

Universal type IIA de Sitter solutions at tree-level

Based on 1003.3590 with Ulf Danielsson and Thomas Van Riet,
and older work: 0812.3551 with Caviezel, Körs, Lüst, Wräse and Zagermann
<http://itf.fys.kuleuven.be/~koerber/talks.html>



Kolymvari, Crete, 16 September 2010

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Motivation I

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Goal: construct a completely explicit and simple dS solution

- Classical in α' , $g_S \rightarrow$ solution of supergravity eoms
- Clear 10d interpretation: compact internal manifold
- Limit ingredients to what is well-understood: avoid e.g. non-geometric fluxes

Motivation II

In this talk:

- Focus on type IIA supergravity
- Derive simple set of conditions on the geometry for obtaining dS solutions
- Present a one-dimensional family of examples on $SU(2) \times SU(2)$

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Do not expect too much!

- Still far from realistic
- In particular: tachyonic modes
- Smeared orientifolds

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But there is more. . .

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- *Flauger, Paban, Robbins, Wrane*: rules out specific models with geometric fluxes; *Caviezel, PK, Körs, Lüst, Tsimpis, Zagermann*: rules out some coset models, retains $SU(2) \times SU(2)$

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Impose **duality** condition

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In IIA: $F_{(0)}$ (Romans mass m), $F_{(2)}$, $F_{(4)}$, $F_{(6)}$, $F_{(8)}$, $F_{(10)}$
Impose **duality** condition
- Two opposite chirality gravitino fields $\Psi_M^{1,2}$ and dilatino fields $\lambda^{1,2}$

Compactification ansatz

- Metric:

$$ds^2 = e^{2A(y)} g_{(4)\mu\nu}(x) dx^\mu dx^\nu + g_{mn}(y) dy^m dy^n ,$$

with $g_{(4)}$ flat Minkowski, AdS_4 metric or dS_4 metric, A warp factor

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- NSNS-flux H: purely internal
- RR-fluxes:

$$F = \sum_l F_l = \hat{F} + e^{4A} \text{vol}_4 \wedge F_{\text{el}} , \quad (F_{\text{el}} = \star_6 \sigma(\hat{F}))$$

where σ reverses the indices of a form

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- Susy implies the existence of nowhere-vanishing internal 6d spinor η_+ (and complex conjugate η_-)

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$$\Omega = \frac{1}{3! \|\eta\|^2} \eta_-^\dagger \gamma_{i_1 i_2 i_3} \eta_+ dx^{i_1} \wedge dx^{i_2} \wedge dx^{i_3}$$

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- J and Ω define SU(3)-structure, not holonomy since generically $dJ \neq 0$, $d\Omega \neq 0$: **geometric flux**
- Susy conditions translate into condition dJ and $d\Omega$
- More general: η^1, η^2 : $SU(3) \times SU(3)$ -structure

$SU(3)$ -structure AdS_4 solutions

Lüst, Tsimpis

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$$dJ = \frac{3}{2} \text{Im}(\mathcal{W}_1 \Omega^*) + \mathcal{W}_4 \wedge J + \mathcal{W}_3$$

$$d\Omega = \mathcal{W}_1 J \wedge J + \mathcal{W}_2 \wedge J + \mathcal{W}_5^* \wedge \Omega$$

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- Geometric flux i.e. non-zero torsion classes:

$$dJ = \frac{3}{2} \text{Im}(\mathcal{W}_1 \Omega^*) \quad \text{with} \quad \mathcal{W}_1, \mathcal{W}_2 \text{ purely imaginary}$$

$$d\Omega = \mathcal{W}_1 J \wedge J + \mathcal{W}_2 \wedge J \quad \Rightarrow \text{put } \mathcal{W}_1 = i\mathcal{W}_1, \mathcal{W}_2 = i\mathcal{W}_2,$$

- Form-fluxes:

AdS_4 superpotential W :

$$H = \frac{2m}{5} e^\Phi \text{Re} \Omega$$

$$e^\Phi \hat{F}_2 = -\frac{\mathcal{W}_1}{4} J - W_2$$

$$e^\Phi \hat{F}_4 = \frac{3m}{10} J \wedge J$$

$$e^\Phi \hat{F}_6 = \frac{9\mathcal{W}_1}{4} \text{vol}_6$$

$$\nabla_\mu \zeta_- = \frac{1}{2} W \gamma_\mu \zeta_+ \quad \text{definition}$$

$$We^{i\theta} = -\frac{1}{5} e^\Phi m - \frac{3i}{4} \mathcal{W}_1$$

$$R_4 = 4\Lambda = -12|W|^2$$

Bianchi identities

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- Leads to

$$e^{2\Phi} m^2 = \mu + \frac{5}{16} (3(W_1)^2 - 2(W_2)^2) \geq 0$$

with $\mu > 0$: net orientifold charge, $\mu < 0$: net D-brane charge

Geometries

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Geometries:

