



GEOMETRY FROM MODULAR CHAOS

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THE LIFE OF QUANTUM INFORMATION

1

storage

Quantum systems store information in their entanglement pattern, e.g. quantum error correcting codes.

Entanglement assists the efficient transport of information through quantum teleportation.

propagation

2

3

processing

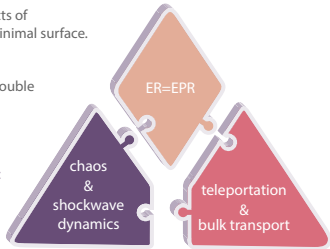
Unitarity and locality impose bounds on the allowed speed of information processing, e.g. maximal chaos bound.

THE GRAVITY- INFORMATION TRIANGLE

1) Features of **entanglement** are mapped to aspects of **geometry**, e.g. Entanglement entropy = Area of minimal surface.

2) The **teleportation** protocol in the thermofield double state of two CFTs holographically **propagates** the excitation through a **bulk traversable wormhole**.

3) The **scrambling** of CFT operators due to **chaotic dynamics** describes the exponential **blueshifting of the infalling bulk particle** due to the black hole's gravitational well.



THE PLAN

Chaos in the
entanglement
pattern

I will present a
chaos bound
characterizing the
modular flow of a
QFT subregion.

Maximal
modular chaos
in
holography

Saturation of
modular chaos bound will
allow us to construct bulk
 $T_{\mu\nu}$ **lightray operators**
from CFT data.

Bulk locality
&
curvature

Bulk **curvature**
and **local Poincare**
algebra is extracted
from
chaos commutators.

MODULAR HAMILTONIAN ESSENTIALS I

Consider a state $|\psi\rangle$ and the algebra $\mathcal{A}(R)$ of observables localized in spacetime region R .

The **modular Hamiltonian** H_{mod} is a Hermitian operator generating an automorphism of $\mathcal{A}(R)$.

It encodes the **pattern of entanglement** between degrees of freedom in R and its complement.

In **lattice systems** it is defined as:

$$H_{\text{mod}} = -\log \rho(R) + \log \rho(R^c)$$

H_{mod} is an unbounded operator. However:

$$\|e^{iH_{\text{mod}}s}O|\psi\rangle\| < \infty$$

for: $O \in \mathcal{A}(R)$

$$-\frac{1}{2} \leq \text{Im}[s] \leq 0$$

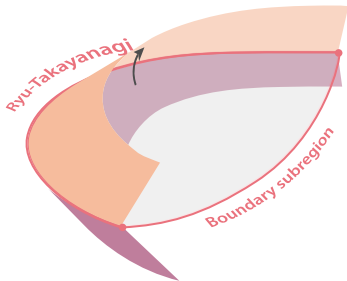
For two subregions $R_1 \subset R_2$ one can prove that:

$$1) \quad H_{\text{mod}}(R_1) \leq H_{\text{mod}}(R_2)$$

$$2) \quad \|e^{-iH_{\text{mod}}(R_2)s}e^{iH_{\text{mod}}(R_1)s}\| \leq 1$$

in the complex strip: $-\frac{1}{2} \leq \text{Im}[s] \leq 0$

MODULAR HAMILTONIAN ESSENTIALS II



Jafferis-Lewkowycz-Maldacena-Suh:

In holographic CFTs, the modular Hamiltonian of a subregion is **equal** to the modular Hamiltonian for the bulk degrees of freedom in the **entanglement wedge**:

$$H_{\text{mod}}^{\text{bulk}} = H_{\text{mod}}^{\text{CFT}}$$

In a **small neighborhood of the RT surface**, generates **homogeneous boosts** along the orthogonal 2-d plane.

$$H_{\text{mod}}^{\text{bulk}} \rightarrow \epsilon^{\alpha\beta} x_{\beta} \partial_{\alpha} + \mathcal{O}(x^2)$$

A BOUND ON MODULAR CHAOS

Suppose we **infinitesimally perturb** the modular Hamiltonian of a subalgebra $H_{\text{mod}} \rightarrow H_{\text{mod}} + \delta H_{\text{mod}}$.

This can be done either by a state perturbation $|\psi\rangle \rightarrow |\psi\rangle + \delta|\psi\rangle$ or by a shape deformation $R \rightarrow R + \delta R$.

The modular chaos bound states that the **matrix elements** of $\delta H_{\text{mod}}(s)$ cannot grow under **modular flow** faster than exponentially with **exponent 2π** .

$$\frac{d}{ds} \log |\langle \chi_i | \delta H_{\text{mod}}(s) | \chi_j \rangle| \leq 2\pi \quad \text{for: } s \gg 1$$

$$\text{where: } \delta H_{\text{mod}}(s) = e^{-iH_{\text{mod}}s} \delta H_{\text{mod}} e^{iH_{\text{mod}}s}$$

MODULAR SCRAMBLING MODES

The modular chaos bound picks out a special set of **operators that saturate** the allowed exponential growth.

