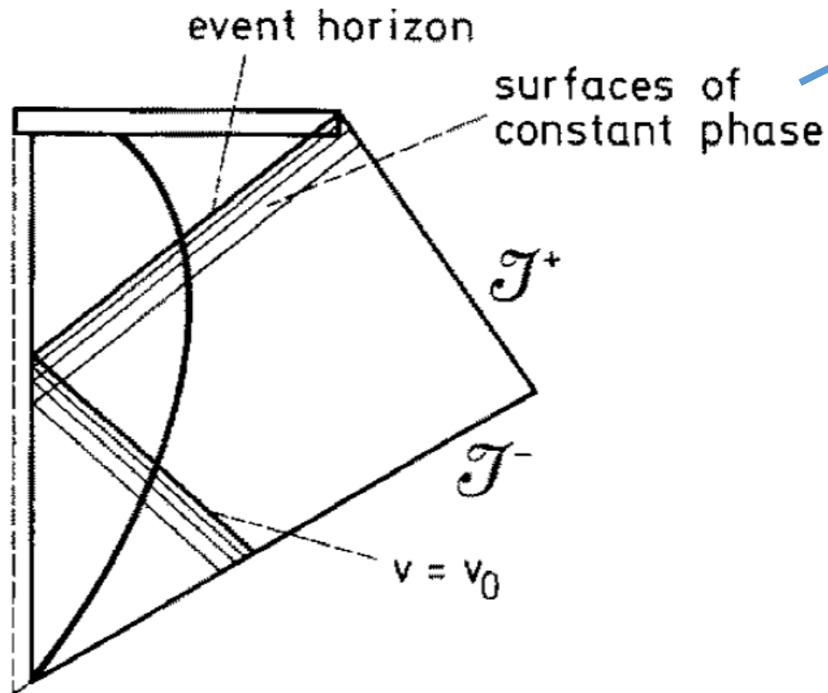
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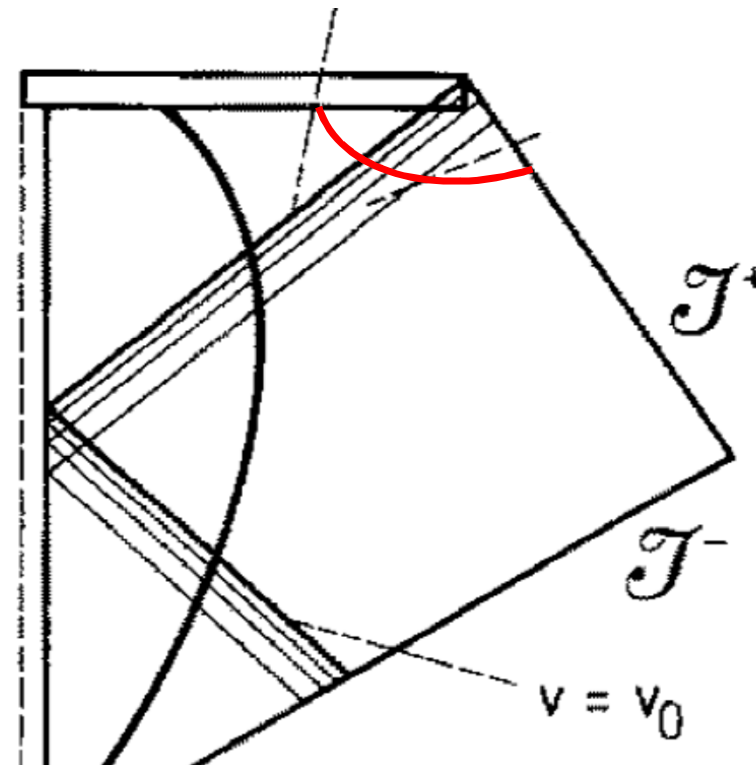
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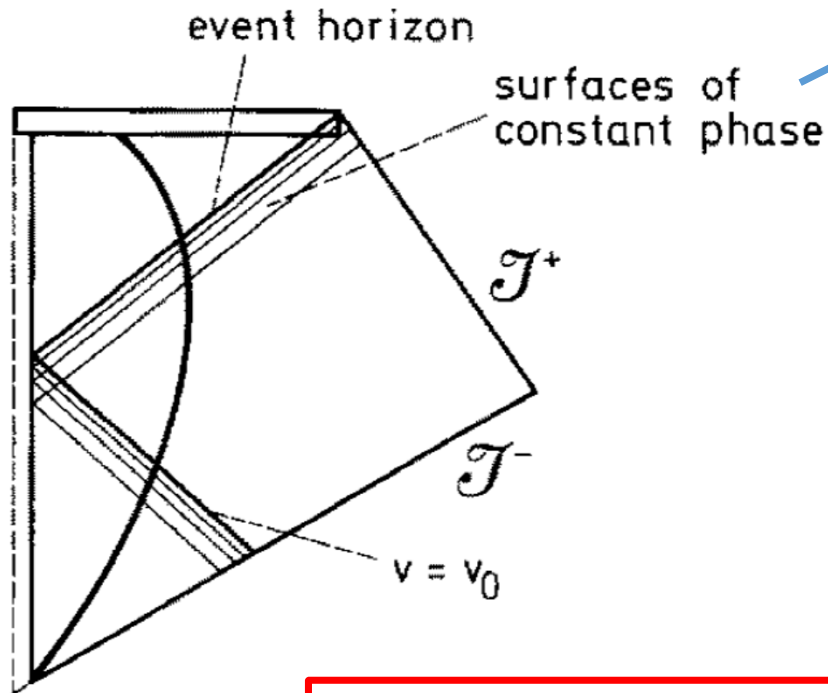
Thermal radiation that looks like





