### Brane Inflation and Moduli Stabilization

Guy Gur-Ari

Weizmann Institute

In collaboration with Ofer Aharony

Crete Regional Meeting, June 2011

◆□> ◆□> ◆目> ◆目> ◆日 ◆ ○ ◆

# Inflation

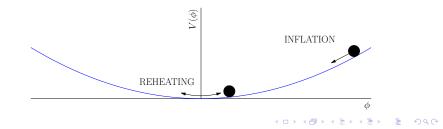
- Motivation: Solve fine-tuning problems in cosmology: Flatness, Isotropy, Baryogenesis, ...
- ► Assume early period of exponential expansion  $a(t) = e^{Ht}$
- To avoid fine-tuning, need  $a(t_{\rm f})/a(t_{\rm i})\gtrsim e^{60}$

The basic setup:

Scalar field with a shallow potential

$$V(\phi) = \lambda \, \phi^{\alpha} \,, \qquad \alpha \lesssim 1$$

Starts at a large value, rolls down slowly



# Large-Field Inflation

To get  $a_f/a_i = e^{60}$ , need initial field value  $\phi_0 > M_{\text{planck}}$ 

- Gravitational waves imprinted in CMB spectrum
- Measurable in next few years

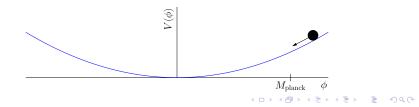
Fine-tuning problems:

Planck-suppressed operators become important, ruin shallow potential:

$$\mathcal{L}_{\rm eff} \supset \frac{\mathcal{O}_4}{M_{\rm pl}^2} \phi^2 + \cdots$$

Large field inflation not natural in effective field theory

Turn to UV complete theory



# Large-Field Inflation

To get  $a_f/a_i = e^{60}$ , need initial field value  $\phi_0 > M_{\text{planck}}$ 

- Gravitational waves imprinted in CMB spectrum
- Measurable in next few years

Fine-tuning problems:

Planck-suppressed operators become important, ruin shallow potential:

$$\mathcal{L}_{\mathrm{eff}} \supset \frac{\mathcal{O}_4}{M_{\mathsf{pl}}^2} \phi^2 + \cdots$$

- Large field inflation not natural in effective field theory
  - Turn to UV complete theory

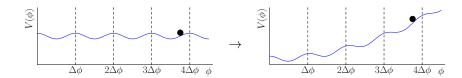


## Large-Field Inflation using Monodromy

Still need control over corrections.

Basic idea:

- Take a periodic field,  $V(\phi) = V(\phi + \Delta \phi)$
- Same quantum corrections in each period, even for large  $\phi$
- Break the periodicity ("monodromy")
- Original symmetry helps control corrections



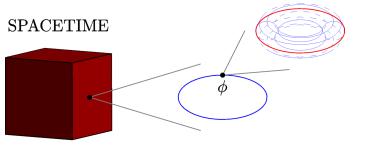
◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Use natural string theory objects:

- Internal manifold: Start with  $T^2 \times S^1 \times X_3$
- Add a D4-brane:
  - Fills spacetime
  - Wraps a 1-cycle on  $T^2$ , localized on base  $S^1$

At low energy,

- ▶ Position of brane on base cycle  $\rightarrow$  periodic scalar field  $\phi(x,t)$
- If brane stretches  $\rightarrow$  potential  $V(\phi)$

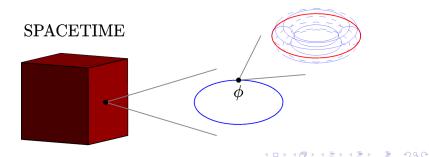


Use natural string theory objects:

- Internal manifold: Start with  $T^2 \times S^1 \times X_3$
- Add a D4-brane:
  - Fills spacetime
  - Wraps a 1-cycle on  $T^2$ , localized on base  $S^1$

At low energy,

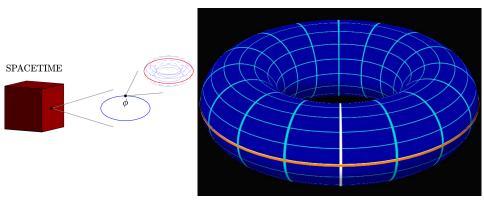
- ▶ Position of brane on base cycle  $\rightarrow$  periodic scalar field  $\phi(x,t)$
- If brane stretches  $\rightarrow$  potential  $V(\phi)$



Introduce monodromy:

- $\blacktriangleright$  Twist the  $T^2\times S^1{:}\ \tau\to \tau+1$  when going around the base circle
- Brane stretches as it moves in  $\phi$ ,

 $V(\phi) \propto \phi^{2/3}$ 

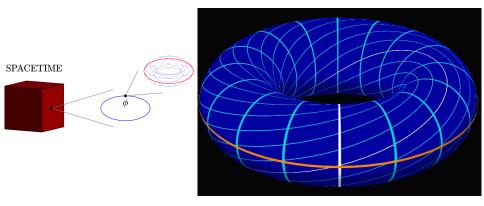


< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

Introduce monodromy:

