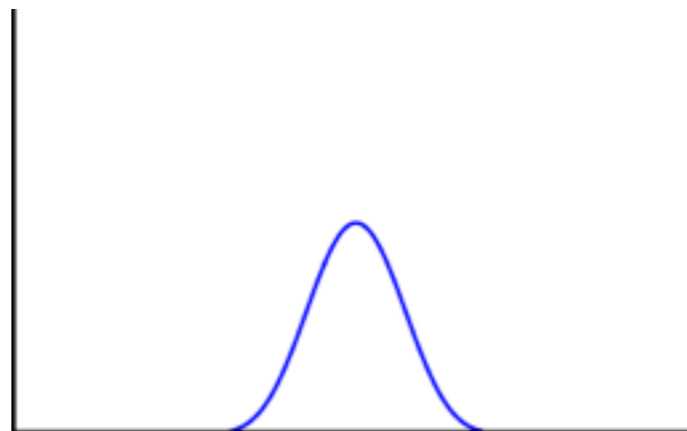


Turbulent strings in AdS/CFT

Takaaki Ishii
(University of Crete)

JHEP06(2015)086 [arXiv:1504.02190]
with Keiju Murata

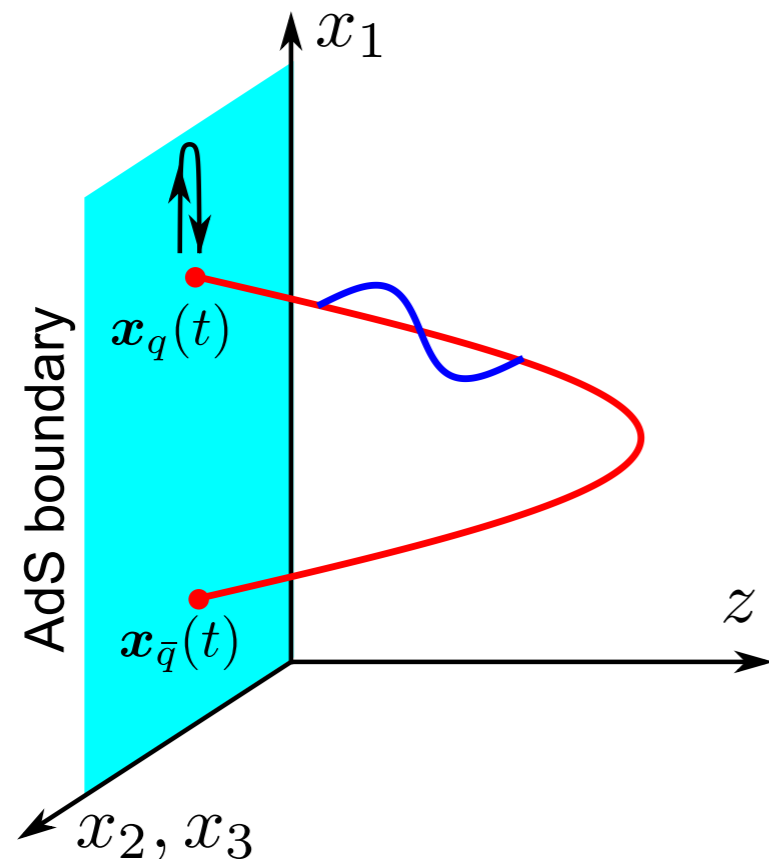


6 July 2015@Nafplion

What I will do

Perturb holographic quark-antiquark potential

Solve nonlinear time evolution



Motivations

- AdS turbulent instability
- Electric field quench on D7
[Hashimoto-Kinoshita-Oka-Murata]
- Dynamical meson melting
[TI-Kinoshita-Murata-Tanahashi]
- Cosmic strings in flat space

Contents

1. Static solution
2. Numerical setup
3. Results

Holographic quark potential

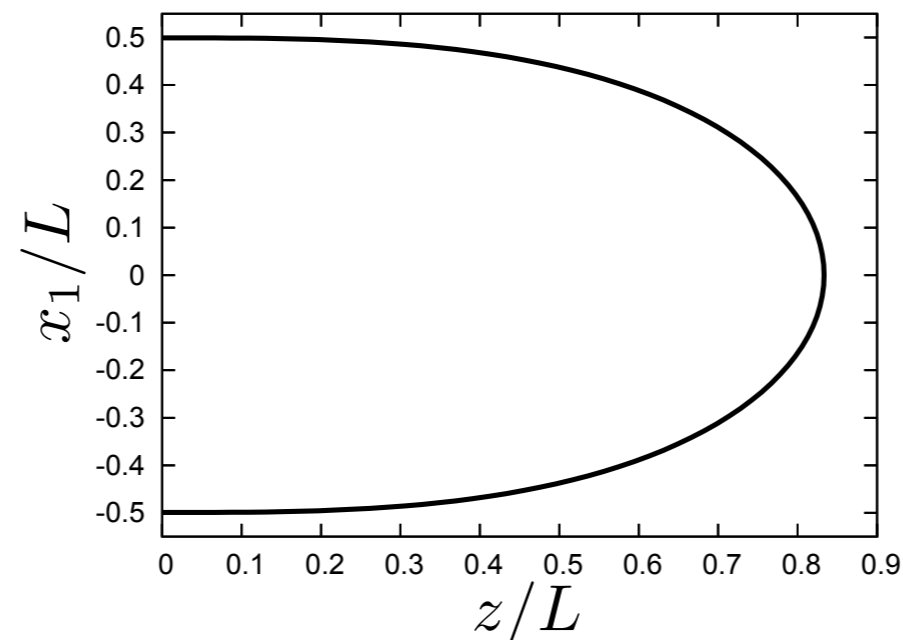
[Maldacena, Rey-Yee]

$$\text{AdS}_5 \times \text{S}^5 \quad ds^2 = \frac{\ell^2}{z^2} (-dt^2 + dz^2 + d\mathbf{x}^2) + \ell^2 d\Omega_5^2$$

Static gauge: $(\tau, \sigma) = (t, z)$

Target space embedding: $x_1 = X_1(z)$

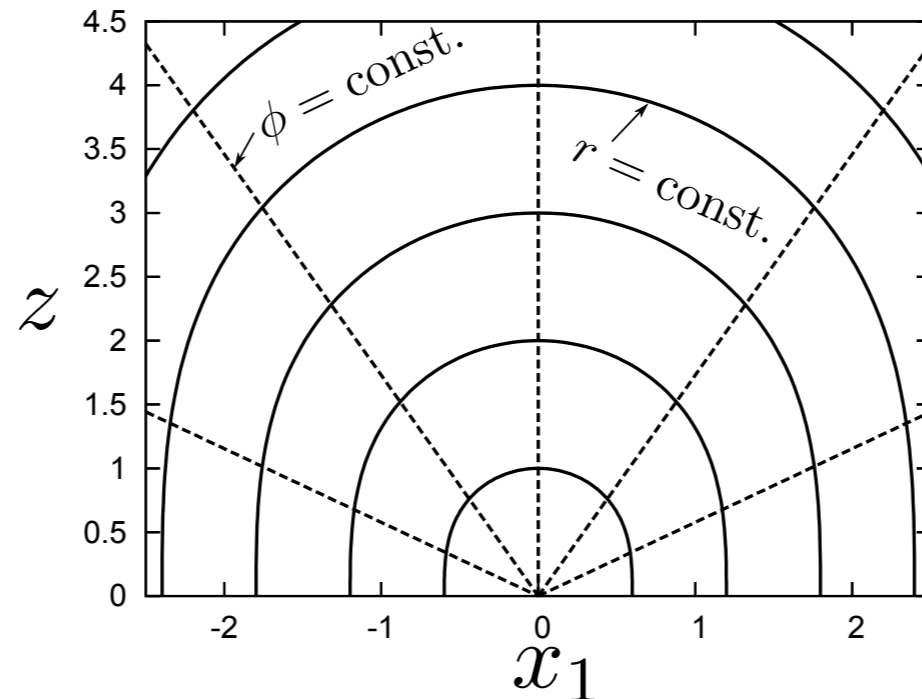
Solution with separation L



A polar parametrization

Use polar-like coordinates (r, ϕ) with $0 \leq \phi \leq \beta_0$

The static solution is at $r = \text{const.}$



$$z = r f(\phi) = r \operatorname{sn}(\phi; i)$$

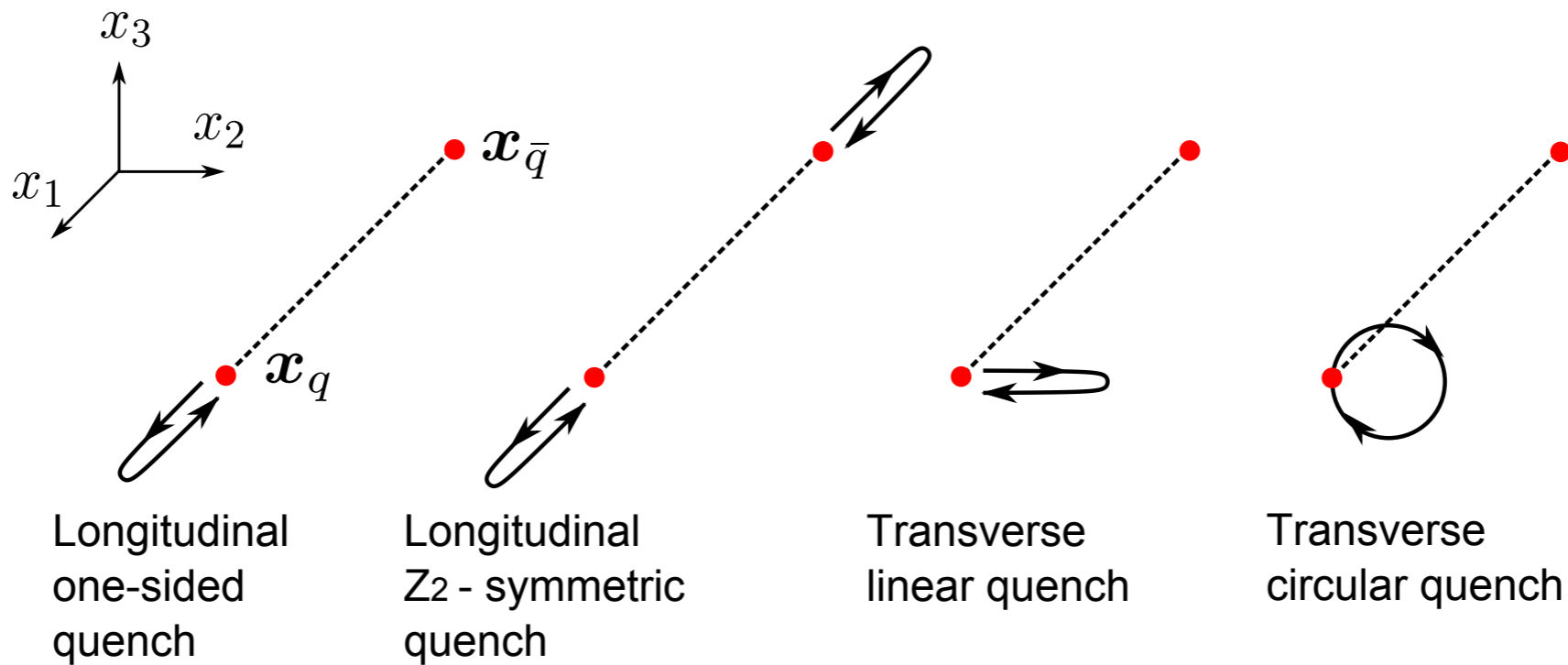
$$x_1 = r g(\phi) = r \begin{cases} \phi - E(\operatorname{sn}(\phi; i); i) + \Gamma_0 & (\phi \leq \beta_0/2) \\ \phi + E(\operatorname{sn}(\phi; i); i) - \Gamma_0 - \beta_0 & (\phi > \beta_0/2) \end{cases} \quad \begin{array}{l} \beta_0 \sim 2.622 \\ \Gamma_0 \sim 0.599 \end{array}$$

