

New susy type II AdS_4 vacua

Dimitrios Tsimpis

Arnold Sommerfeld Center for Theoretical Physics
Ludwig-Maximilians-Universität

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Κολυμπάρι, Κρήτη

Outline

- 1 Introduction
 - Motivation
 - The result
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- 2 General Description of $\text{AdS}_4 \times_w \mathcal{M}_6$
 - $\text{SU}(3)$ structure
 - $\text{SU}(3) \times \text{SU}(3)$ structure
- 3 New AdS_4 vacua
 - Overview
 - Type IIA
 - Type IIB
- 4 Conclusions
 - Summary
 - Open problems

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Motivation

AdS₄/CFT₃ duality

- Nonperturbative definition of String Theory
- Multiple M2 branes

Deviation from CY-ness

- Flux $\neq 0 \implies \text{AdS}_4 \times_w \mathcal{M}_6$

Flux Vacua

- Moduli stabilization
- Starting point for de Sitter
- Susy breaking, ...

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The result

- New $\mathcal{N} = 2$ IIA and $\mathcal{N} = 1$ IIB pure-flux vacua

Ten-dimensional spacetime

$\text{AdS}_4 \times_w \mathcal{M}_6$ with:

$$ds^2(\mathcal{M}_6) = dt^2 + ds_t^2(\mathcal{M}_5)$$

where \mathcal{M}_5 admits a Sasaki-Einstein structure

$$ds_t^2(\mathcal{M}_5) = e^{2B(t)} ds_{KE}^2 + \xi^2(t)(d\psi + A)^2$$



Lüst, DT, JHEP 0904



Lüst, DT, arXiv:0906

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Outline

- $\text{AdS}_4 \times \mathcal{M}_6$ vacua: generalities
- New explicit vacua

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SU(3) structure

Flux = 0, $\mathcal{M}_6 = \text{CY}$

- SUSY: $\nabla\eta = 0$
- Bilinears: $J_{mn} := \eta^\dagger \gamma_{mn} \eta$; $\Omega_{mnp} := \eta \gamma_{mnp} \eta$
- Differential conditions: $d\Omega = 0$; $dJ = 0$

SU(3) structure

Flux $\neq 0$, $\mathcal{M}_6 \neq \text{CY}$

- SUSY: $\nabla\eta \neq 0$
- Bilinears: $J_{mn} := \eta^\dagger \gamma_{mn} \eta$; $\Omega_{mnp} := \eta \gamma_{mnp} \eta$
- Differential conditions: $d\Omega \neq 0$; $dJ \neq 0$

SU(3) structure

[hep-th]

(η, g_{mn})

SU(3) structure

[math.DG]

(J, Ω)

SU(3) structure

Definition

$$\Omega \wedge \bar{\Omega} = \frac{1}{3!} J^3$$

$$J \wedge \Omega = 0$$

$$Sp(6, \mathbb{R}) \cap SL(3, \mathbb{C}) = SU(3)$$

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Spinor ansatz

Two susy parameters

$$\epsilon_i = \zeta \otimes \theta_i + \text{c.c.}; \quad i = 1, 2$$

$$|\theta_1|^2 = |\theta_2|^2 \propto e^A$$

Two unimodular spinors

$$\theta_1 = a \eta_1; \quad \theta_2 = \begin{cases} b \eta_2^* + c^* \eta_1^* & \text{IIA} \\ b \eta_2 + c \eta_1 & \text{IIB} \end{cases}$$

$$a^2 = b^2 + |c|^2$$

$\text{SU}(3) \times \text{SU}(3)$ structure

Local $\text{SU}(2) = \text{SU}(3) \cap \text{SU}(3)$ structure

$$\tilde{J} \wedge \omega = 0$$

$$2\tilde{J} \wedge \tilde{J} = \omega \wedge \omega^* \neq 0$$

$$\iota_K \tilde{J} = \iota_K \text{Re} \omega = \iota_K \text{Im} \omega = 0$$

$\text{SU}(3) \times \text{SU}(3)$ structure

$$J^{(1)} = \frac{i}{2} K \wedge K^* + \tilde{J}; \quad J^{(2)} = \frac{i}{2} K \wedge K^* - \tilde{J}$$

$$\Omega^{(1)} = -i\omega \wedge K; \quad \Omega^{(2)} = i\omega^* \wedge K$$

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The scalar ansatz

Provided certain conditions are obeyed by (K, \tilde{J}, ω)
the scalar ansatz solves the **susy** equations of **IIA/IIB**



Lüst, DT, JHEP 0904

The scalar ansatz

Metric

$$ds^2 = e^{2A} ds^2(AdS_4) + ds^2(\mathcal{M}_6)$$

Three-form

$$H = \frac{1}{24} \left(h_1 \omega^* + h_2 \omega + 2h_3 \tilde{J} \right) \wedge K + \text{c.c.}$$

