

Multi-scale Dynamics on the Baryonic Branch of Klebanov-Strassler

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arXiv:1104.3963 [hep-th]

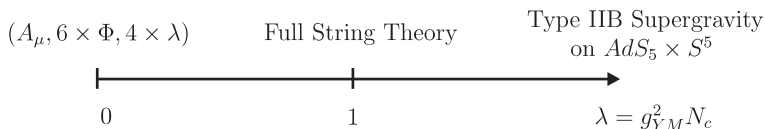
(with Jérôme Gaillard, Carlos Núñez, and Maurizio Piai)

arXiv:0908.2808 [hep-th]

(with Carlos Núñez and Maurizio Piai)

Holography:

- The most familiar and simple example is AdS/CFT where $\mathcal{N} = 4$ SYM is dual to Type IIB String Theory on $AdS_5 \times S^5$:



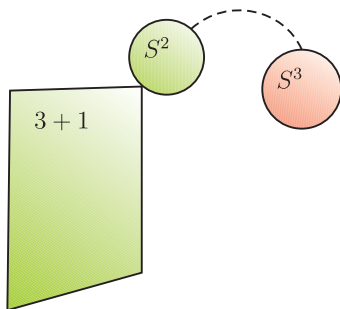
- Many generalizations, for example based on the conifold and its variations, and deformations of $\mathcal{N} = 4$ SYM
- It is interesting to look at backgrounds with less supersymmetry, and for which conformal invariance is broken

Outline of talk:

- D5 System
- Walking Dynamics
- Rotated Solutions and Klebanov-Strassler
- Summary and Outlook

D5 System

Set-up:



- D5-branes wrapped on S^2
- This gives us an $\mathcal{N} = 1$ SUSY field theory

- Type IIB supergravity ansatz (ds^2, F_3, ϕ) :

$$\begin{aligned}
 ds^2 = & e^{\frac{\phi}{2}} \left[dx_{1,3}^2 + e^{2k} d\rho^2 + e^{2h} (d\theta^2 + \sin^2 \theta d\varphi^2) + \right. \\
 & \left. \frac{e^{2g}}{4} \left((\tilde{\omega}_1 + ad\theta)^2 + (\tilde{\omega}_2 - a \sin \theta d\varphi)^2 \right) + \frac{e^{2k}}{4} (\tilde{\omega}_3 + \cos \theta d\varphi)^2 \right], \\
 F_3 = & \frac{N_c}{4} \left[-(\tilde{\omega}_1 + bd\theta) \wedge (\tilde{\omega}_2 - b \sin \theta d\varphi) \wedge (\tilde{\omega}_3 + \cos \theta d\varphi) + \right. \\
 & \left. d\rho \wedge (b'(-d\theta \wedge \tilde{\omega}_1 + \sin \theta d\varphi \wedge \tilde{\omega}_2)) + (1 - b^2) \sin \theta d\theta \wedge d\varphi \wedge \tilde{\omega}_3 \right],
 \end{aligned}$$

with

$$\begin{aligned}
 \tilde{\omega}_1 &= \cos \psi d\tilde{\theta} + \sin \psi \sin \tilde{\theta} d\tilde{\varphi}, \\
 \tilde{\omega}_2 &= -\sin \psi d\tilde{\theta} + \cos \psi \sin \tilde{\theta} d\tilde{\varphi}, \\
 \tilde{\omega}_3 &= d\psi + \cos \tilde{\theta} d\tilde{\varphi}.
 \end{aligned}$$

Finding solutions:

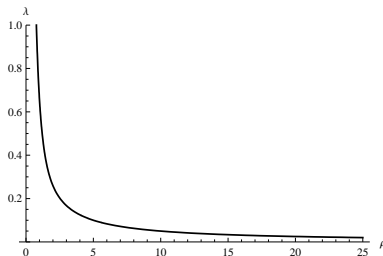
- Write down BPS equations for the background fields $\{g, k, h, \phi, a, b\}$
- These can be repackaged into a single second order differential equation (Hoyos, Nunez, Papadimitriou 2008):

$$P'' + P' \left[\frac{P' + Q'}{P - Q} + \frac{P' - Q'}{P + Q} - 4 \coth(2\rho) \right] = 0,$$
$$Q(\rho) = N_c(2\rho \coth(2\rho) - 1)$$

- Map from P to solutions in Type IIB supergravity:
 $P \rightarrow \{g, k, h, \phi, a, b\}$

Example: Maldacena-Nunez ($P = 2N_c\rho$)