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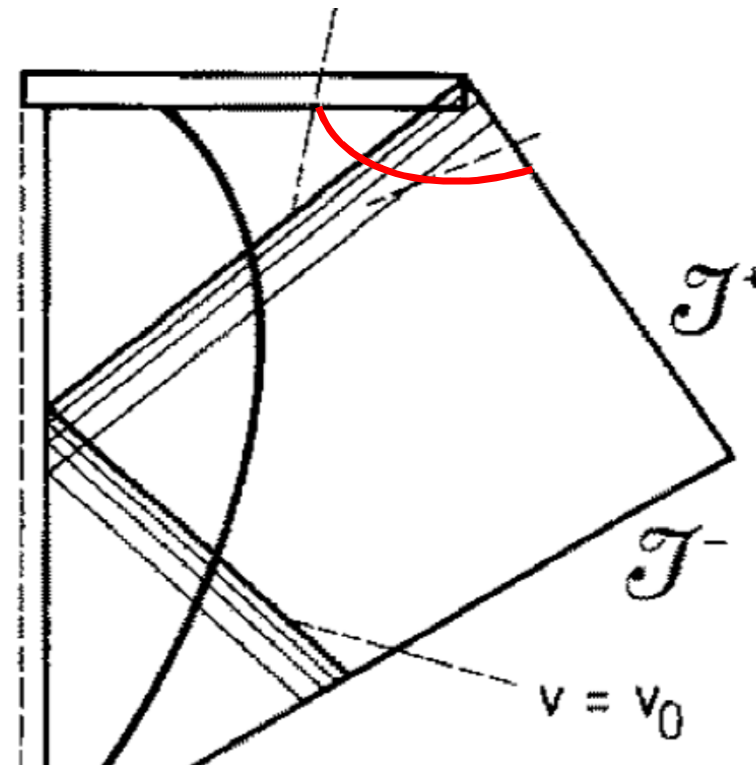
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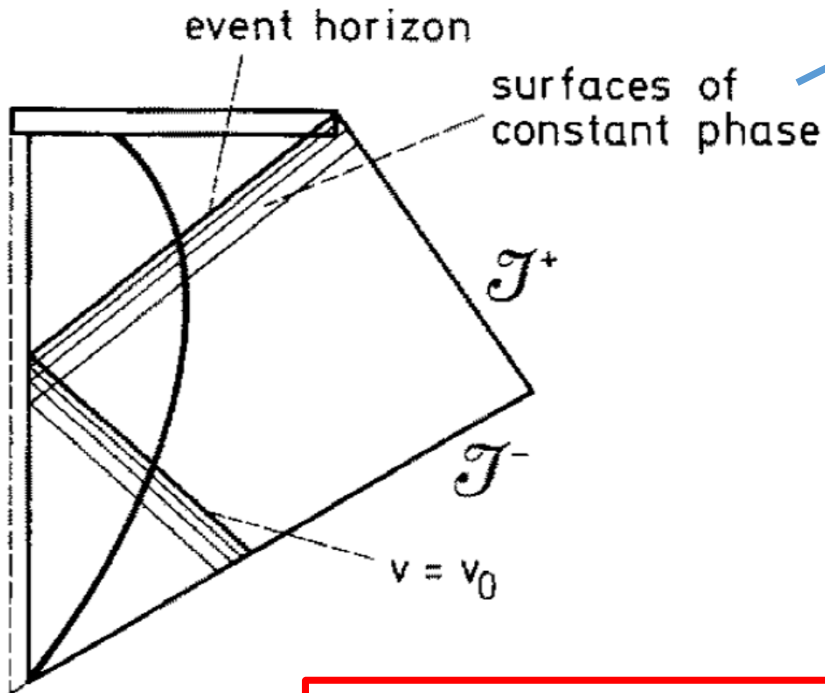
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Information is lost because  