Contents

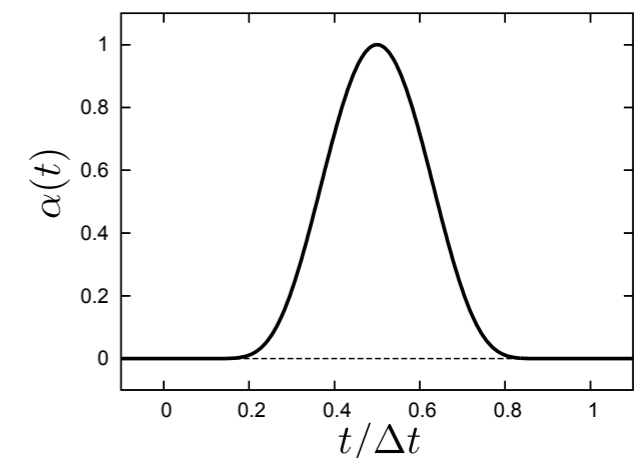
1. Static solution
- 2. Numerical setup**
3. Results

Perturb the string endpoints

We examine 4 representative patterns:



Quench profile: a Gaussian-ish compact C^∞ -function for $0 \leq t \leq \Delta t$



Worksheet double null coordinates

Induced metric $ds_{F1}^2 = -2\gamma_{uv}dudv$

Worksheet: u, v

Target space: $T(u, v), Z(u, v), X_{1,2,3}(u, v)$

$$\gamma_{uv} = \frac{\ell^2}{Z^2} (-T_{,u}T_{,v} + Z_{,u}Z_{,v} + \mathbf{X}_{,u} \cdot \mathbf{X}_{,v})$$

Equations of motion

$$T_{,uv} = \frac{1}{Z} (T_{,u}Z_{,v} + Z_{,u}T_{,v})$$

$$Z_{,uv} = \frac{1}{Z} (T_{,u}T_{,v} + Z_{,u}Z_{,v} - \mathbf{X}_{,u} \cdot \mathbf{X}_{,v})$$

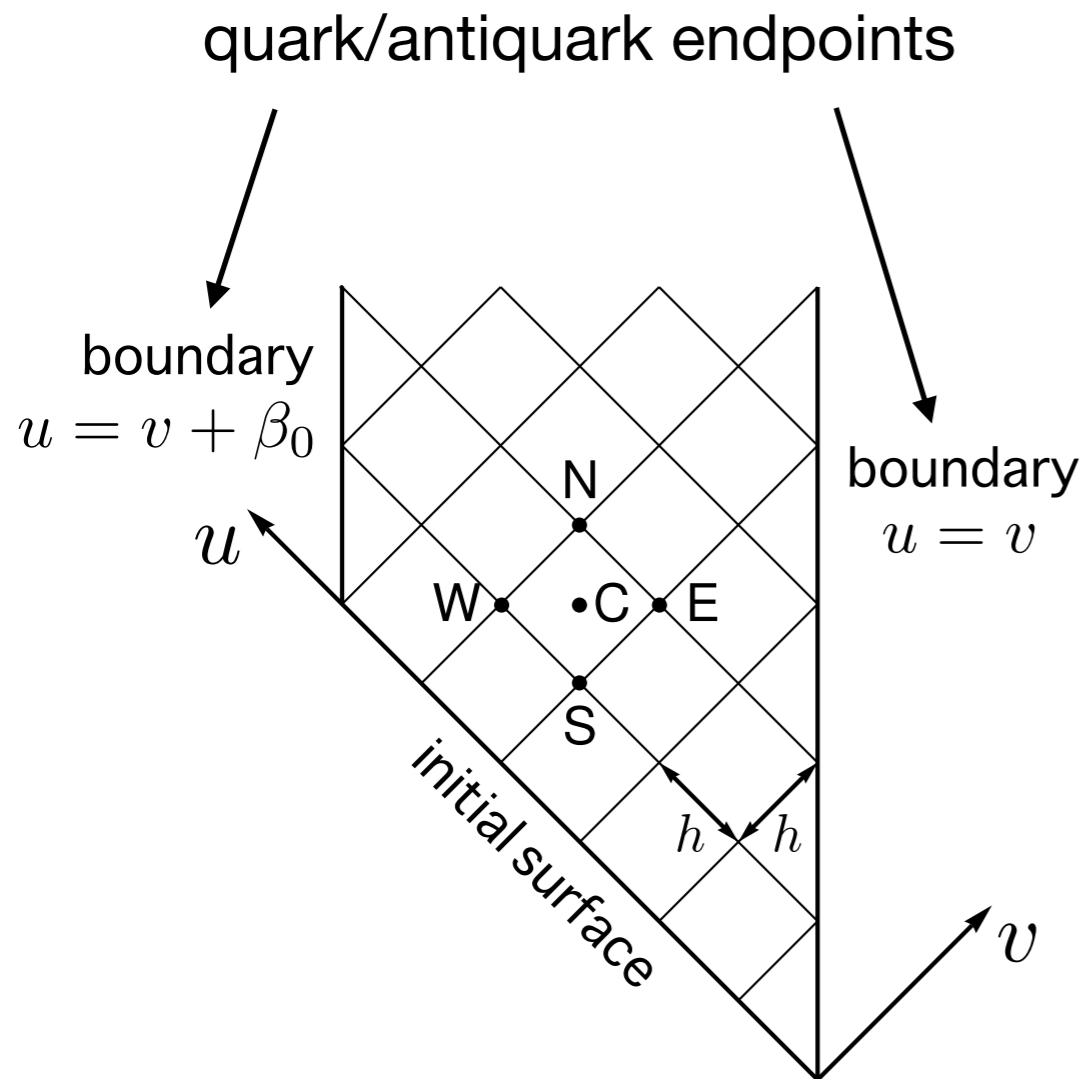
$$\mathbf{X}_{,uv} = \frac{1}{Z} (\mathbf{X}_{,u}Z_{,v} + Z_{,u}\mathbf{X}_{,v})$$

Constraints

$$\gamma_{uu} = \frac{\ell^2}{Z^2} (-T_{,u}^2 + Z_{,u}^2 + \mathbf{X}_{,u}^2) = 0$$

$$\gamma_{vv} = \frac{\ell^2}{Z^2} (-T_{,v}^2 + Z_{,v}^2 + \mathbf{X}_{,v}^2) = 0$$

Discretization



$O(h^2)$ central finite differential

$$\Psi_{,uv}|_C = \frac{\Psi_N - \Psi_E - \Psi_W + \Psi_S}{h^2}$$

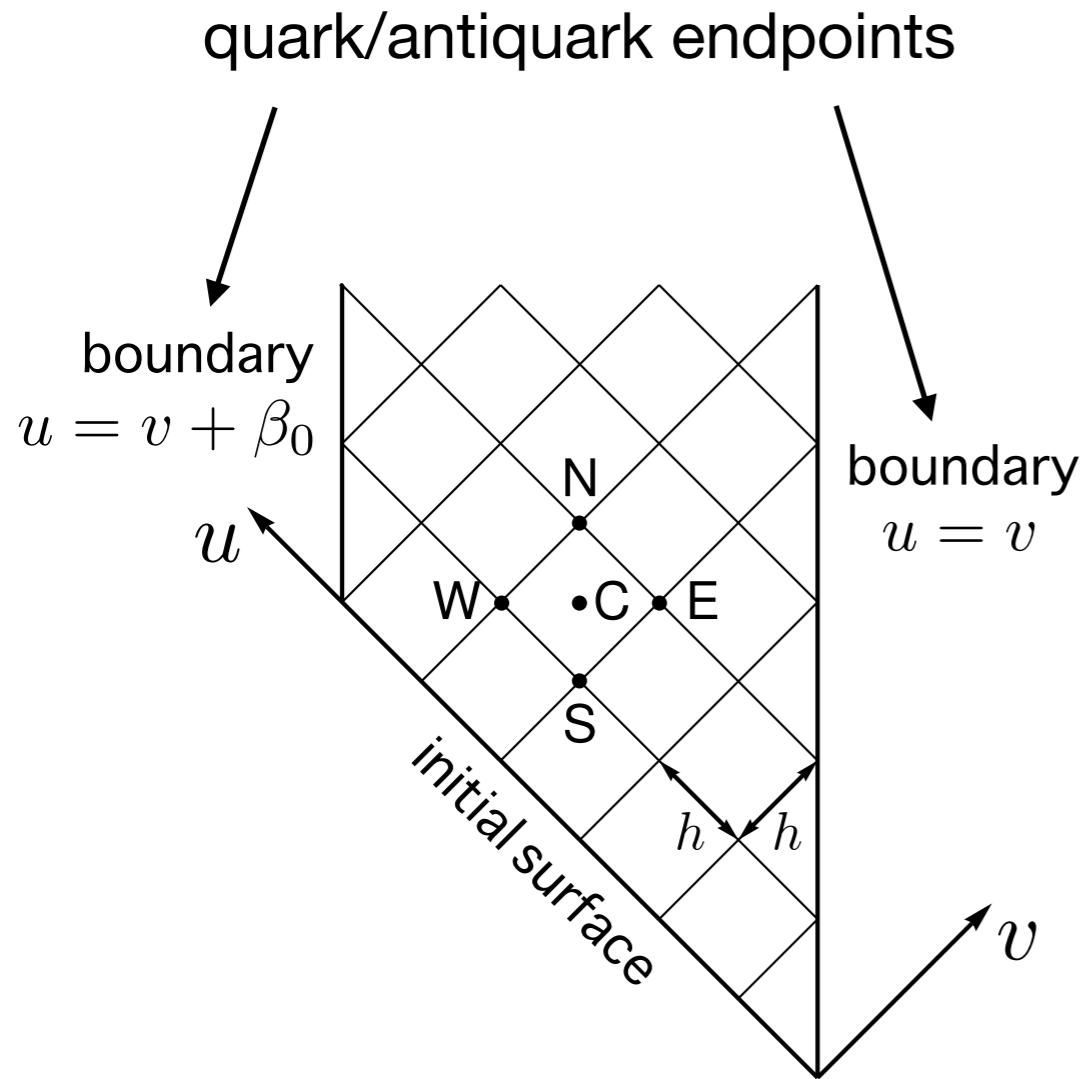
$$\Psi_{,u}|_C = \frac{\Psi_N - \Psi_E + \Psi_W - \Psi_S}{2h}$$

$$\Psi_{,v}|_C = \frac{\Psi_N + \Psi_E - \Psi_W - \Psi_S}{2h}$$

$$\Psi|_C = \frac{\Psi_E + \Psi_W}{2}$$

Compute N by using EWS data

Initial data ($v=0$)



- Gauge: $\phi = u$ at $v = 0$
- Static solution: $Z(u, 0)$, $\mathbf{X}(u, 0)$
- Constraint then determines $T(u, 0)$

$$\gamma_{uu} = \frac{\ell^2}{Z^2} (-T_{,u}^2 + Z_{,u}^2 + \mathbf{X}_{,u}^2) = 0$$

