

**Gauge Theories**  
and  
**Non-Commutative Geometry**

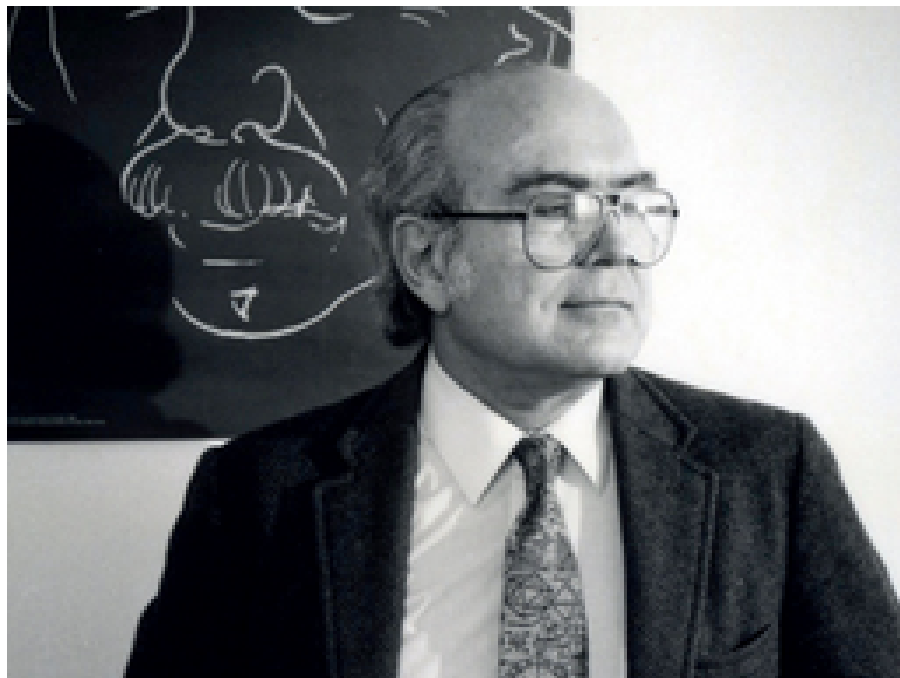
Eighth Crete Regional Meeting

in string theory

Nafplion, July 5-11 2015

John Iliopoulos

ENS Paris



# Motivation

- ▶ 1) Short distance singularities.

# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder

# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder
- ▶ 2) External fluxes.

# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder
- ▶ 2) External fluxes.
- ▶ Landau (1930) ; Peierls (1933)

# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder
- ▶ 2) External fluxes.
- ▶ Landau (1930) ; Peierls (1933)
- ▶ 3) Seiberg-Witten map.

# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder
- ▶ 2) External fluxes.
- ▶ Landau (1930) ; Peierls (1933)
- ▶ 3) Seiberg-Witten map.
- ▶ 4) Large  $N$  gauge theories and matrix models.



# Motivation

- ▶ 1) Short distance singularities.
- ▶ Heisenberg → Peierls → Pauli → Oppenheimer → Snyder
- ▶ 2) External fluxes.
- ▶ Landau (1930) ; Peierls (1933)
- ▶ 3) Seiberg-Witten map.
- ▶ 4) Large  $N$  gauge theories and matrix models.
- ▶ 5) The construction of gauge theories using the techniques of non-commutative geometry.

▶  $[x_\mu, x_\nu] = i\theta_{\mu\nu}$

simplest case:  $\theta$  is constant (canonical, or Heisenberg case).

▶  $[x_\mu, x_\nu] = i\theta_{\mu\nu}$

simplest case:  $\theta$  is constant (canonical, or Heisenberg case).