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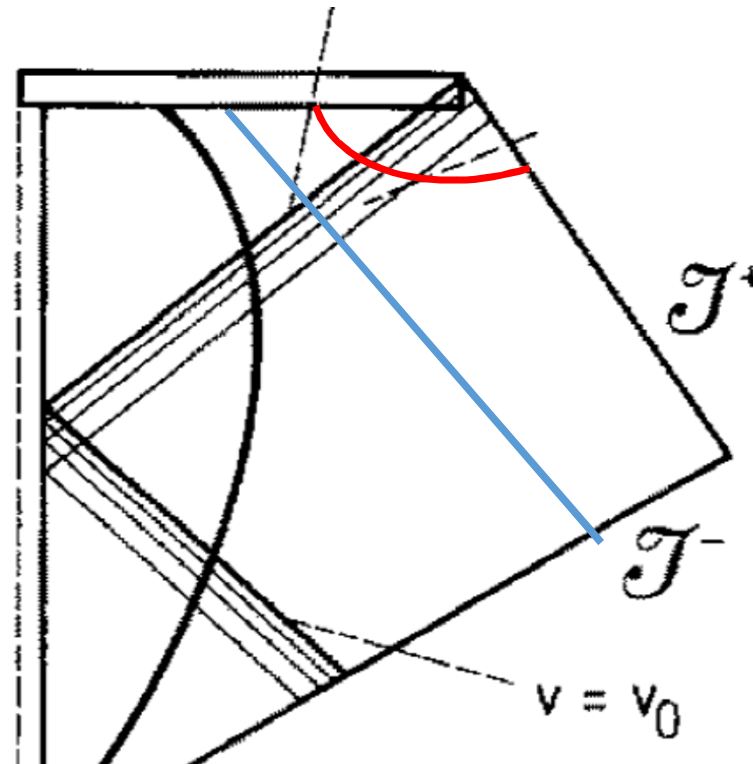
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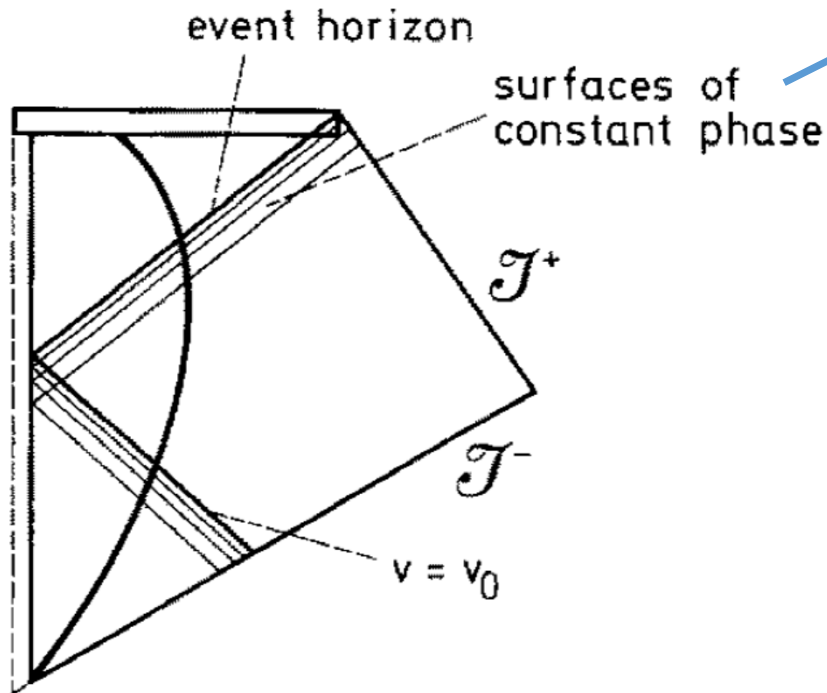
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Hawking calculation including non-pert stringy effects



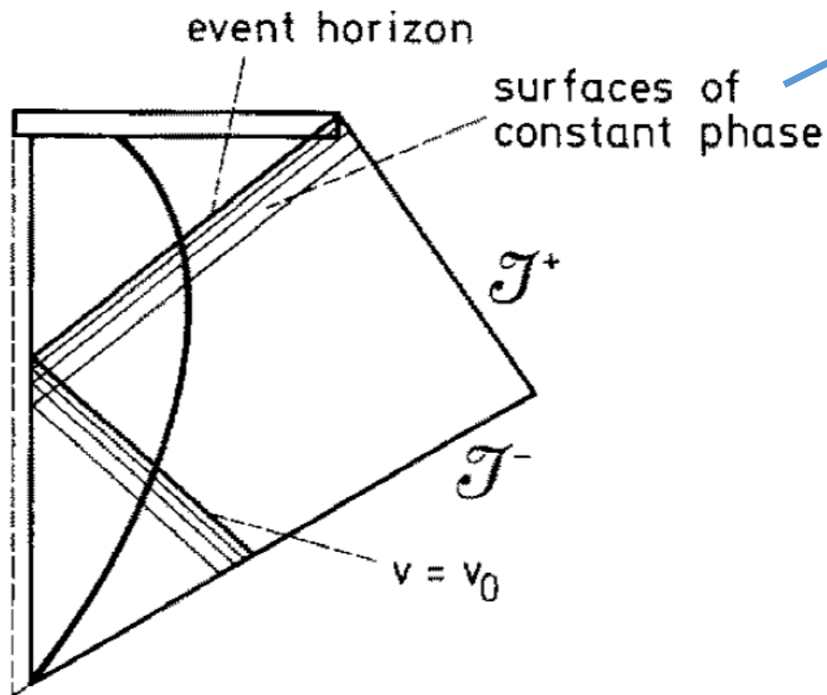