The scalar ansatz

RR fluxes IIA

$$e^\phi F_0 = f_0$$

$$e^\phi F_2 = \frac{1}{8} \left(f_2 \omega^* + f_3 \tilde{\mathbf{J}} + 2if_1 K \wedge K^* \right) + \text{c.c.}$$

$$e^\phi F_4 = \frac{1}{16} g_1 \tilde{\mathbf{J}} \wedge \tilde{\mathbf{J}} + \frac{i}{96} \left(g_2 \omega^* + g_2^* \omega + 2g_3 \tilde{\mathbf{J}} \right) \wedge K \wedge K^*$$

$$e^\phi F_6 = f \text{vol}_6$$

The scalar ansatz

RR fluxes IIB

$$e^\phi F_1 = g_1 K + \text{c.c.}$$

$$e^\phi F_3 = \frac{1}{24} \left(f_1 \omega^* + f_2 \omega + 2f_3 \tilde{J} \right) \wedge K + \text{c.c.}$$

$$e^\phi F_5 = g_2 \star_6 K + \text{c.c.}$$

Equations of motion

- SUSY \oplus (generalized) Bianchi ids \implies All EOM's
- Also in the presence of **calibrated** sources



Koerber, DT, JHEP 0708

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Construction of vacua

IIA Strict $\text{SU}(3)$

- Until recently all known (massive) vacua were of rigid- $\text{SU}(3)$ type
- All can be described in a unifying framework: left-invariant $\text{SU}(3)$ structures on groups, cosets
- Many more should be possible!



Lüst, DT, JHEP 0502



Koerber, Lüst, DT, JHEP 0807

Explicit vacua

IIA Strict SU(3)



Nilsson, Pope, CQG 1984



Sorokin, Tkach, Volkov PLB 1985



Behrndt, Cvetič, NPB 2005



Lüst, DT, JHEP 0502



Graña, Minasian, Petrini, Tomasiello JHEP 0705



Aldazabal, Font, JHEP 0802



Tomasiello, hep-th 0712



Koerber, Lüst, DT, JHEP 0807

Explicit vacua

IIA $SU(3) \times SU(3)$



Gaiotto, Tomasiello, arXiv:0904



Petrini, Zaffaroni, arXiv:0904



Lüst, DT, arXiv:0906

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New $\mathcal{N} = 2$ IIA vacua

\mathcal{M}_5 admits a S-E $\text{SU}(2)$ structure

$$\iota_u \alpha = \iota_u \beta = \iota_u \gamma = 0$$

$$\alpha \wedge \beta = \beta \wedge \gamma = \gamma \wedge \alpha = 0$$

$$\alpha \wedge \alpha = \beta \wedge \beta = \gamma \wedge \gamma \neq 0$$

$$du = -2\gamma; \quad d(\alpha + i\beta) = -3iu \wedge (\alpha + i\beta); \quad d\gamma = 0$$

5d Sasaki-Einstein

Killing spinor

$$\nabla_m \eta = \pm \frac{i}{2} \Gamma_m \eta$$

$$R_{mn} = 4g_{mn}$$

The $\text{SU}(2)$ structure

$$u_m := (\eta^\dagger \Gamma_m \eta)$$

$$\alpha_{mn} + i\beta_{mn} := (\eta \Gamma_{mn} \eta)$$

$$\gamma_{mn} := i(\eta^\dagger \Gamma_{mn} \eta)$$

New $\mathcal{N} = 2$ IIA vacua

Local $\text{SU}(2)$ structure

$$K = e^{B(t)} (dt - i\xi(t)u)$$

$$\begin{pmatrix} \tilde{J} \\ \text{Re}\omega \\ \text{Im}\omega \end{pmatrix} = e^{2C(t)} \mathcal{R}(\theta(t), \chi(t)) \begin{pmatrix} \alpha \\ \beta \\ \gamma \end{pmatrix}$$

Metric on \mathcal{M}_6

$$ds^2(\mathcal{M}_6) = e^{2B} \left(\frac{3}{W} e^{2(C-B)} ds_{KE}^2 + \xi^2 u \otimes u + dt^2 \right)$$

- \mathcal{M}_6 is locally a smooth S^2 bundle over 4D K-E base.

