



Is Weinberg's Theorem violated?

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Weinberg's Theorem

hep-th/0602442

- Gravity + Inflaton + free scalars + free spinors + free vector fields
- Quantum corrections to the primordial power spectrum
- Effects down by $G H^2$
- Largest possible IR enhancement is powers of $\ln[a(t)]$

Result of Ford, Wu and Ng

gr-qc/0608002

- Gravity + Inflaton + Conformal Matter
- Extra contribution due to conf. goes like $[1/a_{\text{beg}}]^2$
- Because $a_{\text{end}} = 1$ (for them) we all GUESS effect $\sim [a(t)/a_{\text{beg}}]^2$
- TWO PUZZLES:
 1. How can Weinberg's bound be violated?
 2. How can conformal matter get a big IR enhancement?



Strength of Quantum Gravity

- Coupling constant $G = 1/M_p^2$
- Corrections go like $GE^2 = (E/M_p)^2$
- TRIVIALY large if $E > M_p$ but
 - Don't know physics at those scales
 - Perturbation theory isn't valid
- Can SOMETIMES get IR enhancement for $E \ll M_p$



Strength of Quantum Loops

- Classical response to virtual particles
- More virtuals = larger response
- TWO FACTORS give density of virtuals
 1. Time virtuals persist (controlled by Energy-Time Uncertainty Principle)
 2. Rate at which they emerge (controlled by kinetic term in Lagrangian)



How does it work for FRW?

- $ds^2 = -dt^2 + a^2(t) dx^i dx^i$
- PERSISTENCE TIME
 - $H(t) > 0$ increases it
 - Maximum effect for $q(t) < 0$ (inflation)
- EMERGENCE RATE
 - Conformal invariance reduces it



E-Time Unc. Principle for Flat Space

- $\Delta t \Delta E > 1$ to resolve ΔE
 - Hence $\Delta t \Delta E < 1$ to NOT resolve
- Virtual pair has $\Delta E = 2(m^2 + k^2)^{1/2}$
 - Hence can last $\Delta t < 1/2(m^2 + k^2)^{-1/2}$
- Eg: Vacuum polarization
 - Most for e^\pm because smallest m
 - Smallest k 's live longest \rightarrow EM stronger at shorter distances (less polarization)



E-Time Unc. Principle for FRW

- $ds^2 = -dt^2 + a^2(t) dx^i dx^i$
 - $k = 2\pi/\lambda$ still conserved
 - But physics depends on $k/a(t)$
 - $E(t) = [m^2 + k^2/a^2(t)]^{1/2}$
- $\Delta t \Delta E \rightarrow \int_t^{t+\Delta t} dt' E(t') < 1$
- $m=0 \rightarrow \int_t^{t+\Delta t} dt'/a(t')$
 - Bounded wrt Δt for inflation!
 - Any $m=0$ virtuals live forever!



What about the EMERGENCE RATE?

- Conformal invariance, a killer symmetry:

$$\mathcal{L} \xrightarrow{g'_{\mu\nu} = \Omega^2 g_{\mu\nu}} \mathcal{L}' \xrightarrow{\text{field redefinition}} \mathcal{L}$$

- EM in 4D

$$\mathcal{L} = -1/4 F_{\mu\rho} F_{\nu\sigma} g^{\mu\nu} g^{\rho\sigma} g^{1/2}$$

$$\mathcal{L}' = -1/4 F_{\mu\rho} F_{\nu\sigma} g^{\mu\nu} g^{\rho\sigma} g^{1/2} \Omega^{D-4}, \quad F'_{\mu\rho} = F_{\mu\rho}$$

- Massless conformal coupled scalar (mcc)

$$-\frac{1}{2} \partial_\mu \phi \partial_\nu \phi \sqrt{-g} - \frac{1}{4} \frac{D-2}{D-1} R \phi^2 \sqrt{-g}$$



Killer Symmetry Suppresses Emergence Rate

- $-dt^2 + a^2 dx^2 = a^2 (-d\eta^2 + dx^2) \rightarrow a d\eta = dt$
- Conformal invariance \rightarrow same locally (in conformal coordinates) as flat
 - Hence $dN/d\eta = \Gamma_{\text{flat}}$
 - Hence $dN/dt = \Gamma_{\text{flat}}/a(t)$
- Any $m=0$, conf. virtuals that emerge DO live forever, but few emerge!



Maximum IR Enhancements

- Requirements:
 - PERSISTENCE TIME \rightarrow $m=0$ + inflation
 - EMERGENCE RATE \rightarrow no conf. invariance
- Realized by:
 - Massless, minimally coupled scalars
 - Gravitons
- NOT by conformal matter
 - How did Wu, Ng and Ford get anything?

They initially used Raychaudhuri Eqn.