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Which at high E gives

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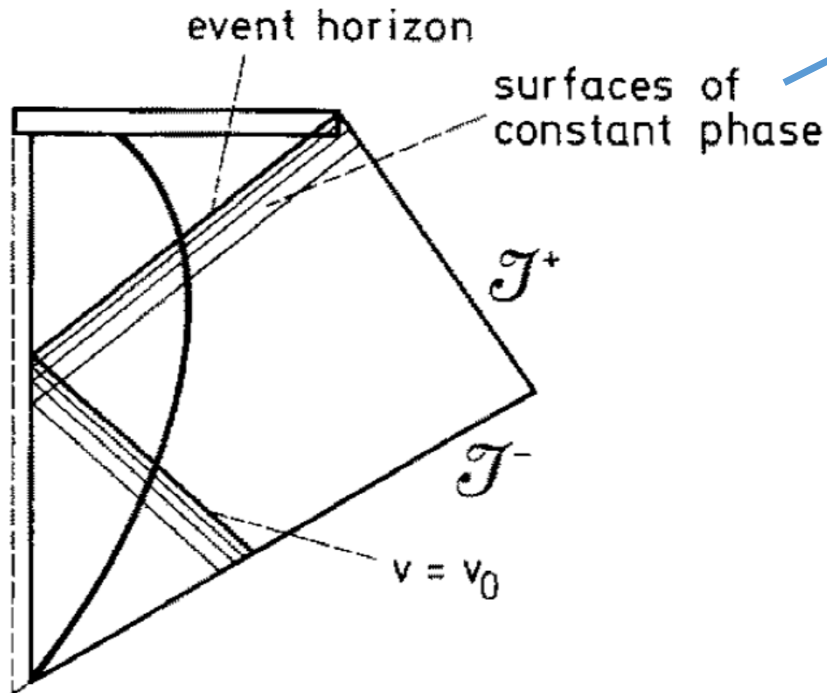
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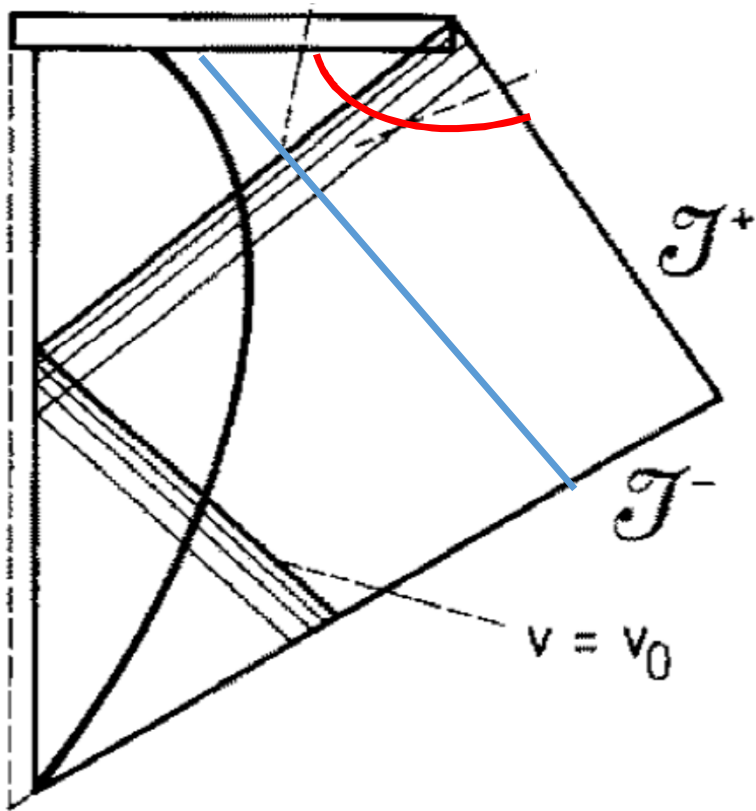