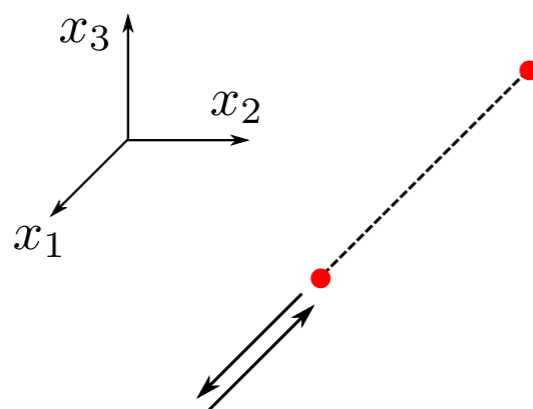
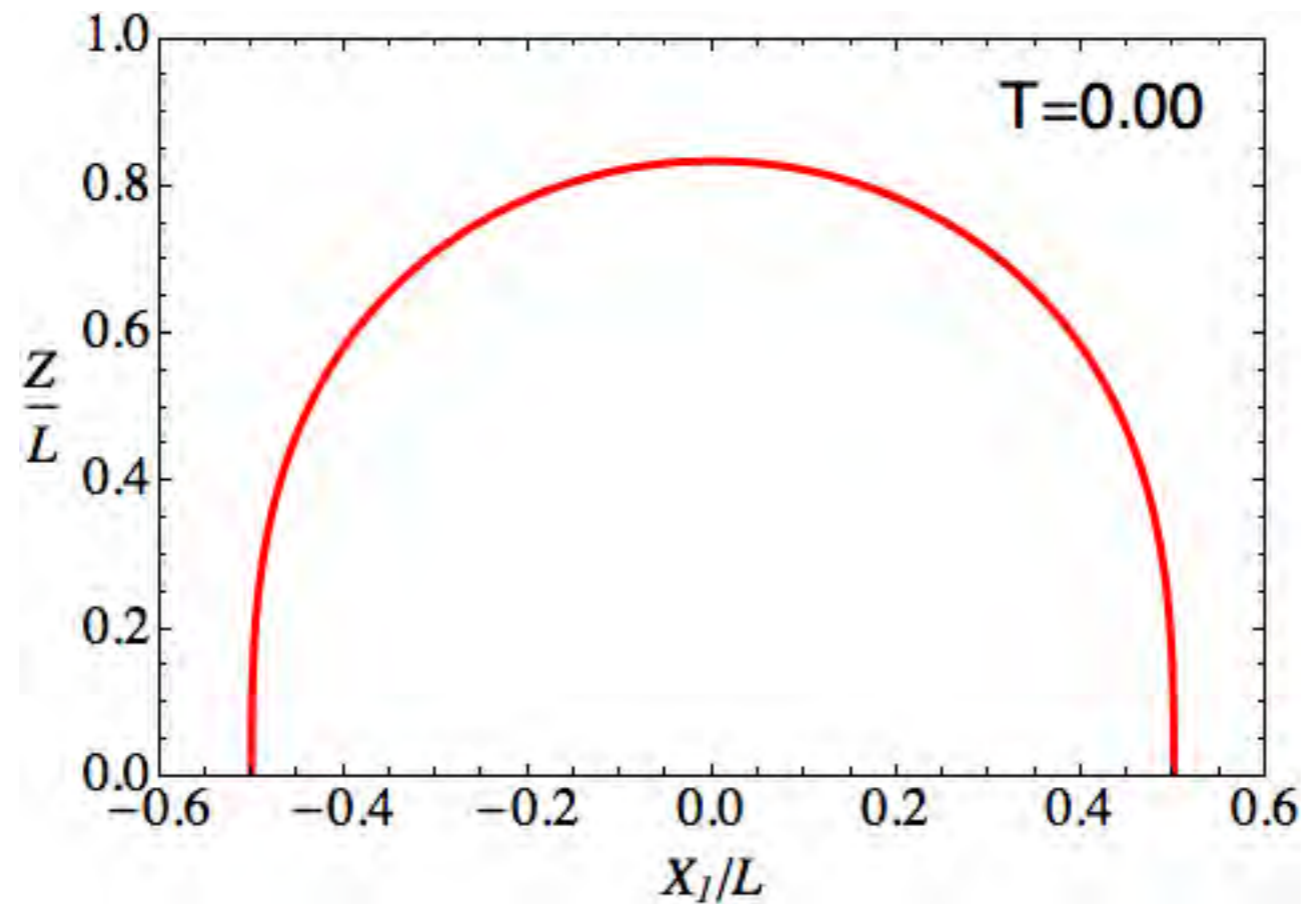
Result (analytic):

$$\begin{aligned} T(u, 0) &= z_0 u \\ Z(u, 0) &= z_0 f(u) \\ X_1(u, 0) &= z_0 g(u) \end{aligned} \quad \frac{L}{2} = z_0 \Gamma_0$$

Contents

1. Static solution
2. Numerical setup
- 3. Results**

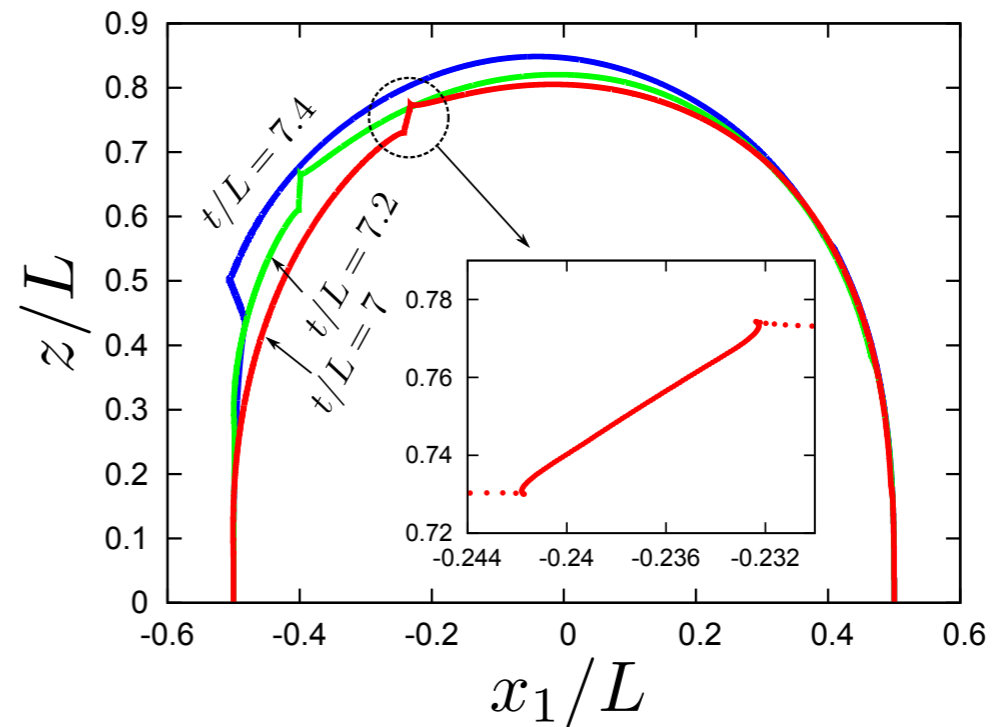
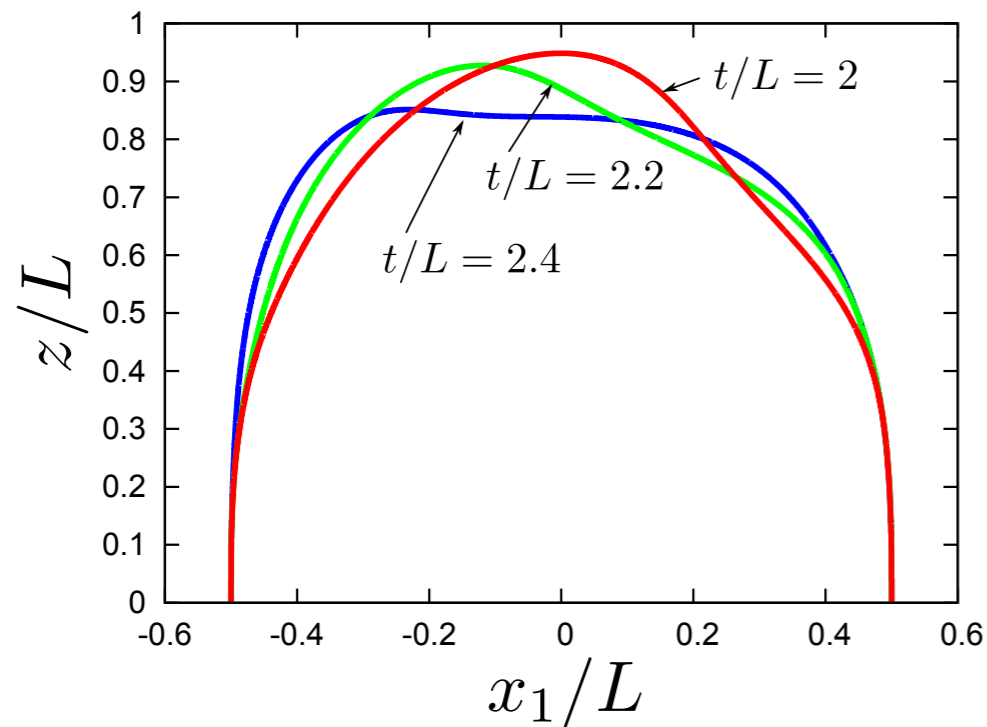
Longitudinal one-sided quench



$$\varepsilon=0.03, \Delta t/L=2$$

Two parameters:
Amplitude: $\varepsilon=\Delta x/L$
Duration: $\Delta t/L$

Cusp formation



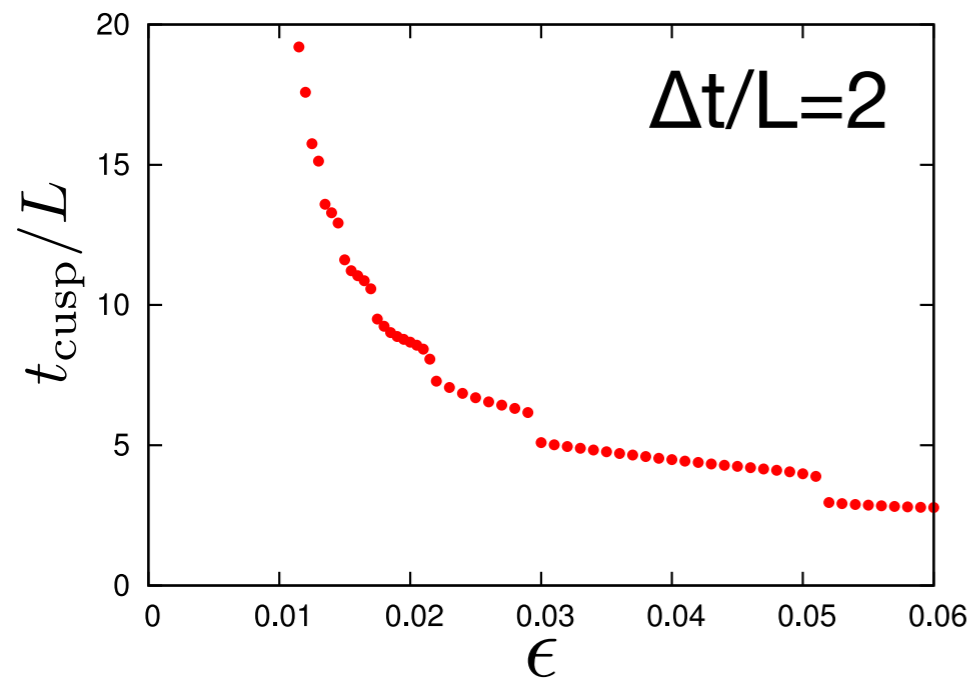
- Worldsheet $T(u,v)$, $Z(u,v)$, $\mathbf{X}(u,v)$ are all regular
- Target space plots show cusps
- Cusps are pair-created (here $t_{\text{cusp}}/L \sim 5$)