- $\blacktriangleright$  Twist the  $T^2\times S^1{:}\ \tau\to \tau+1$  when going around the base circle
- Brane stretches as it moves in  $\phi$ ,

 $V(\phi) \propto \phi^{2/3}$ 

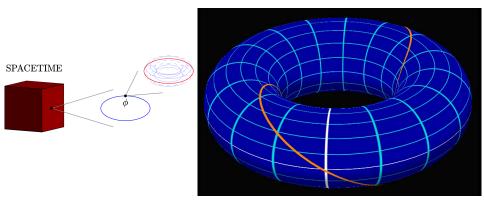


▲□▶▲□▶▲□▶▲□▶ ▲□ ● ● ●

Introduce monodromy:

- $\blacktriangleright$  Twist the  $T^2\times S^1{:}\ \tau\to \tau+1$  when going around the base circle
- Brane stretches as it moves in  $\phi$ ,

 $V(\phi) \propto \phi^{2/3}$ 



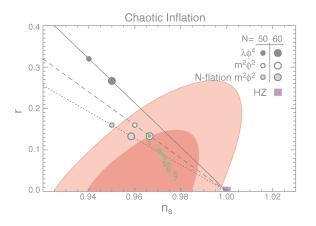
▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで

#### Generalization to Twisted 6-Torus

• We generalized from  $T^2 \times S^1 \rightarrow T^5 \times S^1$ , found more potentials:

 $V(\phi) \propto \phi^{2/3} \longrightarrow \phi^{2/3}, \phi^1, \phi^{6/5}, \phi^{4/3}, \phi^{10/7}, \phi^{3/2}, \phi^2$ 

Exponent is related to cosmological observables n<sub>s</sub>, r



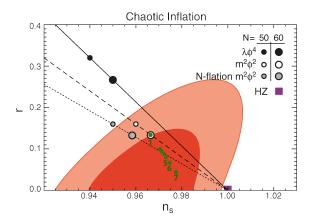
(ロ) (同) (三) (三) (三) (○) (○)

#### Generalization to Twisted 6-Torus

• We generalized from  $T^2 \times S^1 \rightarrow T^5 \times S^1$ , found more potentials:

 $V(\phi) \propto \phi^{2/3} \longrightarrow \phi^{2/3}, \phi^1, \phi^{6/5}, \phi^{4/3}, \phi^{10/7}, \phi^{3/2}, \phi^2$ 

Exponent is related to cosmological observables n<sub>s</sub>, r



(ロ) (同) (三) (三) (三) (○) (○)

### Realistic Construction [Silverstein, 2008]

Strategy for  $\mathcal{M}_3 = T^2 \times S^1$ :

- Manifold: Two copies of twisted torus,  $\mathcal{M}_3 \times \tilde{\mathcal{M}}_3$
- Assume SUGRA limit and smearing approximation
- Arrange de Sitter vacuum with small cosmological constant

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

• Add D4 brane as probe  $\rightarrow$  inflation

Requirements:

- Valid SUGRA limit
- Small backreaction of D4 brane
- No light scalars

### Realistic Construction [Silverstein, 2008]

Strategy for  $\mathcal{M}_3 = T^2 \times S^1$ :

- Manifold: Two copies of twisted torus,  $\mathcal{M}_3 \times \tilde{\mathcal{M}}_3$
- Assume SUGRA limit and smearing approximation
- Arrange de Sitter vacuum with small cosmological constant

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

• Add D4 brane as probe  $\rightarrow$  inflation

Requirements:

- Valid SUGRA limit
- Small backreaction of D4 brane
- No light scalars

The moduli:

- Dilaton  $g \sim e^{\phi}$
- Metric G<sub>ij</sub> (21 modes)
- Axions  $B_{ij}, C_i, \ldots$

Stabilization mechanisms:

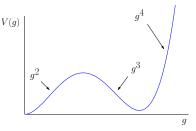
 $\blacktriangleright$  Negative curvature  $\rightarrow$  need negative energy source

(日) (日) (日) (日) (日) (日) (日)

- Add O6 plane,  $\mathcal{M}_3 \leftrightarrow \tilde{\mathcal{M}}_3$ 
  - Orientifold removes 9 metric moduli
- Twist reduces isometry:  $U(1)^3 \rightarrow U(1)$ 
  - KK gauge bosons eat 2 metric moduli

The moduli:

- Dilaton  $g \sim e^{\phi}$
- ▶ Metric G<sub>ij</sub> (21 modes)
- Axions  $B_{ij}, C_i, \ldots$



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

Stabilization mechanisms:

- ► Negative curvature → need negative energy source
- Add O6 plane,  $\mathcal{M}_3 \leftrightarrow \tilde{\mathcal{M}}_3$ 
  - Orientifold removes 9 metric moduli
- Twist reduces isometry:  $U(1)^3 \rightarrow U(1)$ 
  - KK gauge bosons eat 2 metric moduli

Remaining moduli:

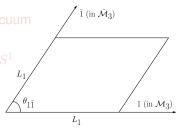
#### $g, \quad L_1, L_2, L_3, \quad \theta_{1\tilde{1}}, \theta_{2\tilde{2}}, \theta_{3\tilde{3}}, \quad G_{12}, G_{1\tilde{2}}, G_{1\tilde{3}}, G_{2\tilde{3}}$

Add ingredients that stabilize diagonal modes L<sub>i</sub>

- Fluxes, KK monopoles, discrete Wilson lines
- ► Tune parameters to form local minimum for dilaton

What about the angular modes?

- We found that  $\theta_{1\tilde{1}}, \theta_{2\tilde{2}}$  destabilize the vacuum
- No obvious way to stabilize
- $\blacktriangleright$  Analogous problems in all twisted  $T^5 \times S^1$



э

・ロット (雪) (日) (日)

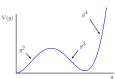
Remaining moduli:

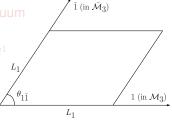
 $g, \quad L_1, L_2, L_3, \quad \theta_{1\tilde{1}}, \theta_{2\tilde{2}}, \theta_{3\tilde{3}}, \quad G_{12}, G_{1\tilde{2}}, G_{1\tilde{3}}, G_{2\tilde{3}}$ 

- Add ingredients that stabilize diagonal modes L<sub>i</sub>
  Fluxes, KK monopoles, discrete Wilson lines
- Tune parameters to form local minimum for dilaton

#### What about the angular modes?

- We found that  $\theta_{1\tilde{1}}, \theta_{2\tilde{2}}$  destabilize the vacuum
- No obvious way to stabilize
- Analogous problems in all twisted  $T^5 \times S^1$





・ロット (雪) (日) (日)

Remaining moduli:

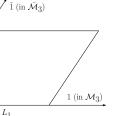
 $g, \quad L_1, L_2, L_3, \quad \theta_{1\tilde{1}}, \theta_{2\tilde{2}}, \theta_{3\tilde{3}}, \quad G_{12}, G_{1\tilde{2}}, G_{1\tilde{3}}, G_{2\tilde{3}}$ 

- ► Add ingredients that stabilize diagonal modes *L<sub>i</sub>* 
  - Fluxes, KK monopoles, discrete Wilson lines
- Tune parameters to form local minimum for dilaton

What about the angular modes?

- We found that  $\theta_{1\tilde{1}}, \theta_{2\tilde{2}}$  destabilize the vacuum
- No obvious way to stabilize
- Analogous problems in all twisted  $T^5 \times S^1$





 $\theta_{1\tilde{1}}$ 

・ロット (雪) (日) (日)

- Large-field inflation is not natural in effective field theory
- Monodromy helps control corrections
- First string theory implementation uses branes on twisted tori
- ... but is unstable due to angular metric modes
- ▶ We found no stable examples, cannot rule out existence
- Conclusion: Moduli stabilization is complicated!

Thank You

(日) (日) (日) (日) (日) (日) (日)

- Large-field inflation is not natural in effective field theory
- Monodromy helps control corrections
- First string theory implementation uses branes on twisted tori
- ... but is unstable due to angular metric modes
- We found no stable examples, cannot rule out existence
- Conclusion: Moduli stabilization is complicated!

Thank You

(ロ) (同) (三) (三) (三) (○) (○)

- Large-field inflation is not natural in effective field theory
- Monodromy helps control corrections
- First string theory implementation uses branes on twisted tori
- ... but is unstable due to angular metric modes
- We found no stable examples, cannot rule out existence
- Conclusion: Moduli stabilization is complicated!

Thank You

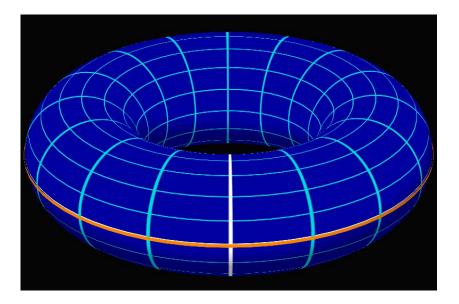
< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

- Large-field inflation is not natural in effective field theory
- Monodromy helps control corrections
- First string theory implementation uses branes on twisted tori
- ... but is unstable due to angular metric modes
- We found no stable examples, cannot rule out existence
- Conclusion: Moduli stabilization is complicated!

Thank You

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

▲□▶ ▲□▶ ▲目▶ ▲目▶ ▲□ ▶ ▲□



▲□▶▲□▶▲≡▶▲≡▶ ≡ のへで