▶  $[x_\mu, x_\nu] = iF_{\mu\nu}^\rho x_\rho$  (Lie algebra case)

$$\blacktriangleright [x_\mu, x_\nu] = i\theta_{\mu\nu}$$

simplest case:  $\theta$  is constant (canonical, or Heisenberg case).

$$\blacktriangleright [x_\mu, x_\nu] = iF_{\mu\nu}^\rho x_\rho \text{ (Lie algebra case)}$$

$$\blacktriangleright x_\mu x_\nu = q^{-1} R_{\mu\nu}^{\rho\sigma} x_\rho x_\sigma \text{ (quantum space case)}$$

$$\blacktriangleright [x_\mu, x_\nu] = i\theta_{\mu\nu}$$

simplest case:  $\theta$  is constant (canonical, or Heisenberg case).

$$\blacktriangleright [x_\mu, x_\nu] = iF_{\mu\nu}^\rho x_\rho \text{ (Lie algebra case)}$$

$$\blacktriangleright x_\mu x_\nu = q^{-1} R_{\mu\nu}^{\rho\sigma} x_\rho x_\sigma \text{ (quantum space case)}$$

$\blacktriangleright$  Definition of the derivative:

$$\partial^\mu x_\nu = \delta_\nu^\mu \quad [x_\mu, f(x)] = i\theta_{\mu\nu} \partial^\nu f(x)$$

$$\blacktriangleright [x_\mu, x_\nu] = i\theta_{\mu\nu}$$

simplest case:  $\theta$  is constant (canonical, or Heisenberg case).

$$\blacktriangleright [x_\mu, x_\nu] = iF_{\mu\nu}^\rho x_\rho \text{ (Lie algebra case)}$$

$$\blacktriangleright x_\mu x_\nu = q^{-1} R_{\mu\nu}^{\rho\sigma} x_\rho x_\sigma \text{ (quantum space case)}$$

$\blacktriangleright$  Definition of the derivative:

$$\partial^\mu x_\nu = \delta_\nu^\mu \quad [x_\mu, f(x)] = i\theta_{\mu\nu} \partial^\nu f(x)$$

$\blacktriangleright$  Define a  $*$  product

$$f * g = e^{\frac{i}{2} \frac{\partial}{\partial x_\mu} \theta_{\mu\nu} \frac{\partial}{\partial y_\nu}} f(x) g(y) \Big|_{x=y}$$

All computations can be viewed as expansions in  $\theta$   
*expansions in the external field*

More efficient ways?

Quantum field theory in a space with non-commutative geometry?  
BRS Symmetry?

# Large $N$ field theories

►  $\phi^i(x)$   $i = 1, \dots, N$  ;  $N \rightarrow \infty$

$$\phi^i(x) \rightarrow \phi(\sigma, x) \quad 0 \leq \sigma \leq 2\pi$$

$$\sum_{i=1}^{\infty} \phi^i(x) \phi^i(x) \rightarrow \int_0^{2\pi} d\sigma (\phi(\sigma, x))^2$$

but

$$\phi^4 \rightarrow (f)^2$$



# Large $N$ field theories

- ▶  $\phi^i(x)$   $i = 1, \dots, N$  ;  $N \rightarrow \infty$

$$\phi^i(x) \rightarrow \phi(\sigma, x) \quad 0 \leq \sigma \leq 2\pi$$

$$\sum_{i=1}^{\infty} \phi^i(x)\phi^i(x) \rightarrow \int_0^{2\pi} d\sigma (\phi(\sigma, x))^2$$

but

$$\phi^4 \rightarrow (f)^2$$

- ▶ For a Yang-Mills theory, the resulting expression is local

# Gauge theories on surfaces

E.G. Floratos and J.I.

- ▶ Given an  $SU(N)$  Yang-Mills theory in a  $d$ -dimensional space

$$A_\mu(x) = A_\mu^a(x) t_a$$

# Gauge theories on surfaces

E.G. Floratos and J.I.