- Nearly-Kähler solutions *Behrndt, Cvetič* $\mathcal{W}_2 = 0$
The only examples known are the homogeneous manifolds:
 $SU(2) \times SU(2)$, $\frac{G_2}{SU(3)} = S^6$, $\frac{Sp(2)}{S(U(2) \times U(1))} = \mathbb{CP}^3$, $\frac{SU(3)}{U(1) \times U(1)}$
- Families of solutions on the above manifolds with $\mathcal{W}_2 \neq 0$
Tomasiello; PK, Lüst, Tsimpis
- $\mu \neq 0$: e.g. Iwasawa manifold

How to construct dS solutions?

Some ideas from the susy AdS solutions seem worth keeping:

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Approximation of constant warp factor & smeared orientifolds

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- Let's try adding $W_3 \rightarrow$ half-flat manifold
- The simplest dS solution we can find: put $W_2 = 0$
- Orientifold sources: like in susy case \rightarrow calibrated

The ansatz II

Geometry: $SU(3)$ -structure with

$$\begin{aligned} dJ &= \frac{3}{2} W_1 \operatorname{Re} \Omega + W_3, \\ d\operatorname{Re} \Omega &= 0, \\ d\operatorname{Im} \Omega &= W_1 J \wedge J. \end{aligned}$$

Put $A = 0$, Φ constant

The ansatz II

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Flux & sources:

$$e^\Phi \hat{F}_0 = f_1$$

$$H = f_5 \operatorname{Re} \Omega + f_6 W_3$$

$$e^\Phi \hat{F}_2 = f_2 J$$

$$e^\Phi j = j_1 \operatorname{Re} \Omega + j_2 W_3$$

$$e^\Phi \hat{F}_4 = f_3 J \wedge J$$

$$e^\Phi \hat{F}_6 = f_4 \operatorname{vol}_6$$

Supergravity equations of motion

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$$d\hat{F}_2 + H\hat{F}_0 = -j, \quad (\text{Bianchi } \hat{F}_2)$$

$$d\star_6 \hat{F}_4 - H \wedge \star_6 \hat{F}_6 = 0, \quad (\text{eom } \hat{F}_4)$$

$$d(e^{-2\Phi} \star_6 H) - (\star_6 \hat{F}_2)\hat{F}_0 - (\star_6 \hat{F}_4) \wedge \hat{F}_2 - (\star_6 \hat{F}_6) \wedge \hat{F}_4 = 0, \quad (\text{eom } H)$$

$$2(R_4 + R_6) - H^2 - e^\Phi \star_6 (\text{Im}\Omega \wedge j) = 0, \quad (\text{eom } \Phi)$$

$$R_4 + e^{2\Phi} \sum_n (\tilde{F}_{(n)}^2) + e^\Phi \star_6 (\text{Im}\Omega \wedge j) = 0, \quad (\text{external Einstein})$$

$$-\tfrac{1}{2}H^2 + \tfrac{1}{4}e^{2\Phi} \sum_n (5-n)\hat{F}_{(n)}^2 + \tfrac{3}{4}e^\Phi \star_6 (\text{Im}\Omega \wedge j) = 0, \quad (\text{trace Einstein/eom } \Phi)$$

$$R_{ij} - \tfrac{1}{2}H_i \cdot H_j - \tfrac{1}{4}e^{2\Phi} \sum_n \left(\hat{F}_{(n)i} \cdot \hat{F}_{(n)j} - \tilde{F}_{(n)i} \cdot \tilde{F}_{(n)j} \right)$$

$$+ \tfrac{1}{4}e^\Phi \left\{ -g_{ij} \star_6 (\text{Im}\Omega \wedge j) + 2\star_6 \left[(g_{k(i} dx^k \wedge \iota_{j)}) \text{Im}\Omega \wedge j \right] \right\} = 0. \quad (\text{Einstein/eom } \Phi)$$