New $\mathcal{N} = 2$ IIA vacua

- All conditions of the scalar ansatz are satisfied
- All BI's and EOM's are satisfied
- All fields are determined

The solution

$$\tan \theta = \frac{\tan \varphi}{\sin \varepsilon} = \sqrt{2} \tan(\sqrt{2}t)$$

$$f = -3We^{-A} \cos \varepsilon \cos \varphi$$

$$f_0 = -We^{-A} (\cos \varphi \sin \varepsilon + \csc \varepsilon \sin \varphi \tan \varphi)$$

$$f_1 = -\cos \varepsilon \cos \varphi \left(We^{-A} + 4e^{-B} A' \cot \varphi \sin \varepsilon \right)$$

$$f_2 = -8e^{-B} A' \cos \varepsilon \cos \varphi$$

New $\mathcal{N} = 2$ IIA vacua

- All conditions of the scalar ansatz are satisfied
- All BI's and EOM's are satisfied
- All fields are determined

The solution

$$g_1 = -8 (\cos \varphi \sin \varepsilon + \sin \varphi \csc \varepsilon \tan \varphi) \\ \left(We^{-A} - 4e^{-B} A' \cot \varphi \sin \varepsilon \right)$$

$$g_2 = 48 \sin \varphi \left(We^{-A} + e^{-B} A' \cot \varphi \sin \varepsilon \right)$$

$$h_1 = -6 \sin^2 \varphi \cot \varepsilon \left(We^{-A} - 2e^{-B} A' \sin \varepsilon \cot \varphi \right)$$

$$\frac{h_1}{h_3} = \frac{h_2}{h_3} = \frac{f_2}{f_3} = \frac{g_2}{g_3} = -\tan \theta .$$

New $\mathcal{N} = 2$ IIA vacua

- All conditions of the scalar ansatz are satisfied
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The solution

$$e^{4A} = \frac{1}{\cos^2 \theta} \tan \varepsilon$$

$$e^{B-A} = -\frac{1}{2W} \cot \theta (\log \tan \varepsilon)'$$

$$e^{\phi-3A} = \cos \varphi \cos \varepsilon$$

$$\xi = \frac{3}{2W} e^{A-B} \sin \theta$$

New $\mathcal{N} = 2$ IIA vacua

- All conditions of the scalar ansatz are satisfied
- All BI's and EOM's are satisfied
- All fields are determined

The solution

$$\begin{aligned}\zeta' &= \frac{1}{2W} e^{2(A-C)} \cos \theta \cot \varepsilon \sin^2 \varphi (\log \tan \varepsilon)' \\ \theta' &= \cot \theta \left(\frac{1}{2W} e^{2(A-C)} \sin^2 \theta - 1 \right) (\log \tan \varepsilon)' \\ C' &= -\frac{1}{4W} e^{2(A-C)} (\sin^2 \varphi + \cos^2 \theta) (\log \tan \varepsilon)'\end{aligned}$$

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New susy IIB vacua

Local SU(2) structure

$$K = e^{A(t)} (idt - 3u)$$
$$\begin{pmatrix} \tilde{J} \\ \text{Re}\omega \\ \text{Im}\omega \end{pmatrix} = e^{2A(t)} \mathcal{R}(\theta(t)) \begin{pmatrix} \alpha \\ \beta \\ \gamma \end{pmatrix}$$

Metric on \mathcal{M}_6

$$ds_E^2 = ds^2(AdS_4) + \frac{6}{5W^2} ds_{KE}^2 + 9u \otimes u + dt^2$$

- \mathcal{M}_6 is the product of squashed 5D S-E and S^1 or \mathbb{R} .

New susy IIB vacua

- All conditions of the scalar ansatz are satisfied
- All BI's and EOM's are satisfied
- All fields are determined

The solution

$$H = \frac{1}{2} W \text{Re} w \wedge dt - \left(2A' \tilde{J} + c e^{-4A} \text{Re} w \right) \wedge u$$

$$e^\phi F_1 = -2c e^{-4A} dt$$

$$e^\phi F_3 = -\frac{1}{2} W \tilde{J} \wedge dt + \left(2A' \text{Re} w - c e^{-4A} \tilde{J} \right) \wedge u$$

$$e^\phi F_5 = \frac{3}{2} W \tilde{J} \wedge \tilde{J} \wedge u$$

New susy IIB vacua

- All conditions of the scalar ansatz are satisfied
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The solution

$$\phi = 4A$$

$$e^\phi = \begin{cases} \frac{2}{\sqrt{5}} \left| \frac{c}{W} \right| \cosh [\sqrt{5} W(t - t_0)] , & c \neq 0 \\ \exp [\sqrt{5} W(t - t_0)] , & c = 0 \end{cases}$$

$$\tan(\theta - \theta_0) = \begin{cases} \tanh \left[\frac{\sqrt{5}}{2} W(t - t_0) \right] , & c \neq 0 \\ 0 , & c = 0 \end{cases}$$

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Conclusions

Summary

- Given any 5D S-E manifold we have constructed corresponding families of new explicit IIA/IIB AdS_4 vacua
- Infinite number of families
- Pure flux
- Given any 4D K-E manifold we have constructed corresponding families of vacua

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Open problems

- Implications for $\text{AdS}_4/\text{CFT}_3$
(specific models)
- Many more vacua? $\mathcal{N}=3$?
(more CFT_3 's than AdS_4 's)
- General properties of flux vacua?
- Phenomenological applications

Thank You