- $d\theta/dt = -R_{\mu\nu}u^\mu u^\nu - 1/3\theta^2 - \sigma_{\mu\nu}\sigma^{\mu\nu} + \omega_{\mu\nu}\omega^{\mu\nu} + (u^\nu{}_{;\mu}u^\mu)_{;\nu}$
- $\theta(t) = \# \int \# dt [R_{\mu\nu}u^\mu u^\nu]_q$
 - $R_{\mu\nu} = 8\pi G(T_{\mu\nu} - 1/2g_{\mu\nu}T)$
 - $(T_{\mu\nu})_{RW} = a^{-2}(T_{\mu\nu})_{flat}$
 - $d\rho/dt + \theta(\rho + p) = 0 \longrightarrow$
 $\delta\rho/\rho_0 = -(1+w) \int d\eta \theta/a(\eta)$
- Cannot ignore $(u^\nu{}_{;\mu}u^\mu)_{;\nu}$ because it cancels parts of $-R_{\mu\nu}u^\mu u^\nu$

Cannot Ignore the Acceleration Term

- Co-moving gauge : $u^\mu = \delta^\mu_0 \longrightarrow g_{00} = -1$

$$ds^2 = -dt^2 + 2ah_{0i}dx^i dt + a^2(\delta_{ij} + h_{ij})dx^i dx^j$$

- $\theta = u^\mu{}_{;\mu} = d/dt \ln(g^{1/2}) = 3H + 1/2 d/dt h_{ii}$
- $\delta\theta = 1/2 d/dt h_{ii}$ due to conformal matter

$$u^\mu = - \frac{g^{\mu\nu} \partial_\nu \varphi}{\sqrt{-g^{\alpha\beta} \partial_\alpha \varphi \partial_\beta \varphi}} .$$

$$u^0 = 1 + O(\Delta^2) , u^i = \frac{h_{0i}}{a} + \frac{\partial_i \delta\varphi}{a^2 \dot{\varphi}_0} + O(\Delta^2) .$$

$$h_{0i}(t, \vec{x}) = -\partial_i \left[\frac{\delta\varphi(t, \vec{x})}{a(t)\dot{\varphi}_0(t)} \right] + O(\Delta^2) .$$

$$R_{00} = -3\dot{H} - 3H^2 - \frac{1}{2} \left[\ddot{h}_{ii} + 2H\dot{h}_{ii} \right] + \frac{1}{a} \left[\dot{h}_{0i,i} + Hh_{0i,i} \right] + O(\Delta^2) .$$

$$(u^\nu{}_{;\mu} u^\mu)_{;\nu} \longrightarrow \frac{1}{\sqrt{-g}} \partial_\nu \left(\sqrt{-g} g^{\nu k} \dot{g}_{k0} \right) = \frac{1}{a} \left[\dot{h}_{0i,i} + Hh_{0i,i} \right] + O(\Delta^2) .$$

Really h_{ij} Mixes with $\delta\varphi$

First order perturbation :

$$\delta\ddot{\varphi} + 3H\delta\dot{\varphi} + V_0''\delta\varphi + \frac{1}{2}\dot{\varphi}_0\dot{h}_{ii} = 0$$

$$-\kappa^2\dot{\varphi}_0\delta\dot{\varphi} + \frac{1}{2}\kappa^2V_0'\delta\varphi - \frac{\nabla^2}{a^2}\frac{\partial}{\partial t}\left[\frac{\delta\varphi}{\dot{\varphi}_0}\right] - \frac{1}{2a^2}\frac{\partial}{\partial t}\left[a^2\dot{h}_{ii}\right] = \frac{1}{2}\kappa^2T_{00}^{\text{conf}}$$

$$(\delta\varphi)_1 = \dot{\varphi}_0 \quad , \quad (\dot{h}_{ii})_1 = 6\dot{H}$$

$$\delta\varphi(x) \equiv \dot{\varphi}_0(t)\Phi(x)$$

$$\dot{h}_{ii} = 6\dot{H}\Phi - \left(4\frac{\ddot{\varphi}_0}{\dot{\varphi}_0} + 6H\right)\dot{\Phi} - 2\ddot{\Phi}$$

$$\left\{ \partial_t^2 + \left(5H + 2\frac{\ddot{\varphi}_0}{\dot{\varphi}_0}\right)\partial_t + 4\dot{H} + 6H^2 + 4H\frac{\ddot{\varphi}_0}{\dot{\varphi}_0} - 2\frac{\ddot{\varphi}_0^2}{\dot{\varphi}_0^2} + 2\frac{\ddot{\varphi}_0}{\dot{\varphi}_0} - \frac{\nabla^2}{a^2} \right\} \dot{\Phi} = \frac{1}{2}\kappa^2T_{00}^{\text{conf}}$$