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Same temperature,  
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So instead of



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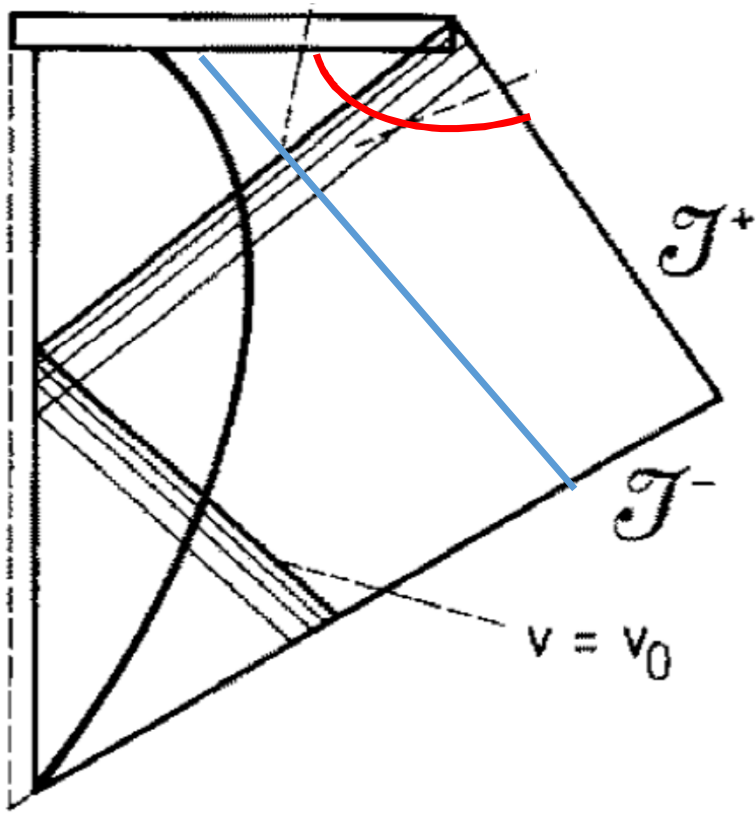
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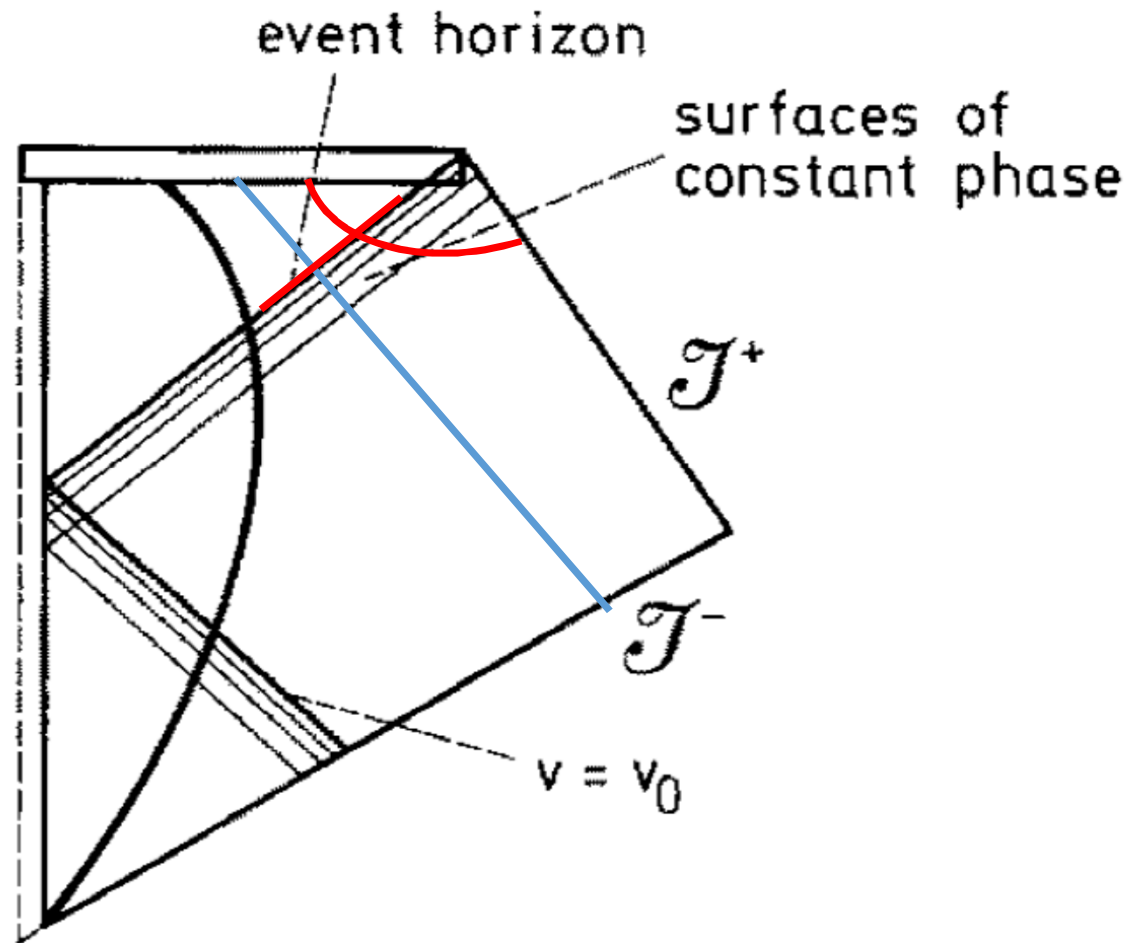
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## Summary:

1. Illustrated UV/IR mixing at the cigar geometry.
2. Argued it should play important role in BH info. puzzle.
3. Very much relevant for LST theory that Kutasov will discuss next.