Analysis 1: Cusp detection

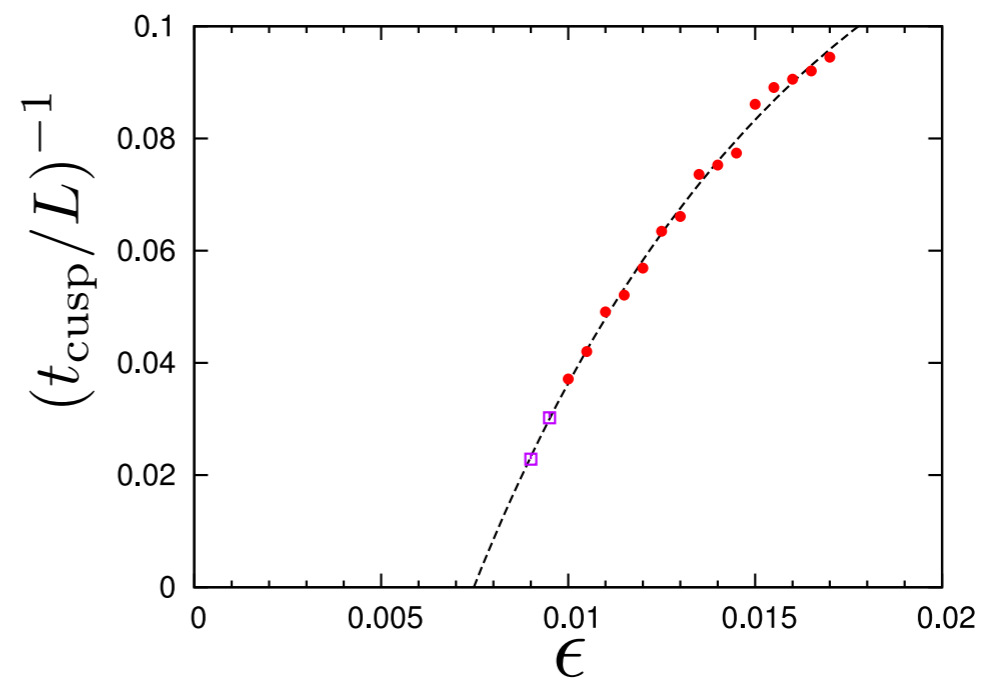
The conditions satisfied at the cusp:

$$J_z \equiv T_{,u}Z_{,v} - T_{,v}Z_{,u} = 0$$

$$J_i \equiv T_{,u}X_{i,v} - T_{,v}X_{i,u} = 0$$



flip



Cusp formation times
when amplitude ϵ is varied

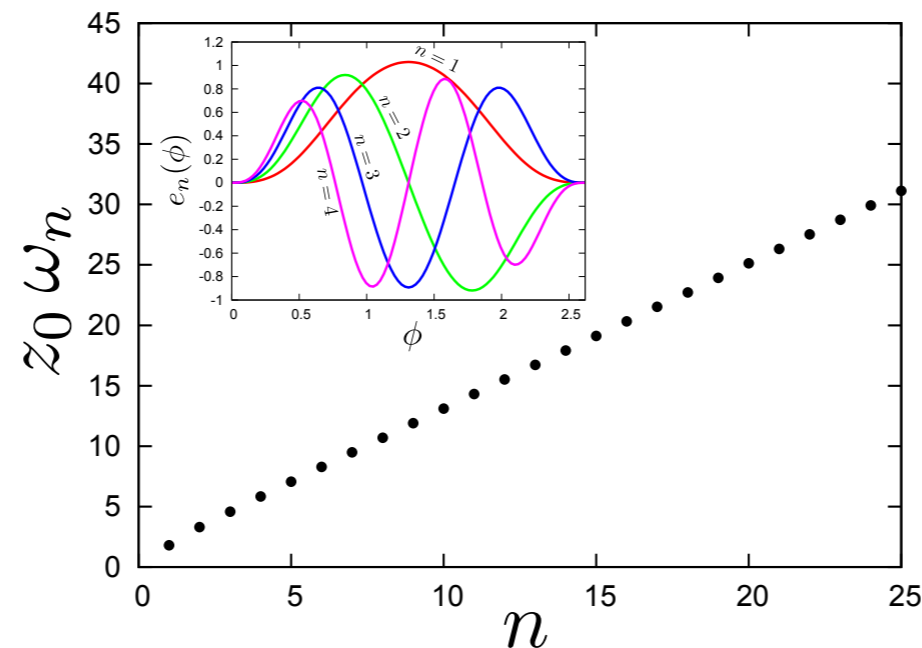
Minimal amplitude is nonzero
An extrapolation: $\epsilon_{\text{crit}} \sim 0.075$

Appendix: Linearized perturbations

[Callan-Guijosa, Klebanov-Maldacena-Thorn]

Linearized fluctuation $X_1 = X_{1(\text{stat})} + \chi_1$

Eigenvalues/functions $(\partial_t^2 + \mathcal{H})\chi_1 = 0$



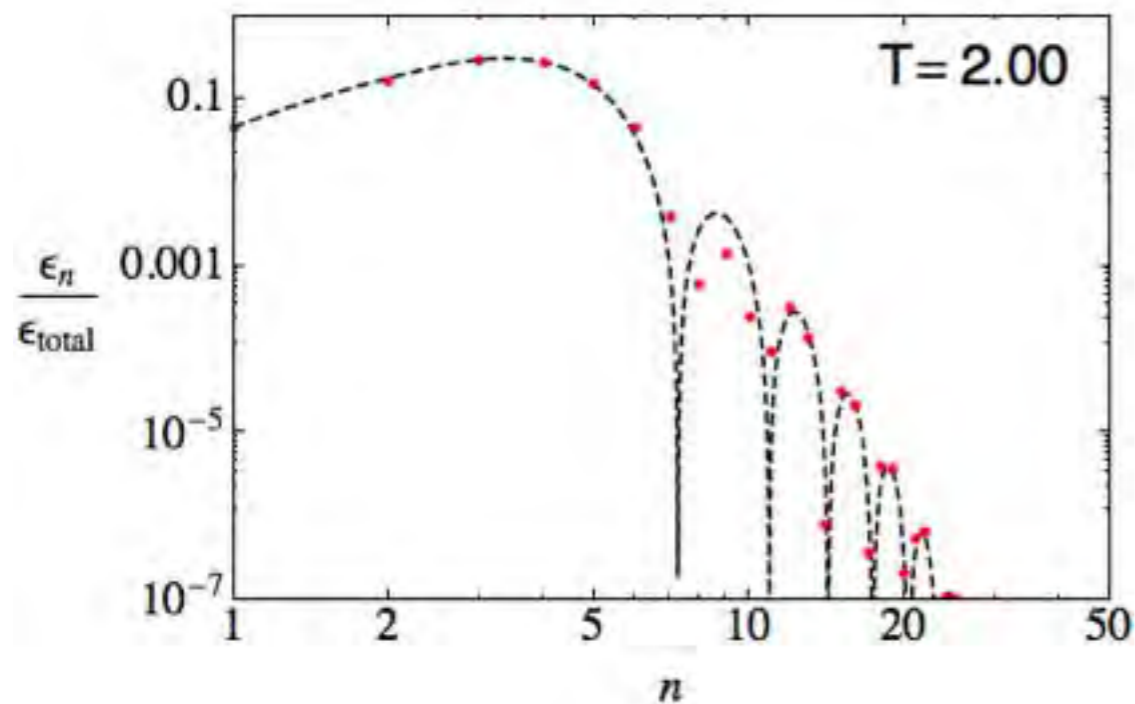
**Transverse modes (x_2, x_3) can be also computed

Analysis 2: Energy spectrum

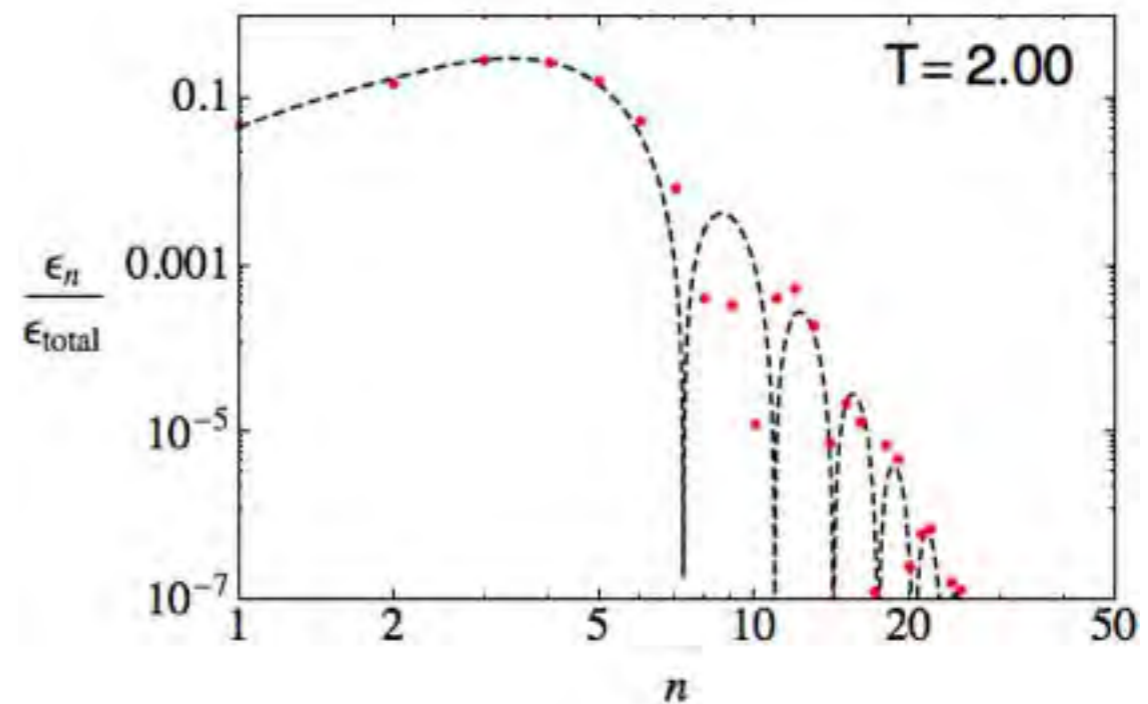
Decompose nonlinear solutions in linear eigenmodes e_n

$$\chi_1 = \sum_{n=1}^{\infty} c_n(t) e_n(\phi) \quad \varepsilon_n(t) = \frac{\sqrt{\lambda} z_0}{4\pi} (\dot{c}_n^2 + \omega_n^2 c_n^2)$$