Green's Function Solution

- The formal solution: $d\Phi/dt \equiv \Psi$

$$\tilde{G}(t, t', k) = \frac{\theta(t-t')}{W(t', k)} \left\{ \Psi_2(t, k)\Psi_1(t', k) - \Psi_1(t, k)\Psi_2(t', k) \right\}$$

$$\dot{\Phi}(x) = \int d^D x' G(x; x') \frac{\kappa^2}{2} T_{00}^{\text{conf}}(x')$$

- It cannot be solved exactly but can be related to solution of Mukhanov equation

$$\left[\partial_t^2 + H\partial_t + \frac{k^2}{a^2} - \left(\frac{\ddot{\theta} + H\dot{\theta}}{\theta} \right) \right] Q_i = 0 \quad , \quad \theta(t) \equiv H/(a\sqrt{-\dot{H}})$$

$$\Psi_i(t, k) = \frac{Q_i(t, k)}{a^2(t)\sqrt{-\dot{H}(t)}}$$

$$Q_{\pm}(t, k) = e^{\pm ik \int^t dt' / a(t')}$$

- Before the 1st horizon crossing :
- After the 1st horizon crossing :

$$Q_s(t) = \theta(t) \quad ,$$

$$Q_\ell(t) = \theta(t) \int^t \frac{dt'}{a(t')\theta^2(t')}$$



Asymptotic Results for $d\Phi/dt$

- $d\Phi/dt$ is BIG at early times
- $d\Phi/dt$ is SMALL at late times
- $\Phi(t) = \Phi_0 + \int_0^t dt' d\Phi/dt'$
- THIS EXPLAINS THEIR BIG $\Phi(\text{end})!$
 - Not from late times when $k/a(t) \ll M_p$
 - From EARLY times when $k/a(t) \gg M_p$



Explaining the Two Puzzles

- No violation of Weinberg's Thm
 - Their $[1/a(\text{beg})]^2$ is NOT $[a(t)/a(\text{beg})]^2$
 - It's really $[k/a(\text{beg})]^2$
- No IR enhancement from conf. matter
 - Effect from early times when $k/a \gg M_p$
 - All modes give big effects in this regime!



Trans-Planckian Issue

- No paradoxes, but is it RIGHT?
- Problems with $k/a(\text{beg}) \gg M_p$
 - We don't know fundamental physics there
 - Perturbation theory isn't valid
- We agree to disagree
 - Wu, Ng & Ford: start at beginning of inflation & use pert. theory → big corrections from early times
 - Miao & Woodard: assume modes emerge from trans-Planck regime in vacuum → only small corrections from late times



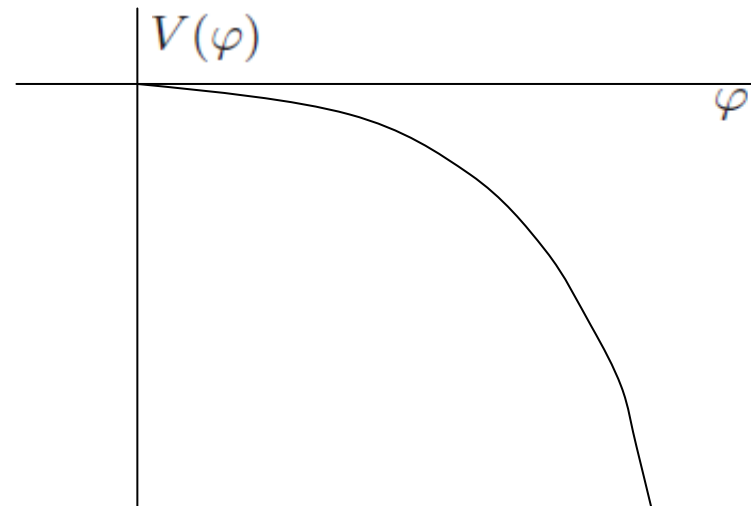
Bua's Bound

- Bua Chaicherdsukul (hep-th/0611352) :
No large effects to power spectrum from fermions or gauge particles, no matter how they couple

Counter example to Bua's bound --- Yukawa theory

$$\mathcal{L} = -\frac{1}{2}\partial_\alpha\varphi\partial_\beta\varphi g^{\alpha\beta}\sqrt{-g} - \frac{1}{2}\delta\xi\varphi^2 R\sqrt{-g} - \frac{1}{4!}\lambda\varphi^4\sqrt{-g} \\ + i\bar{\psi}e^{\beta}_b\gamma^b\mathcal{D}_\beta\psi\sqrt{-g} - f\varphi\bar{\psi}\psi\sqrt{-g}$$

- Integrate out fermions at leading Log. Order and compute V_{eff}
 - Scalar induces Fermion mass
 - Fermion vacuum $E < 0$
 - Unbounded below!



- When can we get big QFT effects?



Conclusions

- No Violation to Weinberg's theorem
- Big quantum corrections due to conformal matter come from highly Trans-Planckian regime
 - Only MMC scalars and gravitons get big IR enhancements
- Agree to disagree about significance
- QFT corrections CAN be big
 - No general rule yet