- ▶ Given an  $SU(N)$  Yang-Mills theory in a  $d$ -dimensional space

$$A_\mu(x) = A_\mu^a(x) t_a$$

- ▶ there exists a reformulation in  $d+2$  dimensions

$$A_\mu(x) \rightarrow \mathcal{A}_\mu(x, z_1, z_2) \quad F_{\mu\nu}(x) \rightarrow \mathcal{F}_{\mu\nu}(x, z_1, z_2)$$

with

$$[z_1, z_2] = \frac{2i}{N}$$

$$[A_\mu(x), A_\nu(x)] \rightarrow \{\mathcal{A}_\mu(x, z_1, z_2), \mathcal{A}_\nu(x, z_1, z_2)\}_{Moyal}$$

$$[A_\mu(x), \Omega(x)] \rightarrow \{\mathcal{A}_\mu(x, z_1, z_2), \Omega(x, z_1, z_2)\}_{Moyal}$$

$$\int d^4x \text{Tr} (F_{\mu\nu}(x) F^{\mu\nu}(x)) \rightarrow \int d^4x dz_1 dz_2 \mathcal{F}_{\mu\nu}(x, z_1, z_2) * \mathcal{F}^{\mu\nu}(x, z_1, z_2)$$

These expressions are defined for *all N!*

Not necessarily integer ???

# I. Large $N$

-A simple algebraic result:

At large  $N$

The  $SU(N)$  algebra  $\rightarrow$  The algebra of the area preserving diffeomorphisms of a closed surface. (sphere or torus).

-The structure constants of  $[SDiff(S^2)]$  are the limits for large  $N$  of those of  $SU(N)$ .

-Alternatively: For the sphere

$$x_1 = \cos\phi \sin\theta, \quad x_2 = \sin\phi \sin\theta, \quad x_3 = \cos\theta$$

$$Y_{l,m}(\theta, \phi) = \sum_{\substack{i_k=1,2,3 \\ k=1,\dots,l}} \alpha_{i_1\dots i_l}^{(m)} x_{i_1}\dots x_{i_l}$$

where  $\alpha_{i_1\dots i_l}^{(m)}$  is a symmetric and traceless tensor.

For fixed  $l$  there are  $2l + 1$  linearly independent tensors  $\alpha_{i_1\dots i_l}^{(m)}$ ,  
 $m = -l, \dots, l$ .

Choose, inside  $SU(N)$ , an  $SU(2)$  subgroup.

$$[S_i, S_j] = i\epsilon_{ijk} S_k$$

A basis for  $SU(N)$ :

$$S_{l,m}^{(N)} = \sum_{\substack{i_k=1,2,3 \\ k=1,\dots,l}} \alpha_{i_1\dots i_l}^{(m)} S_{i_1} \dots S_{i_l}$$

$$[S_{l,m}^{(N)}, S_{l',m'}^{(N)}] = if_{l,m;l',m'}^{(N)} S_{l'',m''}^{(N)}$$



The three  $SU(2)$  generators  $S_i$ , rescaled by a factor proportional to  $1/N$ , will have well-defined limits as  $N$  goes to infinity.

$$\begin{aligned}S_i &\rightarrow T_i = \frac{2}{N} S_i \\ [T_i, T_j] &= \frac{2i}{N} \epsilon_{ijk} T_k \\ T^2 &= T_1^2 + T_2^2 + T_3^2 = 1 - \frac{1}{N^2}\end{aligned}$$

In other words: under the norm  $\|x\|^2 = \text{Tr}x^2$ , the limits as  $N$  goes to infinity of the generators  $T_i$  are three objects  $x_i$  which commute and are constrained by

$$x_1^2 + x_2^2 + x_3^2 = 1$$

$$\frac{N}{2i} [f, g] \rightarrow \epsilon_{ijk} x_i \frac{\partial f}{\partial x_j} \frac{\partial g}{\partial x_k}$$

$$\frac{N}{2i} [T_{l,m}^{(N)}, T_{l',m'}^{(N)}] \rightarrow \{Y_{l,m}, Y_{l',m'}\}$$

$$N[A_\mu, A_\nu] \rightarrow \{A_\mu(x, \theta, \phi), A_\nu(x, \theta, \phi)\}$$

## II. To all orders

We can parametrise the  $T_i$ 's in terms of two operators,  $z_1$  and  $z_2$ .

$$T_+ = T_1 + iT_2 = e^{\frac{iz_1}{2}} (1 - z_2^2)^{\frac{1}{2}} e^{\frac{iz_1}{2}}$$

$$T_- = T_1 - iT_2 = e^{-\frac{iz_1}{2}} (1 - z_2^2)^{\frac{1}{2}} e^{-\frac{iz_1}{2}}$$

$$T_3 = z_2$$

If we assume that  $z_1$  and  $z_2$  satisfy:

$$[z_1, z_2] = \frac{2i}{N}$$

The  $T_i$ 's satisfy the  $SU(2)$  algebra.