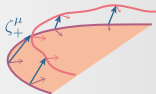
They can be extracted from the **large modular parameter limit** of $\delta H_{\text{mod}}(s)$

$$\begin{aligned} G_+ &= \frac{1}{2\pi} \lim_{s \rightarrow +\Lambda} e^{-2\pi s} \delta H_{\text{mod}}(s) \\ G_- &= -\frac{1}{2\pi} \lim_{s \rightarrow -\Lambda} e^{2\pi s} \delta H_{\text{mod}}(s) \end{aligned}$$

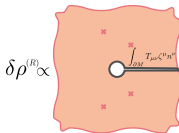
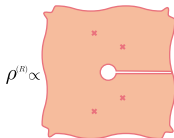
In holographic theories, $\Lambda \sim \log N^2$ and the modular scrambling modes are **stress-tensor lightray operators**

$$G_{\pm} = \int_{-\ell_K}^{\ell_K} dx^{\pm} \int_{RT} d^d y \sqrt{\gamma} \zeta_{(\pm)}^{\mu} (y^i, x^{\pm}) T_{\mu\pm}$$

generating **parallel transport of the RT surface.**



ROAD MAP OF THE DERIVATION



$$\delta H_{\text{mod}}^{(R)} = -\delta \log \rho^{(R)}$$

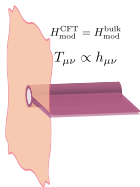
$$= \lim_{b \rightarrow 0} \int_{-\infty}^{+\infty} \frac{ds}{4 \sinh^2(\frac{s+\mu}{2})} \int_{\partial M} \zeta^\mu n^\nu e^{-iH_{\text{mod}} s} : T_{\mu\nu} : e^{iH_{\text{mod}} s}$$

Euclidean path integral
for density matrix

Variation of $\rho = T_{\mu\nu}$
insertion along the cut

Baker-Campbell-Hausdorff

ROAD MAP OF THE DERIVATION



$$\begin{aligned}\delta H_{\text{mod}}^{(R)} &= \int_0^{+\infty} ds e^{-2\pi s} e^{-iH_{\text{mod}}s} T_{+\mu} \zeta^\mu e^{iH_{\text{mod}}s} \\ &\quad - \int_{-\infty}^0 ds e^{2\pi s} e^{-iH_{\text{mod}}s} T_{-\mu} \zeta^\mu e^{iH_{\text{mod}}s} \\ &= \int_0^{\ell_K} dx^+ T_{+\mu} \zeta_{(+)}^\mu - \int_{-\ell_K}^0 dx^- T_{-\mu} \zeta_{(-)}^\mu \\ &\quad + \text{non local}\end{aligned}$$

$$e^{-2\pi|s|} \delta H_{\text{mod}}(s) \sim \int_{-\ell_K}^{\ell_K} dx^\pm T_{\pm\mu} \zeta_{(\pm)}^\mu$$

where: $\ell_K \sim 1/|K_{ij}|\pm|$

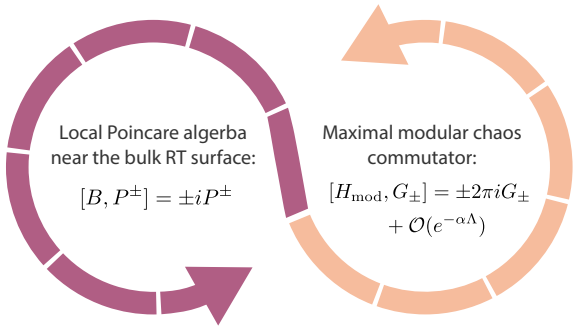
the non-locality scale of the bulk modular H.

Holographic dictionary &
gravitational "Gauss' law"

Split local and non-local
part of modular integral

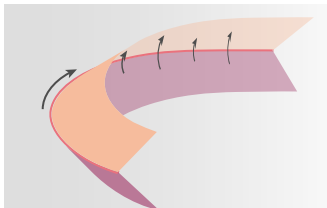
Extract maximally
scrambling piece

POINCARÉ FROM LYAPUNOV



CURVATURE FROM CHAOS COMMUTATORS

$$[G_+, G_-] \in \text{zero mode algebra} \rightarrow \text{large diffeo } \xi^\mu$$



The commutator of scrambling modes is a modular zero mode:

$$[H_{\text{mod}}, [G_+, G_-]] = 0$$

In holographic theories it is an element of the RT surface symmetry algebra that preserve the surface and commute with modular boosts:

$$\xi^\mu \partial_\mu = \omega(y) \epsilon^{\alpha\beta} x_\beta \partial_\alpha + \zeta^i(y) \partial_i + \mathcal{O}(x^2)$$

The **Noether charge** of the large diffeomorphism generated by the **commutator of the scrambling modes** associated to two infinitesimal shape deformations of the modular Hamiltonian **measures the bulk curvature**:

$$Q_{\text{Noether}} = \int_{RT} \sqrt{\gamma} \epsilon_{\alpha\beta} \delta x_1^\alpha \delta x_2^\mu \left(R_\alpha^\mu + K_{\alpha|ij} K^{\alpha|ij} \right)$$

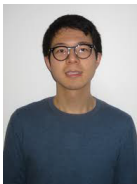
OPEN QUESTIONS

- Can the modular chaos bound be proved in full generality?
- Gravitation as effective theory for scrambling modes?
- Is there a more local version of the story in the bulk?
 - Is modular chaos the right principle for extracting the local Poincare symmetry algebra in an observer's neighborhood?
 - Can we use it to translate ourselves inside a black hole?

IN COLLABORATION WITH



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Dongsheng Ge & L.L.

*A Modular Sewing Kit for
Entanglement Wedges
arXiv: 1903.04493*