Log-log plots



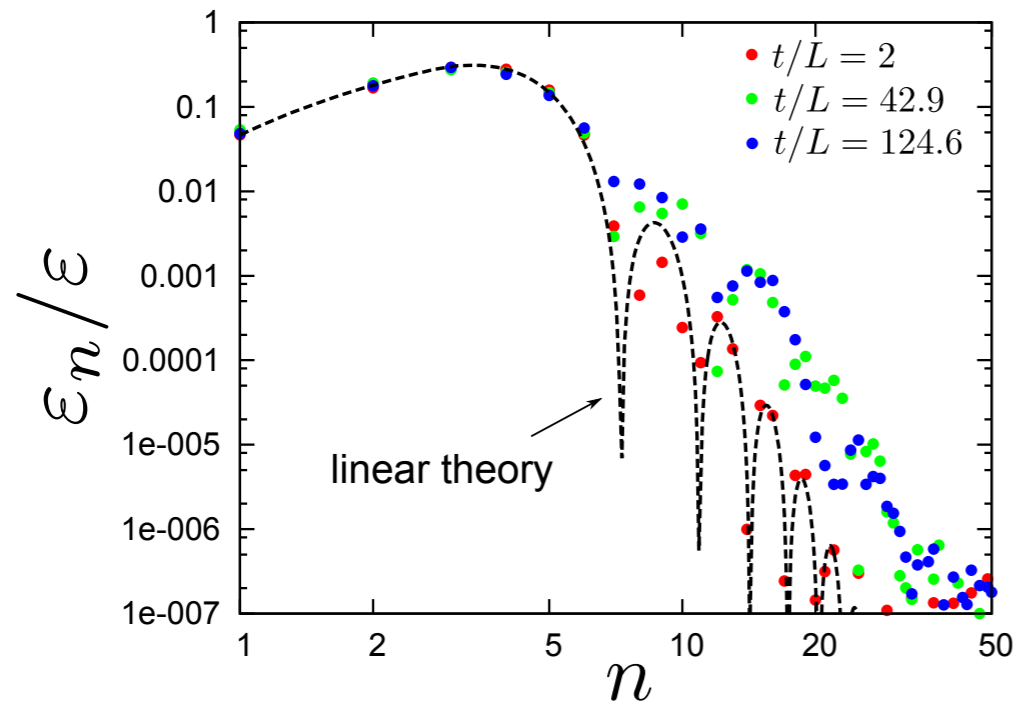
$\varepsilon=0.005, \Delta t/L=2$ (no cusp)



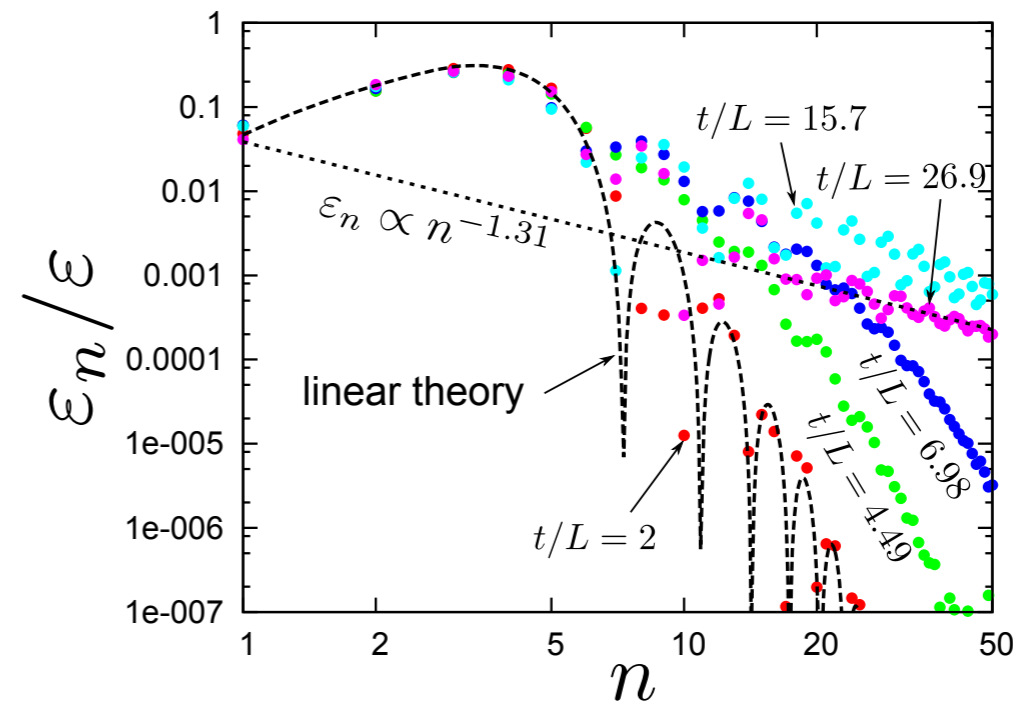
$\varepsilon=0.01$ (cusps $T \sim 27$)

**Dashed lines: from linearized action

Energy cascade



$\varepsilon=0.005, \Delta t/L=2$ (no cusp)



$\varepsilon=0.01$ (cusps $T \sim 27$)

Toward cusp formation: **direct energy cascade**

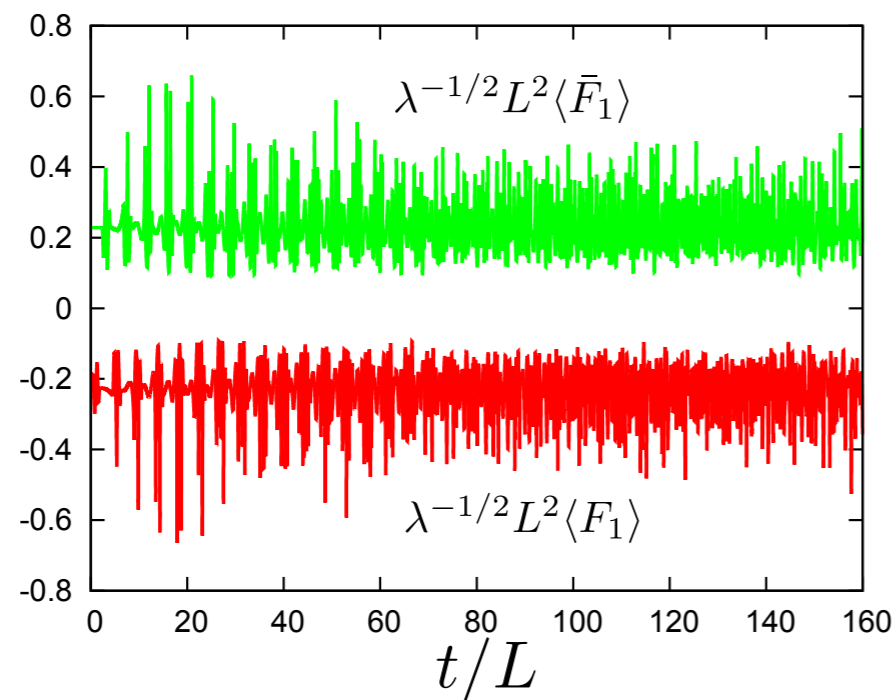
- A fit of right panel at $T \sim 27$: $\varepsilon_n \sim n^{-1.3}$

No cusp (too small ε): no clear power law

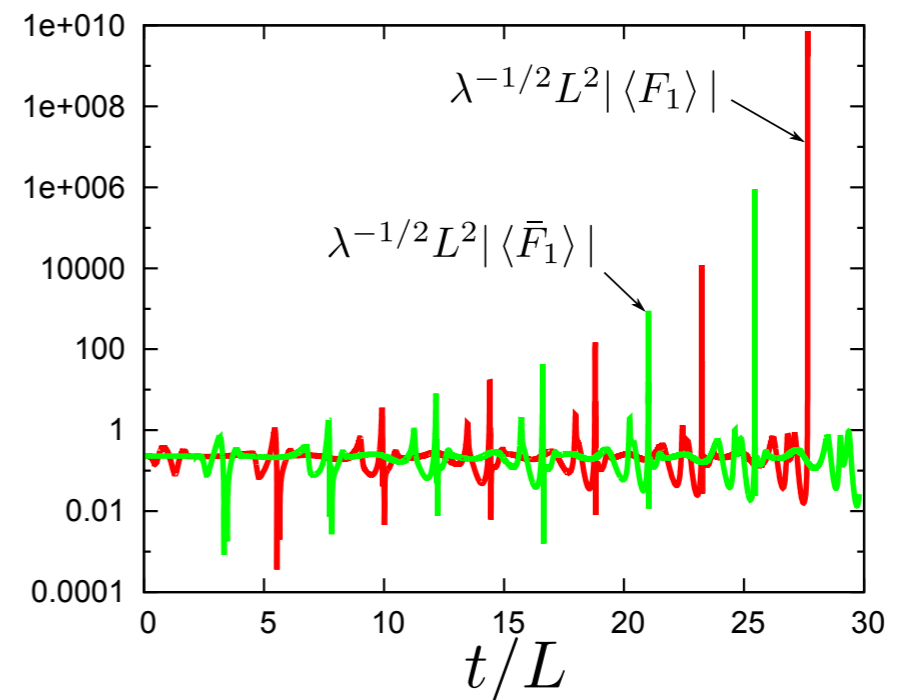
Analysis 3: Forces on the endpoints

$$\langle \mathbf{F}(t) \rangle = \frac{\delta \mathcal{S}_{\text{on-shell}}}{\delta \mathbf{x}_q}$$

Force diverges when a cusp reaches the boundary



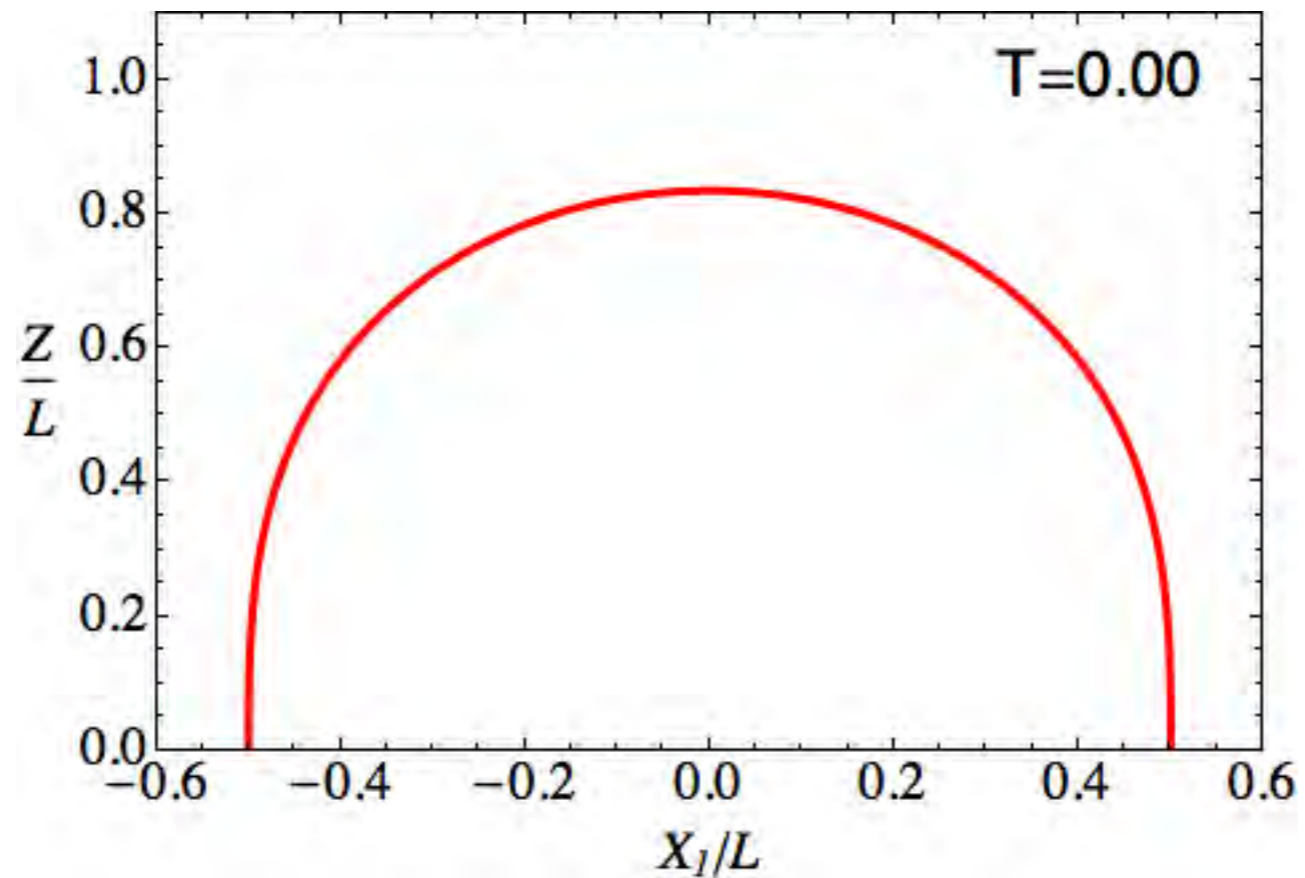
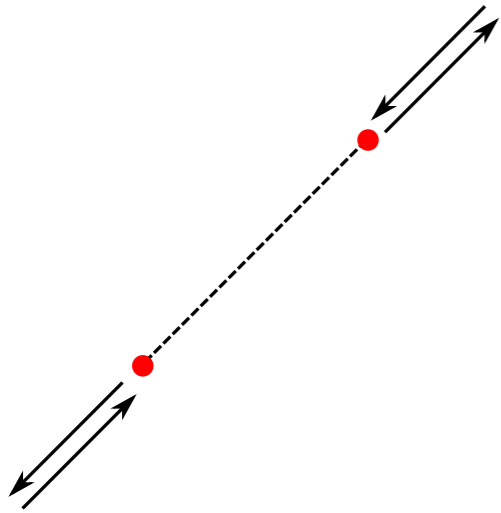
$\varepsilon=0.005, \Delta t/L=2$ (no cusp)