Ricci tensor in terms of $SU(3)$ -structures I

Bedulli, Vezzoni

Decompose R_{ij} $SU(3)$ -reps

$$R_{ij} = \frac{s(R_{ij})}{6}g_{ij} + R_{ij}^+ + R_{ij}^-$$

$s(R_{ij})$ trace

R_{ij}^+ (1,1)-part: **8** of $SU(3)$

R_{ij}^- (2,0)+(0,2)-part: **6 + $\bar{6}$** of $SU(3)$

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Associated primitive two- and three-form to R^+ and R^-

$$P_2(R_{ij}) = \frac{1}{2}J^k{}_i R_{kj}^+ dx^i \wedge dx^j, \quad P_3(R_{ij}) = \frac{1}{2}R_{il}^- \Omega^l{}_{jk} dx^i \wedge dx^j \wedge dx^k$$

Ricci tensor in terms of $SU(3)$ -structures II

$$s(R_{ij}) = \frac{15}{2}(W_1)^2 - \frac{1}{2}(W_2)^2 - \frac{1}{2}(W_3)^2,$$

$$P_2(R_{ij}) = -\frac{1}{4} \star_6 (W_2 \wedge W_2) - \frac{1}{2} \star_6 d \star_6 \left(W_3 - \frac{1}{2} W_1 \text{Re} \Omega \right),$$

$$P_3(R_{ij}) = 2W_1 W_3|_{(2,1)} + 2dW_2|_{(2,1)} - \frac{1}{4} Q(W_3, W_3)$$

with

$$Q(W_3, W_3) = \left(\Omega^{ijk} \iota_j \iota_i W_3 \wedge \iota_k W_3 \right)_{(2,1)},$$

Conditions

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Conditions on torsion classes:

$$W_2 = 0,$$

$$d \star_6 W_3 \propto J \wedge J,$$

$$(W_{3i} \cdot W_{3j})^+ = 0,$$

$$Q(W_3, W_3) = q(W_3)_{2,1}$$

And equations for the constants f_i, j_i in the ansatz

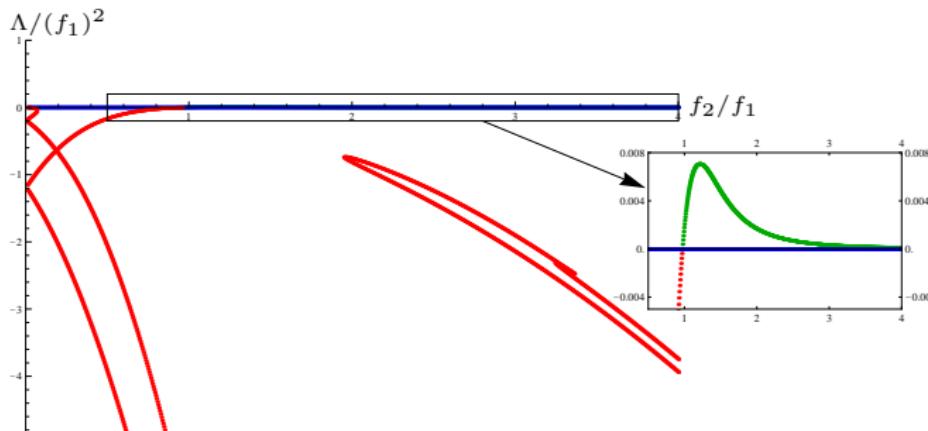
Solutions I

Standard orientifold ansatz sets $q/w_3 = 8/\sqrt{3}$

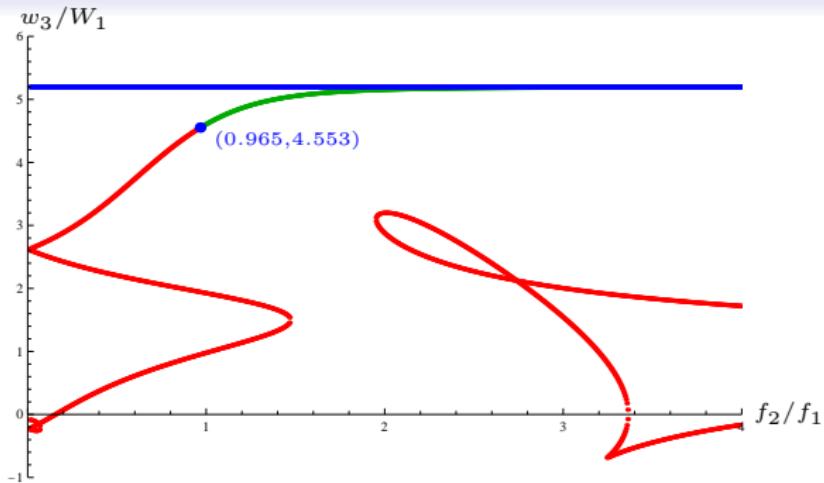
We then only find dS solutions with $f_3 = f_4 = 0$