If we assume that the  $T_i$ 's satisfy the  $SU(2)$  algebra, the  $z_i$ 's satisfy the Heisenberg algebra

# The techniques of non-com. geometry

- ▶ Gauge transformations are:

# The techniques of non-com. geometry

- ▶ Gauge transformations are:
- ▶ Diffeomorphisms *space-time*

# The techniques of non-com. geometry

- ▶ Gauge transformations are:
- ▶ Diffeomorphisms *space-time*
- ▶ Internal symmetries

# The techniques of non-com. geometry

- ▶ Gauge transformations are:
- ▶ Diffeomorphisms *space-time*
- ▶ Internal symmetries
- ▶ Question: Is there a space on which Internal symmetry transformations act as Diffeomorphisms?



# The techniques of non-com. geometry

- ▶ Gauge transformations are:
- ▶ Diffeomorphisms *space-time*
- ▶ Internal symmetries
- ▶ Question: Is there a space on which Internal symmetry transformations act as Diffeomorphisms?
- ▶ Answer: Yes, but it is a space with non-commutative geometry.  
A space defined by an algebra of matrix-valued functions

▶ SO WHAT?

▶ **SO WHAT?**

- ▶ A possible way to unify gauge theories and Gravity???

▶ **SO WHAT?**

- ▶ A possible way to unify gauge theories and Gravity???
- ▶ A possible connection between gauge fields and scalar fields.

▶ **SO WHAT?**

- ▶ A possible way to unify gauge theories and Gravity???
- ▶ A possible connection between gauge fields and scalar fields.
- ▶ The actual implementation brings us back to flat space calculations.

▶ **SO WHAT?**

- ▶ A possible way to unify gauge theories and Gravity???
- ▶ A possible connection between gauge fields and scalar fields.
- ▶ The actual implementation brings us back to flat space calculations.
- ▶ *New predictions for the B.E.H. mass?*

# Is the S.M. reducible?

- ▶ Can we impose a condition of the form

$$\frac{m_\phi}{m_Z} \quad \text{or} \quad \frac{m_\phi}{m_W} = C \quad ?$$

# Is the S.M. reducible?

- ▶ Can we impose a condition of the form

$$\frac{m_\phi}{m_Z} \quad \text{or} \quad \frac{m_\phi}{m_W} = C \quad ?$$

- ▶ Answer: NO! There is no fixed point in the renormalisation group equations.



# Is the S.M. reducible?

- ▶ Can we impose a condition of the form

$$\frac{m_\phi}{m_Z} \quad \text{or} \quad \frac{m_\phi}{m_W} = C \quad ?$$

- ▶ Answer: NO! There is no fixed point in the renormalisation group equations.
- ▶ Related question: Is there a B.R.S. symmetry for this model?

# Conclusions

- ▶ Non-Commutative Geometry has come to stay!

# Conclusions

- ▶ Non-Commutative Geometry has come to stay!
- ▶ Whether it will turn out to be convenient for us to use is still questionable.

# Conclusions

- ▶ Non-Commutative Geometry has come to stay!
- ▶ Whether it will turn out to be convenient for us to use is still questionable.
- ▶ It will depend on our ability to simplify the mathematics sufficiently, or to master them deeply, in order to get new insights

# Conclusions

- ▶ Non-Commutative Geometry has come to stay!
- ▶ Whether it will turn out to be convenient for us to use is still questionable.
- ▶ It will depend on our ability to simplify the mathematics sufficiently, or to master them deeply, in order to get new insights
- ▶ We need somebody with knowledge and imagination