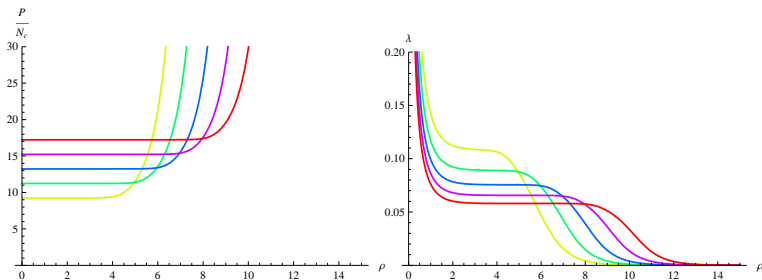
- Non-singular: in the IR, S^2 shrinks to zero size, while the size of S^3 stays finite (like the deformed conifold)
- 4d gauge coupling constant $\lambda = \frac{g^2 N_c}{8\pi^2} = \frac{N_c \coth \rho}{P}$



- One scale: set by gaugino condensate

Walking Dynamics

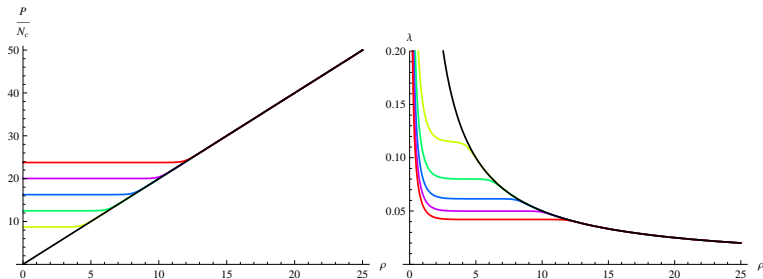
Walking backgrounds (UV behaviour: $P \sim e^{4\rho/3}$):
(Nunez, Papadimitriou, Piai 2008)



Two scales: gaugino condensate and end of walking region ρ_*

Walking Dynamics

Walking backgrounds (UV behaviour: Maldacena-Nunez):
(DE, Nunez, Piai 2009)



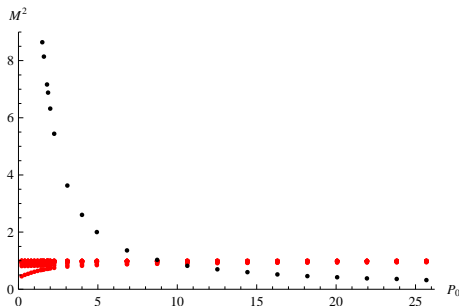
Two scales: gaugino condensate and end of walking region ρ_*

Why is the walking dynamics interesting?

- These models have similarities with Walking Technicolor
- Because of the strongly coupled dynamics, holography is a natural framework
- Question: do we have spontaneous breaking of approximate scale invariance implying a light scalar (the dilaton, pseudo-Goldstone boson of dilatations)?

Walking Dynamics

Spectrum of scalar glueballs for different values of $P_0 \simeq 2\rho_*$ (in units of $g_s\alpha'N_c$):



Light scalar whose mass is suppressed by the length of the walking region

UV behaviour of the walking backgrounds:

- Not asymptotically AdS (either MN or dim-8 operator)
- The dictionary is less well-defined
- It is not easy to identify the QFT that is dual to a particular geometry

Solution generating technique (Maldacena, Martelli 2009):

- Start with a solution to the D5 system (ds^2, F_3, ϕ)
- Generate a new (rotated) solution $(ds^{(r)2}, F_3^{(r)}, H_3^{(r)}, F_5^{(r)}, \phi^{(r)})$:

$$ds^{(r)2} = e^{\phi/2} \left[\left(1 - \kappa^2 e^{2\phi}\right)^{-1/2} dx_{1,3}^2 + \left(1 - \kappa^2 e^{2\phi}\right)^{1/2} ds_6^2 \right],$$

$$\phi^{(r)} = \phi,$$

$$F_3^{(r)} = F_3,$$

$$H_3^{(r)} = -\kappa e^{2\phi} *_6 F_3,$$

$$F_5^{(r)} = -\kappa(1 + *_{10})\text{vol}_{(4)} \wedge d\left(e^{-2\phi} - \kappa^2\right)^{-1}$$