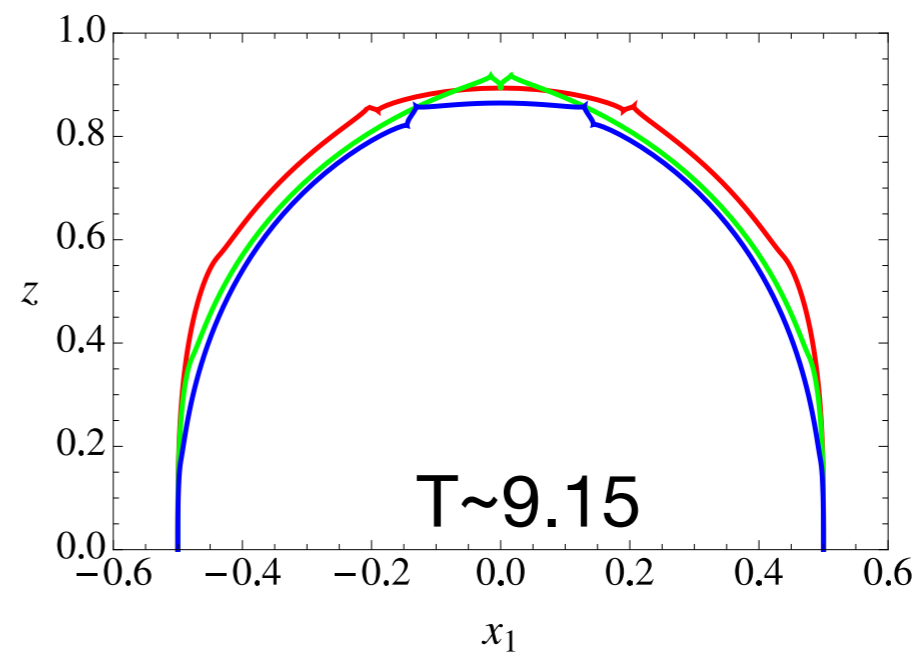
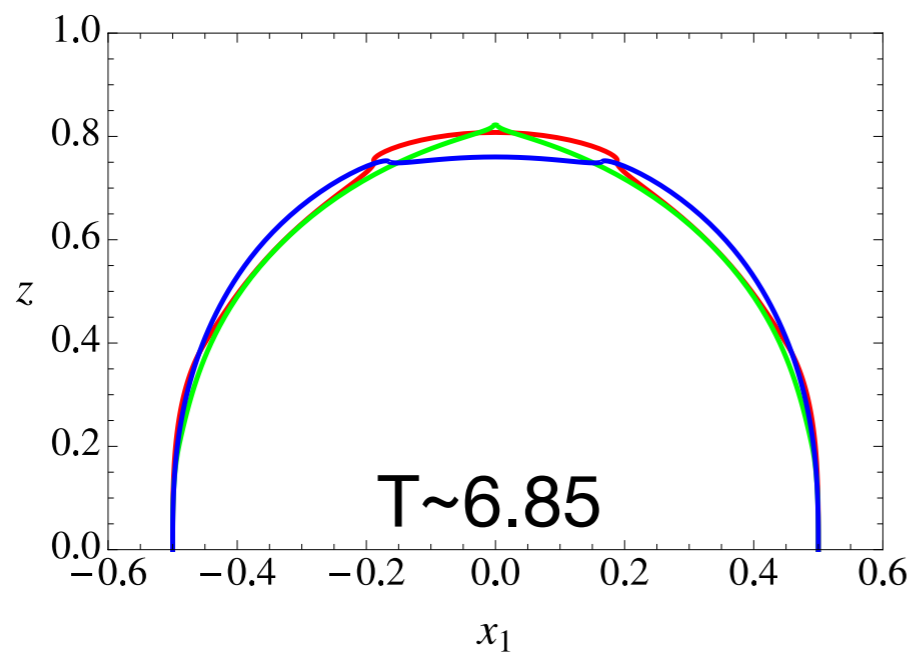
$\varepsilon=0.01$ (cusps $T \sim 27$)

***Red: $x=L/2$, green: $x=-L/2$

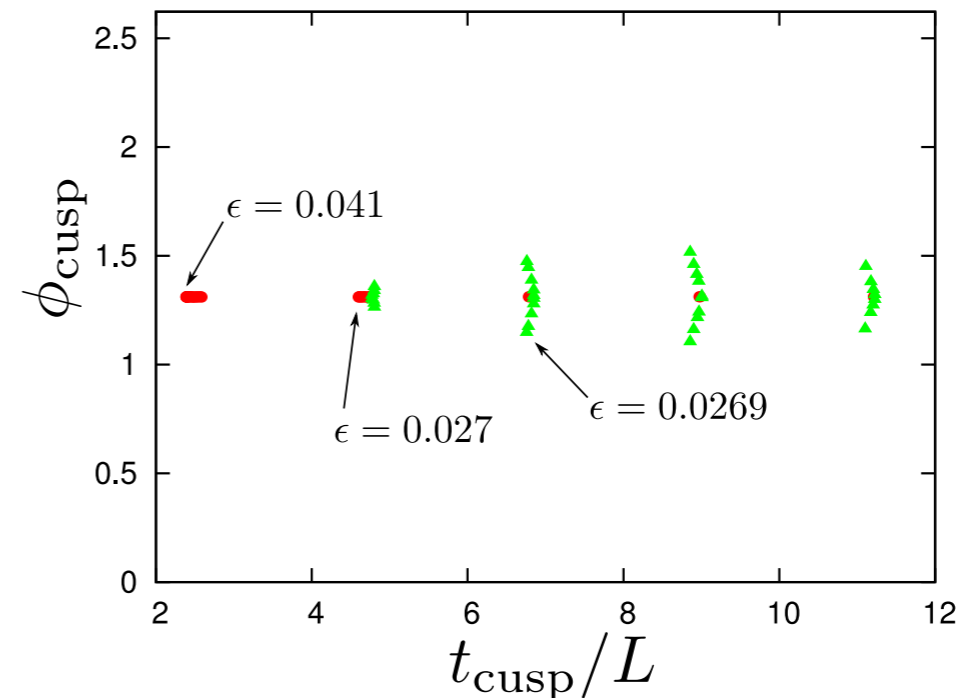
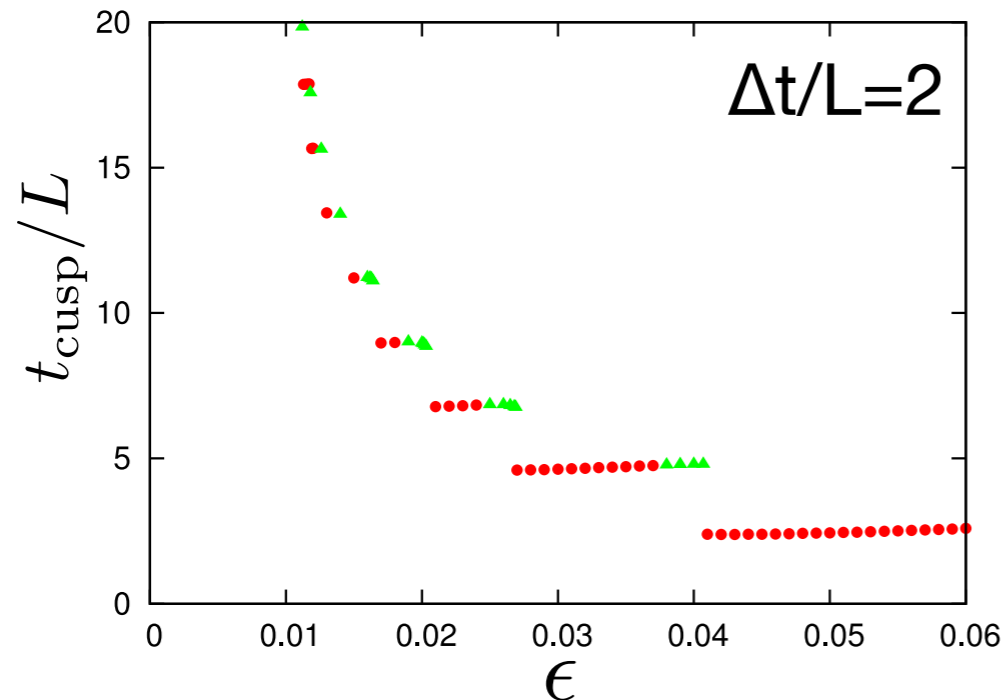
Z₂-symmetric quench



