# BACKREACTION OF THE HAWKING RADIATION

Three recent papers\* by G.A. VILKOVISKY

Presented by

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A semi-classical derivation in 1 + 3 dimensions for spherically symmetric, uncharged, collapsing matter of mass M >>  $\mu$ (Planck mass) having compact spatial support.

As a black hole forms and evolves, its Hawking process decreases its initial mass by about 10%, and creates a vacuum-induced "charge" equal to the remaining mass, corresponding to a stress-energy tensor of the type associated with a long-range field.

\* hep-th/0511182, hep-th/0511183, hep-th/0511184

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#### INTRODUCTION

Toy models (1 + 1 dimension) miss essential physics

# VILKOVISKY PAPERS

Kinematics of evaporating black holeshep-th/0511182Radiation equations for black holeshep-th/0511183Backreaction of the Hawking radiationhep-th/0511184

## **Definitions**:

★ Classical, Semiclassical, and Ultraviolet Regions of a macroscopic collapsing mass — based on the correspondence principle and the concept of causality.

The semiclassical region covers the black hole evolution from macroscopic to microscopic scale.

Charge: the strength of a long-range field having its source in a compact domain.

**Bondi charges**  $\mathcal{M}$  and  $Q^2$  (not necessarily an electric charge): the coefficients of a weak field metric at  $\mathcal{I}^+$ . In the (u,r) coordinates:

$$(\nabla r)^{2}\Big|_{q^{+}} = 1 - \frac{2\mathcal{M}(u)}{r} + \frac{Q^{2}(u)}{r^{2}} + \dots$$
$$u\Big|_{q^{+}} = ct - r_{*}, \quad r_{*} \coloneqq r + r_{g} \ln \left|\frac{r}{r_{g}} - 1\right|, \quad r_{g} \coloneqq \frac{2GM}{c^{2}}$$

**Conclusion**: The collapse metric in the semiclassical region can be written in terms of the 2 arbitrary Bondi charges.

#### **RADIATION EQUATIONS for black holes**

**Bondi charges equations:** 

 $\frac{d\mathcal{M}}{du} = -\frac{\mu^2}{48\pi}\mathcal{K}^2(1+\Gamma)$  $\frac{dQ^2}{du} = \frac{\mu^2}{24\pi}\mathcal{K}$ 

μ Planck mass, M initial mass,  $\frac{\mu}{M} \ll 1$  $\mathcal{K} \coloneqq \frac{(\mathcal{M}^2 - Q^2)^{1/2}}{r_{AH}^2}$  for u > u<sub>0</sub>

Apparent horizon  $r_{AH}$  calculated to satisfy  $1 - \frac{2\mathcal{M}(u)}{r_{AH}} + \frac{Q^2}{r_{AH}^2} = 0$ 

 $\Gamma$  function of  $\mathcal{M}$  and Q; uniformly bounded, small:  $\Gamma \leq 27/160$ 

## **Derivation:**

The collapse metric satisfies the WKB approximation of the solution of expectation-value equation

$$R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R = 8\pi T_{source}^{\mu\nu} + 8\pi \left\langle T_{vac}^{\mu\nu} \right\rangle$$

The vacuum is modeled as a massless scalar field  $\Phi$  in the collapse metric. Its stress-energy tensor is computed at  $\mathcal{I}^+$ :

$$\left\langle T_{vac}^{\mu\nu}\right\rangle\Big|_{q^{+}} = \left\langle invac\Big|T^{\mu\nu}\Big|invac\right\rangle\Big|_{q^{+}} - \left\langle outvac\Big|T^{\mu\nu}\Big|outvac\right\rangle\Big|_{q^{+}}$$

 $\Phi$  is expanded in spherical harmonics.

 $\Phi$  in 4 dimensions is a sequence of  $\varphi_{\ell}$  in 2 dimensions.

2-d models are not 4-d models truncated at  $\ell = 0$ .

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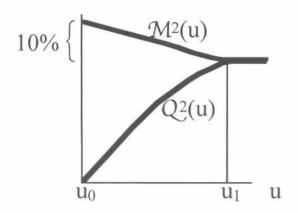
#### **BACKREACTION of the HAWKING RADIATION**

The Bondi charges equations are solved with initial conditions at  $u_0$  when the apparent horizon is tangent to an outgoing ray.

 $\mathcal{M}(u_0) = M$ ,  $Q(u_0) = \mathcal{O}$  (vanishes as  $\mu/M \to 0$ ) Let  $u_1$  be the value of u such that

$$\mathcal{M}^2(\mathbf{u}_1) = Q^2(\mathbf{u}_1)$$

At  $u_1$  the apparent horizon is tangent to an incoming light ray.



Beyond  $u_1$  the calculation is not valid; the WKB approximation remains valid.

Approximate numerical values.

$$u_1 - u_0 = 96\pi \frac{M^3}{\mu^2}$$
,  $0.098 < \frac{M - M(u_1)}{M} < 0.112$ 

The mass decreases by 10%: The Hawking process liberates more than half of the energy of the black hole.

A small part (10%) goos away as thermal radiation