- D5 system \rightarrow D5/D3 system
- Preserves SUSY

Rotated Solutions and Klebanov-Strassler

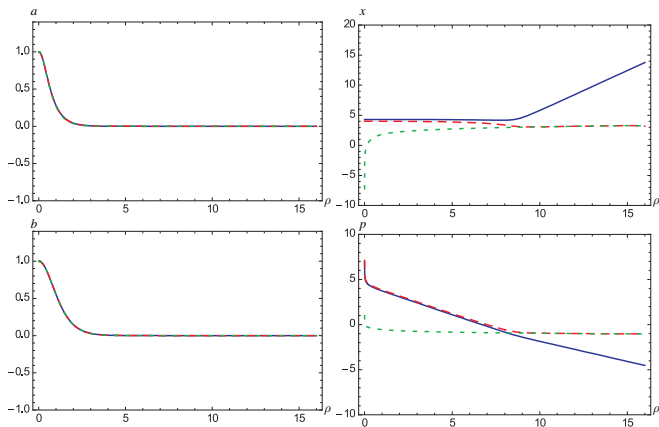
- Apply to the walking backgrounds with $P \sim e^{4\rho/3}$ in the UV (dim-8 operator)
- For $\kappa = e^{-\phi_{UV}}$, the rotated backgrounds behave asymptotically like Klebanov-Strassler in the UV
- The field theory dual to KS is more well-understood:
 $SU(N+M) \times SU(N)$ gauge group, bifundamental matter A_i and B_i ($i = 1, 2$) in representations $(N+M, \bar{N})$ and $(\bar{N+M}, N)$, superpotential $\mathcal{W} = \lambda_1 \text{Tr}(A_i B_j A_k B_l) \epsilon^{ik} \epsilon^{jl}$

Properties of the rotated solutions ($\kappa = e^{-\phi_{UV}}$):

- The dim-8 operator is no longer present, making the UV well-defined
- There is a dim-3 VEV, the gaugino condensate
- There is a dim-2 VEV, $\langle \text{Tr}(A\bar{A} - B\bar{B}) \rangle \neq 0$, signalling that we are on the baryonic branch of Klebanov-Strassler
- There is a dim-6 VEV, $\langle \text{Tr} W^2 \bar{W}^2 \rangle \neq 0$
- Compute Wilson loops \Rightarrow confinement

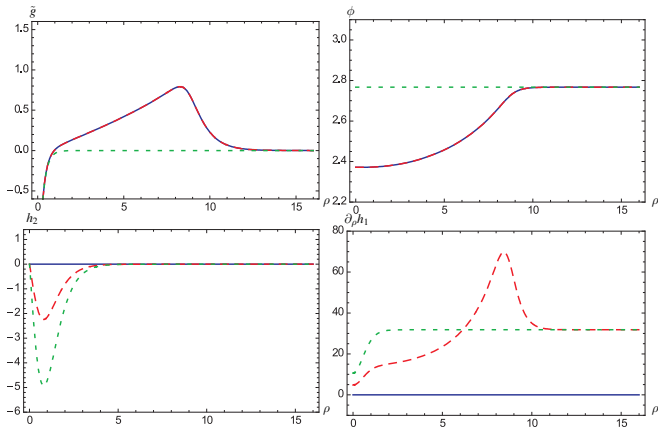
Rotated Solutions and Klebanov-Strassler

Background functions (Unrotated, Rotated, KS):



Rotated Solutions and Klebanov-Strassler

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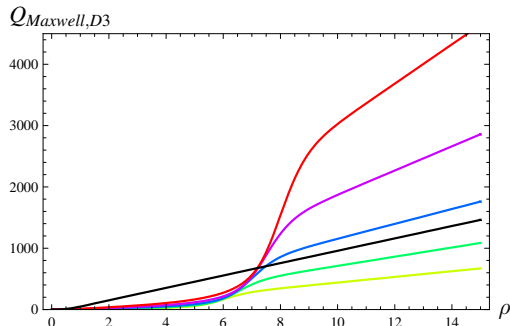


Rotated Solutions and Klebanov-Strassler

- The scale of both the dim-2 and dim-6 VEVs is at ρ_*
- Below this scale, the rotated and unrotated solutions are almost the same
- This suggests that the unrotated solution is an effective description of the rotated solution, valid at scales below ρ_*

Rotated Solutions and Klebanov-Strassler

Maxwell charge ($Q_{Maxwell,D3} \sim \int_{\Sigma_5} F_5$):



Sudden drop in the number of degrees of freedom at scale ρ_*

Picture (suggested from gravity solution):

- Flowing from the UV, the theory undergoes a duality cascade $SU(N + M) \times SU(N) \rightarrow SU(N) \times SU(N - M) \rightarrow \dots$
- This proceeds until the scale ρ_* where, being on the baryonic branch, the gauge group is Higgsed:
 $SU(\tilde{N} + \tilde{M}) \times SU(\tilde{N}) \rightarrow SU(\tilde{N})$
- Consistent with 1. the sudden drop in the number of degrees of freedom at ρ_* , and 2. that below ρ_* the theory is described by the D5 system which has a single gauge group
- Integrating out heavy modes generates the dim-8 operator of the (effective) unrotated theory

- We have described a system of wrapped D5-branes with solutions that have walking dynamics
- Applying a solution generating technique to these, we have found new Type IIB SUGRA solutions whose UV is similar to Klebanov-Strassler
- We have proposed the interpretation that these ‘rotated’ solutions are the UV completion of the ‘unrotated’ ones, which should be thought of as effective descriptions valid below the scale ρ_*

- Compute spectrum
- There is (mild) singularity in the IR — can it be resolved?
- More general backgrounds in which the scales of the dim-2 and dim-6 VEVs are not the same