$\varepsilon=0.025$
 $\Delta t/L=2$

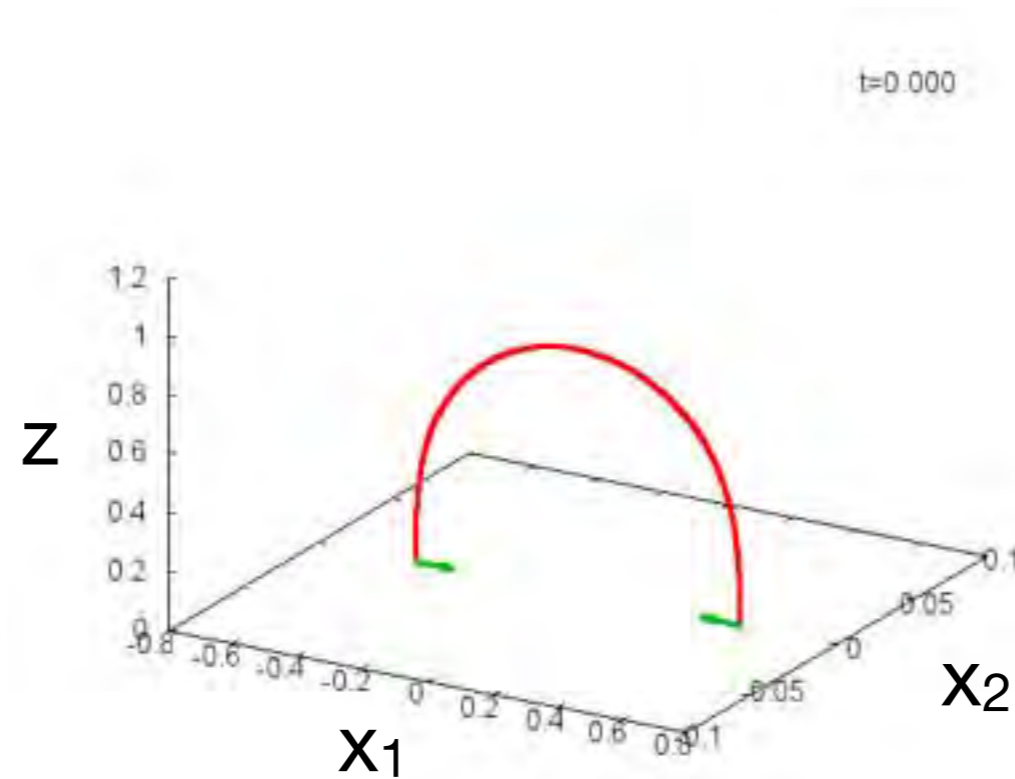
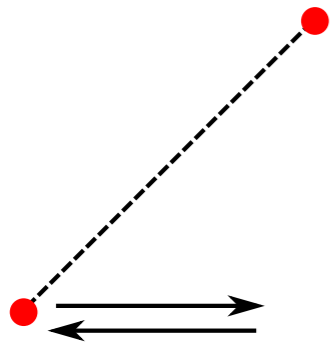


Cusp formation in the Z_2 -case



- Formation times are discretized by wave collisions
- First cusp formations on wave collisions (red ●).
The cusps are pair-created and annihilated.
- Traveling cusps can be formed first (green ▲)

Transverse linear quench

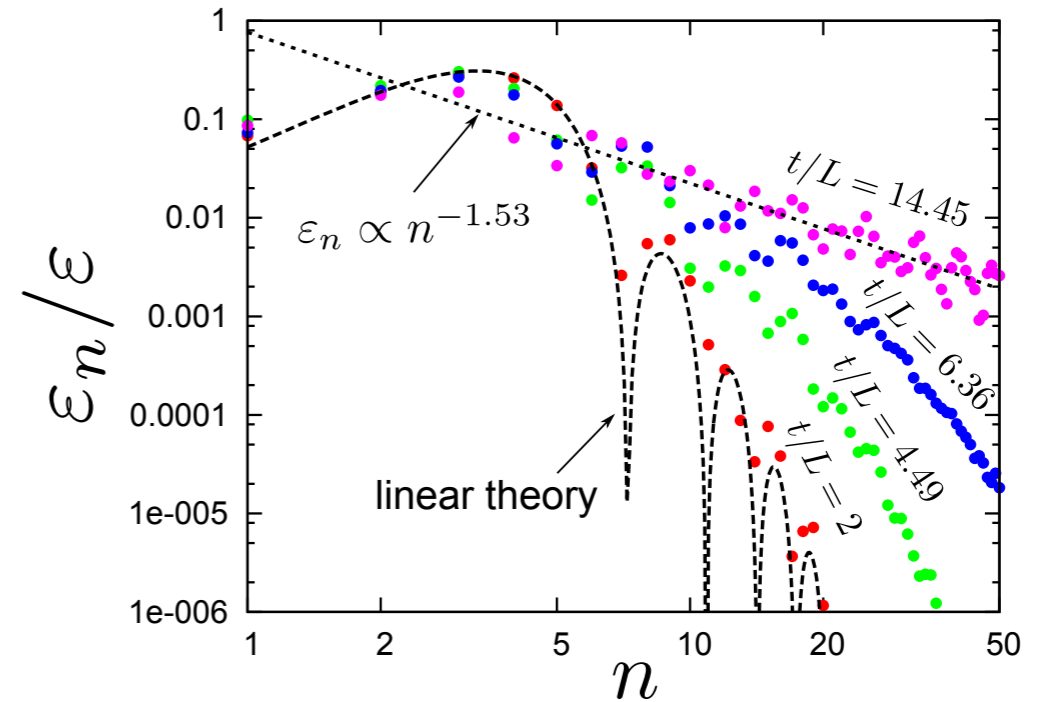
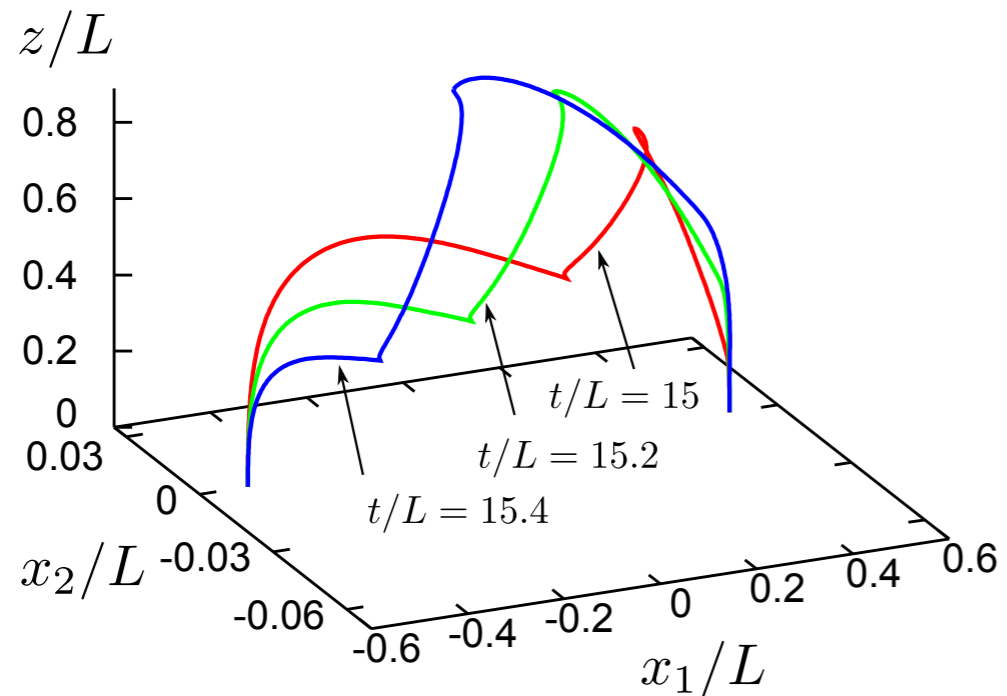


$$\varepsilon=0.03, \Delta t/L=2$$

***Green arrows: forces

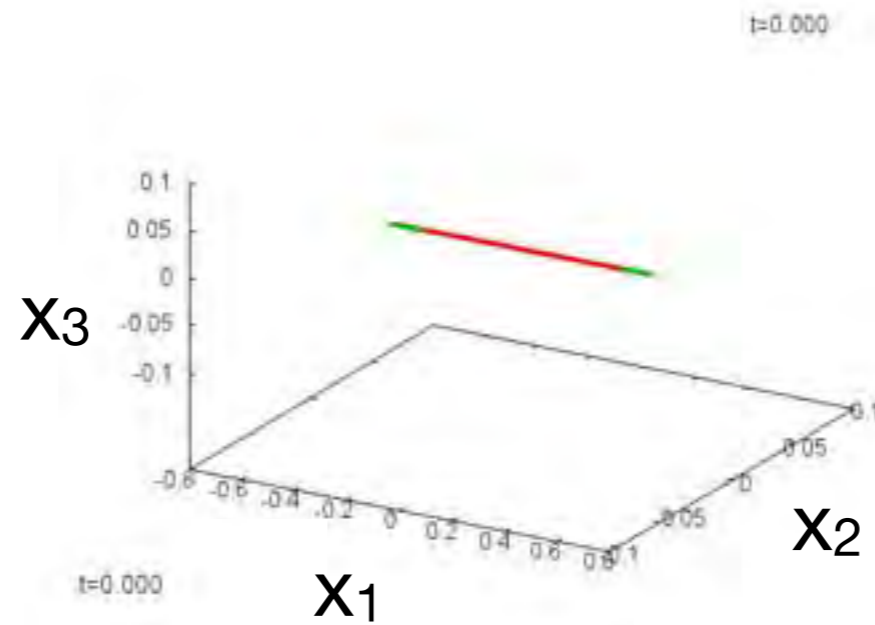
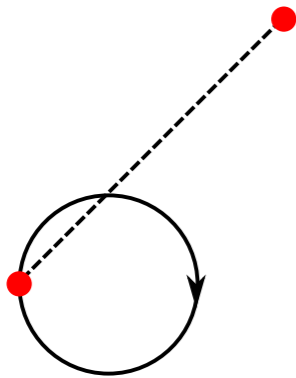
String oscillates in 1+3 dim (t, z, x_1, x_2)

Transverse linear quench

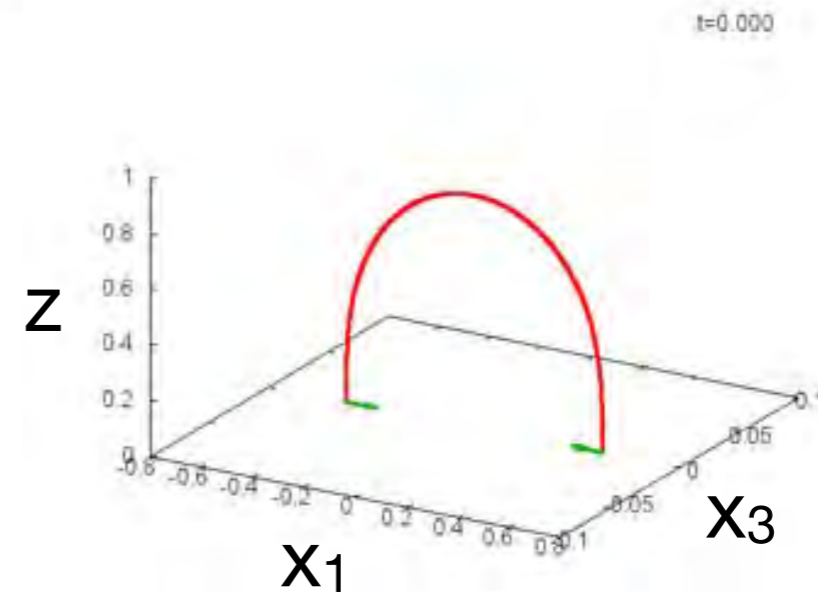
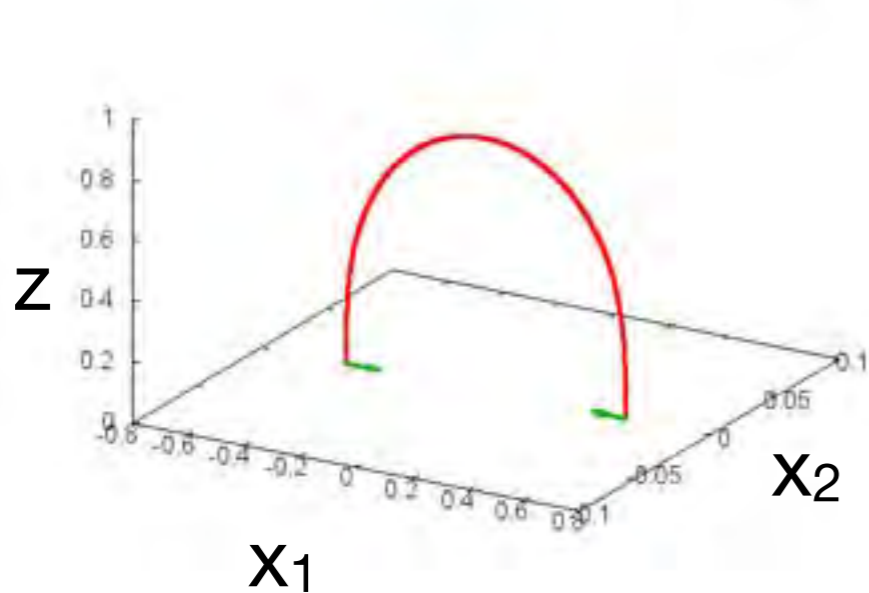