Parameters: overall scale, dilaton and one extra parameter

Take e.g. f_2/f_1



Solutions II



dS solutions with $4.553 < w_3/W_1 < 3\sqrt{3}$ 

Example: $SU(2) \times SU(2)$

Example on group manifold $SU(2) \times SU(2)$

On group manifold: left-invariant forms e^i :

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Four orientifolds (compatible $SU(3)$ -structure)

Solution I

Solution on $SU(2) \times SU(2)$ reads

$$J = a(e^{16} - e^{24} + e^{35}),$$

$$\text{Re}\Omega = v_1(e^{456} + e^{236} + e^{125}) + (a^6/(v_1)^3) e^{134}$$

$$\text{Im}\Omega = (a^3/v_1)(e^{123} + e^{145} - e^{346}) - ((v_1)^3/a^3) e^{256}$$

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It follows:

$$g = \text{diag}(a^4/(v_1)^2, (v_1)^2/a^2, a^4/(v_1)^2, a^4/(v_1)^2, (v_1)^2/a^2, (v_1)^2/a^2),$$

$$W_1 = \frac{a^6 + v_1^4}{4a^5v_1}, \quad W_2 = 0,$$

$$W_3 = \frac{a^6 - 3(v_1)^4}{8a^5(v_1)^4} [(v_1)^4(e^{456} + e^{236} + e^{125}) - 3a^6e^{134}]$$

$$\Rightarrow w_3 = \pm|W_3| = -\frac{\sqrt{3}(a^6 - 3(v_1)^4)}{4a^5v_1}, \quad q/w_3 = 8/\sqrt{3}$$

Solution II

One easily checks conditions torsion classes:

$$W_2 = 0,$$

$$d \star_6 W_3 \propto J \wedge J,$$

$$(W_{3i} \cdot W_{3j})^+ = 0,$$

$$Q(W_3, W_3) = 8/\sqrt{3} w_3 (W_3)_{2,1}$$

Furthermore

$$w_3/W_1 = \frac{\sqrt{3}(3(v_1)^4 - a^6)}{a^6 + (v_1)^4},$$

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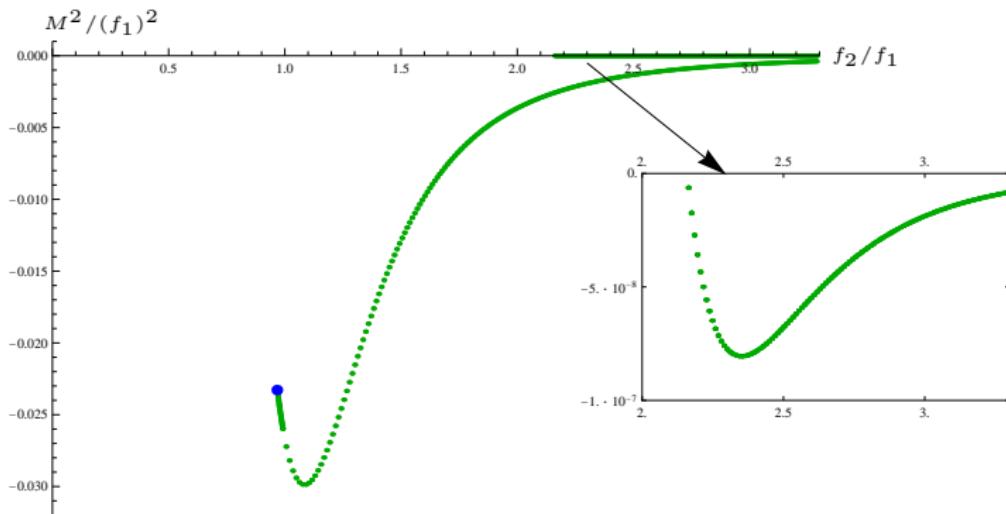
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- $SU(2) \times SU(2)$ seems to know about upper limit $3\sqrt{3}$
- Minkowski solution not present (we would need: $v_1/a^{3/2} \rightarrow \infty$)

Spectrum

Using 4D $N = 1$ supergravity technology we can calculate spectrum of left-invariant fluctuations. Contains tachyonic modes:



Conclusions

- We found simple conditions on geometry for having classical dS solutions
- We found geometry on $SU(2) \times SU(2)$ satisfying these conditions
- Problems: tachyonic modes & smeared orientifolds
- Are tachyons generic? (e.g. *Gómez-Reino, Louis, Scrucca*)
- Fully localized solutions with warp factor?
- Further work: systematic scan manifold & include $W_2 \neq 0$
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The end... The end... The end...