- Cusp formation at $T \sim 14.5$ ($\varepsilon_n \sim n^{-1.5}$)
- Direct energy cascade

Transverse circular quench

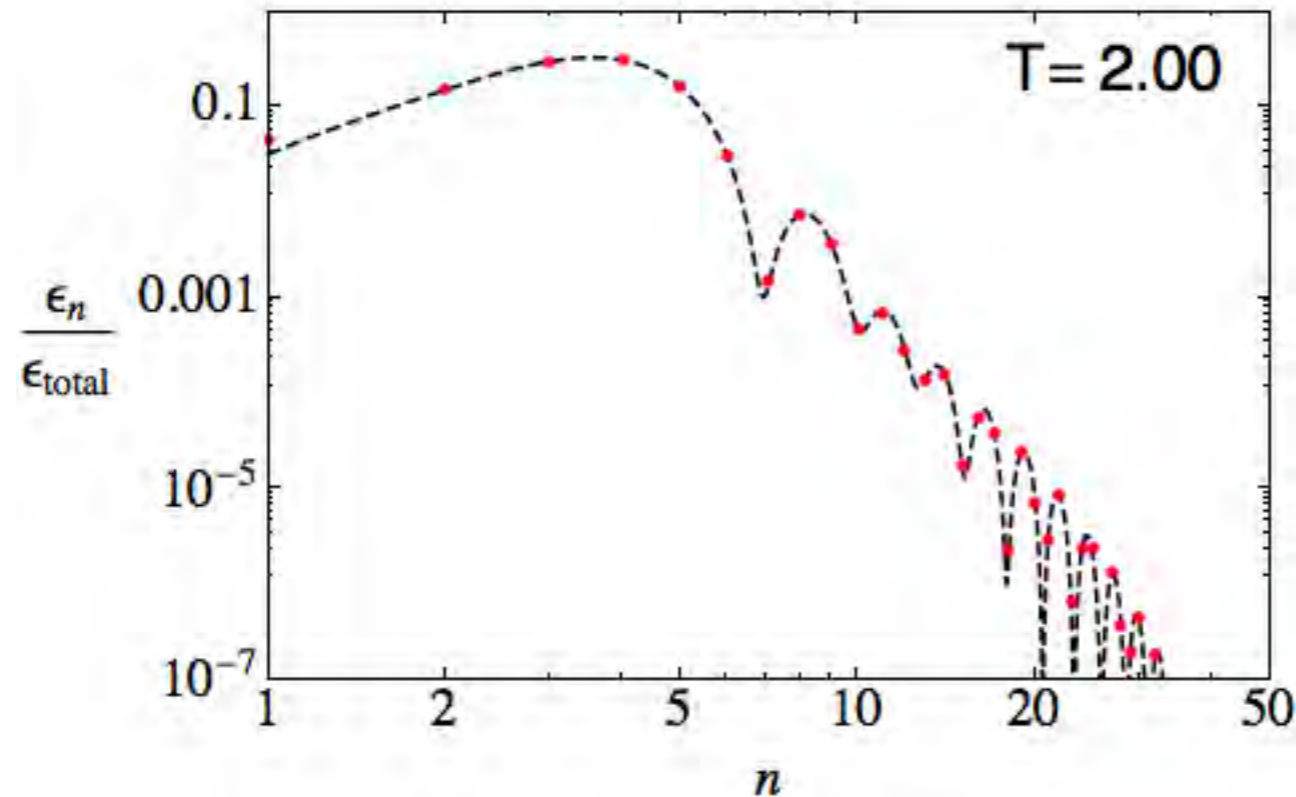


$\epsilon=0.02, \Delta t/L=2$



String oscillates in all 1+4 dim (t, z, x_1, x_2, x_3)

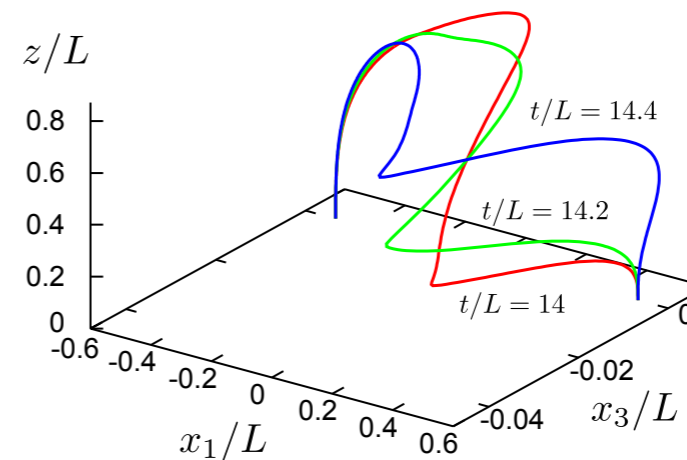
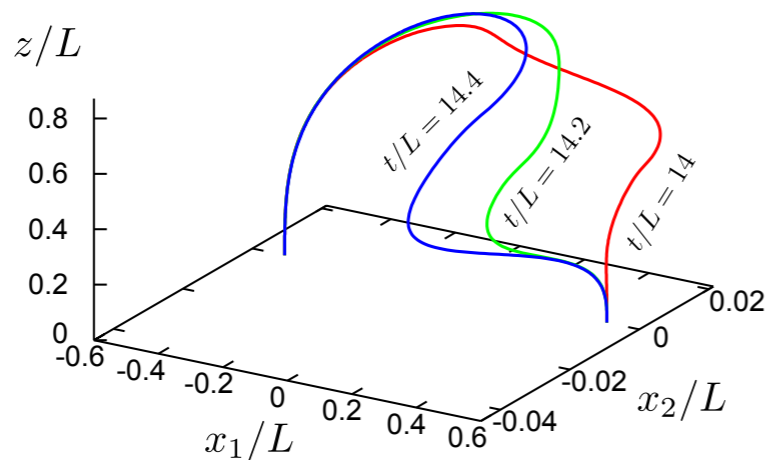
Transverse circular energy spectrum



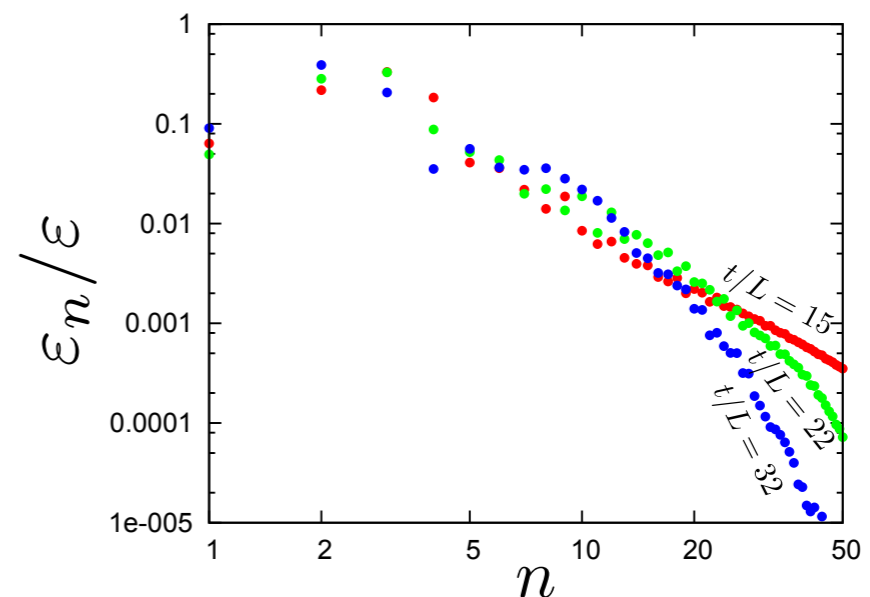
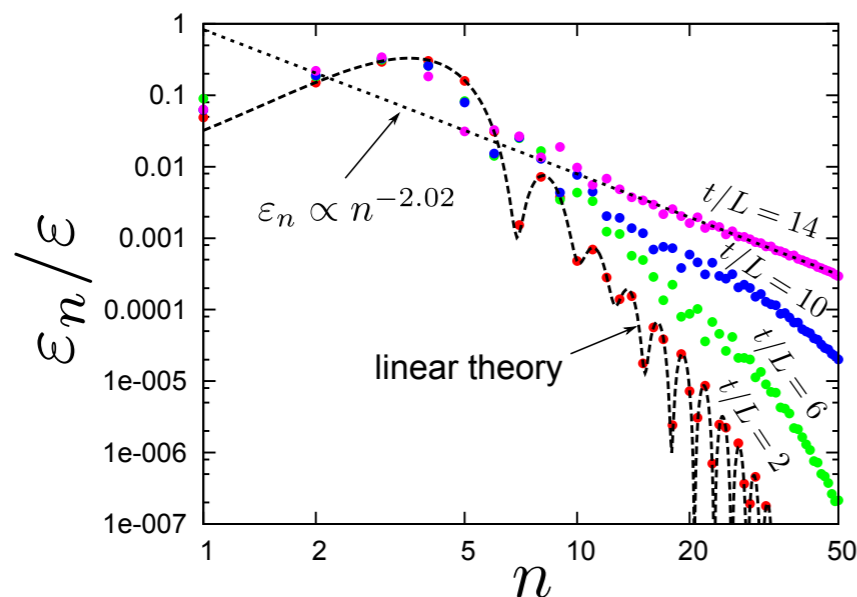
No cusp: no sustaining power law

c.f.) Probability of cusp formation is zero if $\text{dim} > 4$

Transverse circular quench



Cuspy, but not real cusps



Direct cascade $\rightarrow (\epsilon_n \sim n^{-2}) \rightarrow$ **inverse cascade**

Summary

We computed nonlinear dynamics of the quark-antiquark fundamental string in AdS

- Cusps and turbulent behavior in $\leq 1+3$ dim
- No cusp and direct/inverse cascades in $1+4$ dim

Boundary interpretation: Nonlinearity in YM flux tube might squeeze/de-squeeze the energy

Discussion

Backreaction may be necessary at (near) cusps

- Curvature diverges at the cusps
- AdS gravitational wave bursts?
- Boundary: gluon emission from the flux tube?

Future works

- Large amplitude/finite temperature
- Non-conformal backgrounds
- Application to drag force



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